

Research Plan for Spin Physics at RHIC

Abstract

In this report we present the research plan for the RHIC spin program. The report covers 1) the science of the RHIC spin program in a world-wide context; 2) the collider performance requirements for the RHIC spin program; 3) the detector upgrades required, including timelines; 4) time evolution of the spin program.

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1 Executive Summary

An action item from the June 30-July 1, 2004 DOE Office of Nuclear Physics Science and Technology Review of the Brookhaven National Laboratory (BNL) Relativistic Heavy Ion Collider (RHIC) written Report, dated September 13, 2004, was that *"BNL should prepare a document that articulates its research plan for the RHIC spin physics program. A copy should be submitted to DOE by January 31, 2005."* This document is submitted to the DOE Office of Nuclear Physics on behalf of the Laboratory, in response to that action item.

We provide here a plan that addresses: 1) the science of the RHIC spin program in a world-wide context; 2) the collider performance requirements for the RHIC spin program; 3) the detector upgrades required, including timelines; 4) time evolution of the spin program. The RHIC Spin Plan Group was charged to formulate the plan by Thomas Kirk, BNL Associate Director for High Energy and Nuclear Physics.

The importance of the study of nucleon spin to nuclear physics and its co-evolution with heavy ion physics in the RHIC facility is discussed in the first section of this report. Here, we note that, of all the basic properties of the nucleons (protons and neutrons), the intrinsic spin angular momentum aspects are least well known. In particular, the contribution of the gluons to the net spin of the nucleons, as well as the role of orbital angular momentum among the quarks and gluons is almost totally unmeasured. Our plan will argue that the RHIC facility will provide the most comprehensive and precise measurements possible in the present and foreseen worldwide program of spin physics and will be able to accomplish this in a proposed program of 10 weeks (on average) of data taking per year through the year 2012. This scenario is labeled "the technically limited BNL spin plan". If budgetary constraints on the spin program evolution are imposed, the program will be stretched out accordingly and the results will be available later and at higher total cost.

The RHIC 100-250 GeV colliding beams facility of highly polarized protons, with arbitrary spin orientation at two interaction points around the RHIC ring (at the PHENIX and STAR detectors), is unique in the world and likely to remain so for the foreseeable future. The maximum polarization in each beam is expected to rise from 45% to 70% in the 2005-2006 time frame. Exploiting this powerful accelerator and experimental detector capability, together with a steadily improving average luminosity at the interaction points, a unique and important program of spin physics measurements can be completed over the time period 2005-2012. Determination of the gluon's contribution to the proton spin will be an early achievement of the program, followed by the direct measurement of the quark and anti-quark contributions to the proton spin, sorted by quark flavor. These results will be achieved in a later phase of the spin physics program via W-boson production measurements. In the same time frame, the power of lattice gauge physics, carried out on BNL supercomputers, is expected to be able to relate the spin results to the predictions of Quantum Chromo-Dynamics, now believed to be the fundamental theory of the strong nuclear interactions.

We summarize our findings on the four areas in the charge:

Science: Gluon polarization will be measured at RHIC using several independent methods: π^0 , jet, direct γ and γ +jet, and heavy quark production. Results from the different methods will overlap to allow us to test our understanding of the processes involved and expand the range

of momentum fraction for the measurements. We want to learn both the average contribution to the proton spin of the gluons as well as a detailed map. We use first the higher cross section processes, π^0 and jet production and, as we reach higher luminosity and polarization, the clean but rarer process of direct γ production. We plan to emphasize these measurements for $\sqrt{s}=200$ GeV collisions from 2005-2008. At that time, we expect to have reached a precision that can clearly distinguish between zero gluon polarization and a minimal ("standard") gluon polarization. A large gluon polarization, consistent with the gluon carrying most of the spin of the proton, would be precisely measured. In this period we will also pursue the question of the transverse spin structure. Gluons, massless spin 1 particles, cannot contribute to the transverse spin. Large transverse spin asymmetries have been seen for deeply inelastic scattering (DIS) and now for RHIC, so this topic is also a potential window into a new understanding of the structure of the nucleons.

Production of W bosons, the carrier of the weak interaction, has an inherent handedness. At RHIC we plan to use this "parity violation" signal to directly measure the polarization of the quark and anti-quark that form the W boson. To do this we will run at the top RHIC energy, $\sqrt{s}=500$ GeV. This will provide the first direct measurements of anti-quark polarization in the proton, with excellent sensitivity. We plan to begin these measurements in 2009. The W measurements will require the prior completion of certain detector improvements for both PHENIX and STAR.

The RHIC spin results will provide precise measurements of gluon and anti-quark polarization. With these results we will also understand the role of the remaining contributor to the proton spin, orbital angular momentum. We will have also explored our understanding of the interconnected results from the different RHIC spin probes, and from the DIS measurements. The sensitivity at RHIC for gluon polarization is shown in Figure 1, where we also include the sensitivity for the ongoing DIS experiment at CERN, which measures gluon polarization by the production of hadrons. From the figure, we see that RHIC will provide precise results over a large range in momentum fraction, characterizing the gluon contribution to the proton spin.

Figure 1: $\Delta G/G(x)$ vs. $\log x$ with a model showing the x range for various RHIC processes with expected uncertainties; 200 and 500 GeV data are shown with experimental uncertainties in RHIC and COMPASS for $Q^2 \geq 1$.

The sensitivity of RHIC for anti-quark polarization is shown in Figure 2. We will measure the \bar{u} and \bar{d} anti-quark polarizations to about 0.01, as well as u and d quark polarizations. The measurement is direct and very clean, using parity violating production of W bosons. DIS measurements also yield the anti-quark polarizations. The method has the disadvantage of theoretical uncertainties in modeling the fragmentation and the advantage that the method is accessible today. The RHIC and DIS methods probe the proton at very different distances, (Q^2 range) where RHIC corresponds to X Fermi and DIS to Y Fermi, compared to the proton radius of 1 Fermi. The theory of QCD prescribes how to connect the results from different probing distances—the description of unpolarized DIS results over a very large distance range is one of the major successes of QCD. Both the anti-quark and the gluon results from RHIC and DIS test the QCD assumption of universality, namely that the physics for both proton and lepton processes can be described with the same underlying quark, anti-quark, and gluon distributions.

Figure 2: $\Delta q/q(x)$ vs. $\log x$ with a model for u , d , \bar{u} and \bar{d} ; with RHIC and DIS expected

uncertainties.

We emphasize that the planned RHIC program will make major contributions to our understanding of matter. Our results will complement the DIS measurements, completed and planned. We include in our report expectations from a next stage of DIS-colliding polarized electrons with polarized protons and neutrons which probes still further into the structure of matter. As we develop theoretical tools to apply QCD to understand this structure, these spin results will provide a deep test of our understanding of the fundamental building blocks of matter.

Collider Performance Requirements: The spin physics program requires RHIC beams with high polarization, and high integrated luminosity. For the sensitivities assumed above, we have used $P=0.7$ and luminosity 300 pb^{-1} at $\sqrt{s}=200 \text{ GeV}$ and 800 pb^{-1} at $\sqrt{s}=500 \text{ GeV}$. (Note that this refers to "delivered" luminosity, while the sensitivities in this report use the luminosity expected to be recorded by each experiment.)

The polarization level is presently $P=0.45$ and is expected to reach 70% polarization by 2006. This improvement is expected using new Siberian Snake Magnets installed in the AGS in 2004 and 2005. The average luminosity at store must be increased by a factor 15 to reach the integrated luminosity goals for three years of running at 10 physics weeks per year. To achieve this level of performance will require completion of the planned vacuum improvements in RHIC, expected for 2007. The luminosity increase then comes from reaching a bunch intensity of 2×10^{11} . A limit will be caused by beam-beam interactions that change and broaden the betatron tune of the machine, moving part of the beam into a beam resonance region where beam is then lost. Work in 2004 discovered a new betatron tune for RHIC that greatly improves losses from the beam-beam interaction. RHIC at our luminosity goal will operate near or above previously achieved beam-beam parameters, and will be the first hadron collider in the strong-strong beam-beam regime.

Reaching these goals requires "learning by doing". We plan to study limits and develop approaches to improve the polarization and luminosity during physics runs by including beam studies for one shift per day. It is also important that a sustained period of running and development be followed, if possible each year. It is this approach that has led to the major improvements for heavy ion luminosity and to our improvements to this date in polarization and proton luminosity.

Experimental Detector Capabilities: The PHENIX and STAR detectors are complete for the gluon polarization program. Improvements to both detectors are required to carry out the W physics program. Both experiments also plan upgrades that benefit both the heavy ion and spin programs, significantly extending the range of physics probes for spin.

W Program Upgrades–PHENIX The present online event selection for muons, the channel used for W physics, will need to be improved for the W luminosity. New resistive plate chambers (RPC) are being proposed to provide the tighter event selection, along with electronics changes to the muon tracking readout. The RPC proposal was submitted to NSF in January 2005, with a cost estimate of \$1.8M. The tracking readout proposal has been submitted to the Japan Society for the Physical Sciences, with a cost of \$1.0M. The planned timeline is to complete both for the 2008 run.

W Program Upgrades–STAR New tracking for forward electrons from W decay is necessary for the W program. It is planned to propose this upgrade in 2006 to DOE, with an estimated cost of \$5M, although research and development on the technology (GEM detectors) is proceeding

and the cost estimate is rough at this time. The forward tracking detector is to be completed for the 2010 run, with part of the detector in place earlier.

Heavy-Ion/Spin Program Upgrades PHENIX plans a barrel micro vertex detector that provides access to heavy quark states and to jet physics based on tracking. The heavy quark data will add a new probe for gluon polarization at lower momentum fraction (shown on Fig. 1). The jet information will be used in correlated (γ +jet) measurements, which better determine the sub-process kinematics for gluon polarization measurement. A second upgrade being planned is to change the brass "nose cones", used as a filter for the muon arms, to active calorimeters that will measure photons, π^0 and jet energy. The nose cone calorimeters would provide a larger momentum fraction range for the gluon polarization measurements. Both are important upgrades for the heavy ion physics program. The vertex detector is planned for the 2008 run, and the nose cone calorimeter proposal is being developed now.

STAR has proposed expanded forward calorimeters to NSF in January 2005. The calorimeters will measure the gluon density for proton-gold collisions and will also provide very significant spin measurements. With the calorimeters, forward π^0 , direct γ , and jet events can be observed, yielding sensitivity to gluon polarization at lower momentum fraction, as shown on Fig. 1. A second upgrade driven by the heavy ion program, a barrel inner tracker, will give access to heavy quark measurements for spin. The forward calorimeters are to be in place for the 2007 run. The barrel inner tracker is to be completed for the 2010 run.

Time Evolution: In order to indicate the pace of spin program evolution under different budget assumptions, we provide two examples with an average of 10 weeks and 5 weeks of spin physics data taking per year. We note that the actual running plan will continue to be developed, year-by-year, from the experiment beam-use proposals. The time evolution cases provided here indicate paths for reaching the spin program's strategic goals under two example scenarios.

We show in Fig. 3, the impact of 10 and 5 averaged spin physics weeks per year. The "technically-driven schedule" represents the luminosity used for the sensitivities shown in the figures above. With 10 weeks per year, we achieve the $\sqrt{s}=200$ GeV target in 3 years, assuming that we successfully climb the accelerator learning curve to reach the target store luminosity. The 500 GeV running target is also expected to be achieved in 3 years (there is a natural luminosity improvement for 500 GeV of a factor of 2.5 over 200 GeV from the smaller cross section beams).

With 10 averaged spin physics weeks per year, technically driven, our proposed target sensitivities can be reached running at $\sqrt{s}=200$ GeV from 2005-2008, and at $\sqrt{s}=500$ GeV from 2009-2012. We have assumed that 2009 will be an engineering run for W physics (new detector systems and higher luminosity), with a full detector in place for PHENIX and a partial detector for STAR. 2010-2012 will provide full-functionality W physics runs at high luminosity for both detectors. This is the preferred scenario from BNL.

As can be seen in Fig. 3, running 5 averaged spin physics weeks per year (we have interpreted this as running 10 spin physics weeks every two years to reduce end effects), each program, 200 GeV and 500 GeV, takes more than 6 years. Under this scenario RHIC would run roughly 25% of the year and both the heavy ion and spin programs would be stretched a factor of greater than two in calendar time. We hope that such a slow and relatively inefficient program can be avoided.

Fig. 3: pp luminosity projections for 10 and 5 physics weeks per year (5=10/2).

2 The science case for RHIC Spin

Spin is one of the most fundamental concepts in physics, deeply rooted in Poincaré invariance and hence in the structure of space-time itself. All elementary particles we know today carry spin, among them the particles that are subject to the strong interactions, the spin-1/2 quarks and the spin-1 gluons. Spin, therefore, plays a central role also in our theory of the strong interactions, *Quantum Chromodynamics (QCD)*, and to understand spin phenomena in QCD will help to understand QCD itself. To contribute to this understanding is the primary goal of the spin physics program at RHIC.

It is a remarkable property of QCD, known as *confinement*, that quarks and gluons are not seen in isolation, but only bound to singlet states of the strong “color” charge they carry. At the heart of investigating confinement in QCD is the study of the inner structure of strongly-interacting particles in nature that are composed of quarks and gluons. Among these, the proton and neutron are clearly special as they make up all nuclei and hence most of the visible mass in the universe. Their detailed study is therefore of fundamental interest. The proton and neutron also carry spin-1/2, which immediately brings the central role of spin in nucleon structure to the fore. It is worth recalling that the discovery of the fact that the proton has structure— and hence really the birth of strong interaction physics— was due to spin, through the measurement of a very unexpected “anomalous” magnetic moment of the proton by O. Stern and collaborators in 1933 [1]. Today, after decades of ever more detailed studies of nucleon structure, a prime question is how the proton spin-1/2 is composed of the average spins and orbital angular momenta of quarks and gluons inside the proton. Polarization has become an essential tool in the investigation of the strong interactions through nucleon structure.

Quarks were originally introduced simply based on symmetry considerations [2], in an attempt to bring order into the large array of strongly-interacting particles observed in experiment. In order to satisfy the Pauli exclusion principle for baryons such as the Δ^{++} or the Ω^- which are made up of three quarks of the same flavor, the spin-1/2 quarks had to carry a new quantum number [3], later termed “color”. A modern rendition of Rutherford’s experiment has shown us that quarks are real. This experiment is the deeply-inelastic scattering (DIS) of electrons (or, later, muons) off the nucleon, a program that was started in the late 1960’s at SLAC [4]. A high-energy electron interacts with the nucleon, via exchange of a highly virtual photon. For virtuality of $\sqrt{Q^2} > 1$ GeV distances < 0.2 fm are probed in the proton. The proton breaks up in the course of the interaction. The early DIS results compelled an interpretation as elastic scattering of the electron off pointlike, spin-1/2, constituents of the nucleon [5, 6], carrying fractional electric charge. These constituents, called “partons” were subsequently identified with the quarks. The existence of gluons was proved indirectly from a missing $\sim 50\%$ contribution [7] to the proton momentum not accounted for by the quarks. Later on, direct evidence for gluons was found in three-“jet” production in electron-positron annihilation [8]. From observed angular distributions of the jets it became clear that gluons have spin one [9].

The so successful parton interpretation of DIS assumed that partons are practically free (i.e., non-interacting) on the short time scales set by the high virtuality of the exchanged photon. This

implied that the underlying theory of the strong interactions must actually be relatively weak on short time or, equivalently, distance scales [10]. In a groundbreaking development, Gross, Wilczek and Politzer showed in 1973 that the non-abelian theory “QCD” of quarks and gluons, which had just been developed a few months earlier [11], possessed this remarkable feature of “asymptotic freedom” [12], a discovery for which they were awarded the 2004 Nobel Prize for Physics. The interactions of partons at short distances, while weak in QCD, were then predicted to lead to visible effects in the experimentally measured DIS structure functions known as “scaling violations” [13]. These essentially describe the response of the partonic structure of the proton to the resolving power of the virtual photon, set by its virtuality Q^2 . It has arguably been *the* triumph of QCD that the predicted scaling violations have been observed experimentally and verified with great precision. Deeply-inelastic scattering thus paved the way for our theory of the strong interactions, QCD.

Over the following two decades or so, studies of nucleon structure became ever more detailed and precise. Partly this was due to increased luminosities and energies of lepton machines, culminating in the HERA ep collider. Also, hadron colliders entered the scene. It was realized that, again thanks to asymptotic freedom, the partonic structure of the nucleon seen in DIS is universal in the sense that it can also be studied in very inelastic reactions in proton-proton scattering [14, 15, 16]. This offered the possibility to learn about other aspects of nucleon structure (and hence, QCD), for instance about its gluon content which is not primarily accessed in DIS. Being known with more precision, nucleon structure also became a tool in the search of new physics, the outstanding example perhaps being the discovery of the W^\pm and Z bosons at CERN’s Sp \bar{p} S collider [17]. The Tevatron collider today and LHC in the near future are continuations of this theme.

A further milestone in the study of the nucleon was the advent of *polarized* electron beams in the early seventies [18]. This later on allowed to perform DIS measurements with *polarized* lepton beam and nucleon target [19], offering for the first time the possibility to study whether for example quarks and antiquarks have on average preferred spin directions inside a spin-polarized nucleon. The program of polarized DIS has been continuing ever since and has been an enormously successful branch of particle physics. Its single most important result is the finding that quark and antiquark spins provide very little – only about $\sim 20\%$ – of the proton spin [20, 21]. In parallel, starting from the mid 1970’s, there also was a very important line of research on polarization phenomena in hadron-hadron reactions in fixed-target kinematics. In particular, unexpectedly large single-transverse spin asymmetries were seen [22] which, as will be discussed later, may tell us about further fundamental spin-related properties of the nucleon, but have defied a complete understanding in QCD so far.

In the context of the exploration of nucleon structure achieved so far, it is clear that the RHIC spin program is the logical continuation. Very much in the spirit of the unpolarized hadron colliders in the 1980’s, RHIC enters the field to start from where polarized DIS has taken us so far. Here, too, asymptotic freedom of QCD, accessible because of the high energy of RHIC’s polarized beams, is the tool to investigate the partonic structure of the proton. Experiments with polarization at RHIC will probe the proton spin in new profound ways [23], complementary to polarized DIS. We will learn about the polarization of gluons in the proton and about details of the flavor structure of the polarized quark and antiquark distributions. RHIC will probe the structure of transversely polarized protons, and it will unravel the origin of the transverse-spin asymmetries mentioned above. RHIC will also investigate polarization phenomena in high-energy *elastic*

scattering of protons, an equally uncharted area of QCD. Finally, if circumstances are very favorable, knowledge gathered about the spin structure of the proton could conceivably be used to turn RHIC into a discovery machine for New Physics, or a machine that probes the *chiral* structure, inaccessible in unpolarized pp collisions, of new interactions possibly to be found at the LHC.

The field of nucleon structure thrives on the complementarity of information obtained in lepton-nucleon and nucleon-nucleon scattering. We shall see examples of this throughout this report. We see a collider with polarized electrons and protons as the next logical step after RHIC in our quest to explore the spin structure of the nucleon and spin phenomena in QCD.

After a brief review of where we currently stand in this field, the subsequent sections will address the most exciting aspects of the RHIC spin physics program in more detail.

2.1 Synopsis of results from polarized DIS

Spin physics at RHIC has been motivated by the exciting results from the experimental program on polarized DIS over the last ~ 30 years [21]. Most of the DIS measurements were performed with longitudinal polarization of the lepton beam and the nucleon target. The difference of cross sections for the case where the lepton and nucleon have aligned spins or opposite spins then gives access to the spin-dependent structure function $g_1(x, Q^2)$ of the nucleon. Here Q^2 is as before the virtuality of the exchanged photon, and x is the Bjorken variable, $x = Q^2/(2P \cdot q)$ with P and q the nucleon and photon momenta, respectively. The left part of Fig. 1 shows a recent compilation [24] of the world data on $g_1(x, Q^2)$. Information from both proton and neutron targets is available. The importance of g_1 lies in the fact that it has a simple interpretation in the parton model, equivalent to considering the lepton-nucleon interaction as a scattering of polarized leptons off polarized free partons. In the parton model, and including the dominant part of the QCD scaling violations mentioned above, g_1 may be written as

$$g_1(x, Q^2) = \frac{1}{2} \sum_q e_q^2 [\Delta q(x, Q^2) + \Delta \bar{q}(x, Q^2)] . \quad (1)$$

Here the Δq , $\Delta \bar{q}$ are the helicity distribution functions of quarks and antiquarks in the nucleon. For example,

$$\Delta q(x, Q^2) = q^+(x, Q^2) - q^-(x, Q^2) \quad (2)$$

counts the number densities of quarks with the same helicity as the nucleon, minus opposite. It contains information on the *spin* structure of the proton. The kinematic Bjorken variable x is identified with the proton momentum fraction carried by the struck quark. The Q^2 -dependence of the parton distributions is precisely the dependence on the “resolving power” mentioned earlier, quantitatively predictable in QCD perturbation theory, thanks to asymptotic freedom. It is also known as Q^2 -“evolution” of the parton distributions [25]. Physically, it expresses the fact that as Q^2 increases one has higher resolution of the partons, so that it is more likely that a struck quark has radiated one or more gluons so that it is effectively resolved into several partons, each with lower momentum fraction. Similarly, a struck quark may have originated from a gluon splitting into a quark-antiquark pair. This picture explains another remarkable feature of the DIS scaling violations: the Q^2 -dependence of the quark densities, and hence of the structure function $g_1(x, Q^2)$, is partly driven by the gluon density in the proton, despite the fact that the gluon

density does not appear in Eq. (1). The polarized “helicity” gluon density is defined in analogy with Eq. (2) as

$$\Delta g(x, Q^2) = g^+(x, Q^2) - g^-(x, Q^2) . \quad (3)$$

Thus scaling violations in polarized DIS allow, in principle, to determine not only the $\Delta q + \Delta \bar{q}$ combinations for various flavors, but also Δg .

Extensive analyses of polarized-DIS data in terms of the polarized parton distributions have been performed by several groups, taking into account a “state-of-the-art” theoretical framework that includes additional non-leading (“higher-order”) corrections to the framework described above [26, 27, 28, 29, 30]. A typical result, taken from [26], is shown in the right part of Fig. 1. The results refer to a Q^2 scale of 1 GeV², which is a typical scale from which perturbative evolution as described above could be used to calculate the distributions at higher Q^2 . The first panel shows the sum of all polarized quark and antiquark distributions distribution,

$$\Delta \Sigma = \Delta u + \Delta \bar{u} + \Delta d + \Delta \bar{d} + \Delta s + \Delta \bar{s} \quad (4)$$

as a function of x . As can be seen, it is known to a fair accuracy, except at the lower x , where we have shown a typical present uncertainty by a shaded band. The right-hand plot displays the polarized gluon density Δg . Evidently, we know very little, if anything, about gluon polarization in the nucleon. The latter result is not surprising: as we pointed out earlier, the only information on Δg from polarized DIS comes from scaling violations. Since all experiments performed so far have been with fixed targets, the available energy, and hence the reach in Q^2 , have been very limited, resulting in a virtually unconstrained Δg .

It has been possible to extrapolate the results shown in Fig. 1 to $x \rightarrow 0$. There are important insights into nucleon structure that could be gained from this. First of all, there is a venerable sum rule by Bjorken [31] – that actually predates QCD – which remarkably relates the integrals over all x of the high- Q^2 DIS polarized structure functions for the proton and the neutron to the decay constant $g_A \approx 1.273$ in low-energy neutron β -decay:

$$\int_0^1 dx (g_1^p(x, Q^2) - g_1^n(x, Q^2)) = \frac{1}{6} g_A + \mathcal{O}(\alpha_s(Q^2)) , \quad (5)$$

where we have indicated that there are perturbative-QCD corrections to the relation, known to very high accuracy. This sum rule, which was the original motivation for performing measurements in polarized DIS, has been verified experimentally at the 10% level [21].

Using further information from baryon β -decays, it was also possible to determine the x -integral over the combination $\Delta \Sigma$ shown in Fig. 1. This has resulted in one of the most renowned – and debated – results in recent Nuclear and Particle Physics [20]. The importance of the integral of $\Delta \Sigma$, also known as the nucleon “axial charge”, lies in the fact that it yields the average of all *quark and antiquark helicity contributions to the proton helicity*:

$$\langle S_q \rangle = \frac{1}{2} \int_0^1 \Delta \Sigma(x, Q^2) dx . \quad (6)$$

This follows from the definition of the spin-dependent quark distribution functions in Eq. (2); the factor 1/2 is because quarks carry spin-1/2. Experimentally [21],

$$\langle S_q \rangle \approx 0.1 , \quad (7)$$

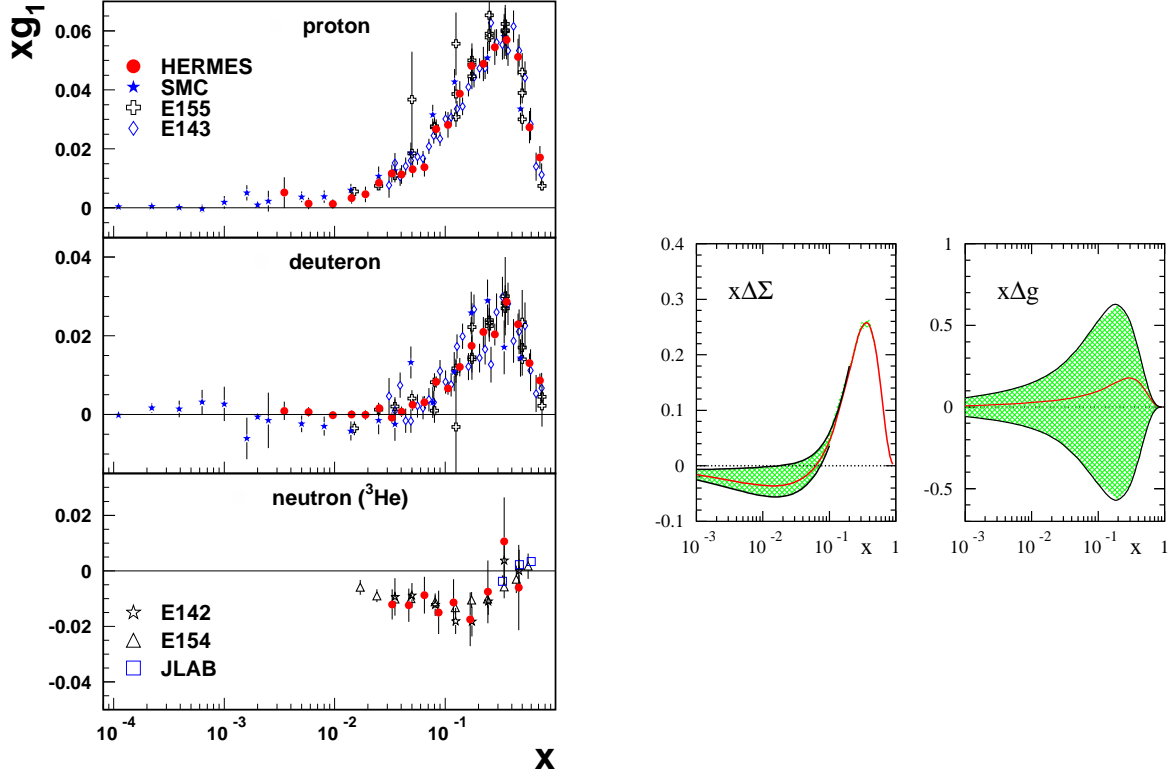


Figure 1: Left: data on the spin structure function g_1 , as compiled and shown in [24]. Right: results from an analysis [26] of polarized DIS in terms of spin-dependent nucleon parton densities $\Delta\Sigma(x, Q^2)$ and $\Delta g(x, Q^2)$ at $Q^2 = 1 \text{ GeV}^2$. $\Delta\Sigma$ is the total quark and antiquark helicity distribution (see Eq. (4)) and Δg the gluon helicity distribution defined in Eq. (3). The shaded bands indicate typical current uncertainties in our knowledge of the distributions.

with an error of about 50%. Despite its large error, the fact that $\langle S_q \rangle \ll 0.5$ implies that very little of the proton spin is carried by that of the quarks. This result is in striking contrast with predictions from constituent quark models and has therefore been dubbed “proton spin crisis”. Even though the identification of nucleon with parton helicity is not a prediction of QCD, such models have enjoyed success in describing hadron magnetic moments and spectroscopy. In any case, polarized DIS teaches us that we must look elsewhere for the proton spin!

2.2 Compelling questions in spin physics

The results from polarized inclusive DIS clearly called for further investigation of the nucleon spin. What are first of all the other candidates for carrying the nucleon spin? An examination of angular momentum in QCD equates the spin-1/2 of the proton by contributions from quark spins, gluon spins, and quark and gluon orbital angular momenta [32, 33, 34]:

$$\frac{1}{2} = \langle S_q \rangle + \langle S_g \rangle + \langle L_q \rangle + \langle L_g \rangle . \quad (8)$$

We have suppressed a dependence of each of the terms on the resolution scale Q^2 . The gluon spin contribution is directly obtained from the gluon helicity distribution in Eq. (3):

$$\langle S_g \rangle(Q^2) = \int_0^1 \Delta g(x, Q^2) dx . \quad (9)$$

Equation (8) motivates a substantial part not only of RHIC spin physics, but of virtually all major current activities in the field of high-energy spin physics. More specifically, the compelling questions are:

How do gluons contribute to the proton spin? There are good reasons to be interested in $\Delta g(x, Q^2)$. First of all, its integral could well be a major contributor to the proton spin. In fact, it is a remarkable feature of QCD that at momentum scales relevant to RHIC physics $\langle S_g \rangle(Q^2)$ may well *be* significant, perhaps even large compared to the “1/2” on the right-hand-side of Eq. (8). The reason is that the integral of $\Delta g(x, Q^2)$ evolves as $1/\alpha_s(Q^2)$ [35], that is, rises logarithmically with Q . This peculiar evolution pattern is a very deep prediction of QCD, related to its so-called axial anomaly. It has inspired ideas that a reason for the smallness of the quark spin contribution should be sought in a “shielding” of the quark spins due to a particular perturbative part of the DIS process $\gamma^* g \rightarrow q\bar{q}$ [35]. The associated contributions arise only at order $\alpha_s(Q^2)$; however, the peculiar evolution of $\langle S_g \rangle(Q^2)$ would compensate this suppression. To be of any practical relevance, such models would require a large positive gluon spin contribution, $\langle S_g \rangle > 1$, even at low “hadronic” scales of a GeV or so. A very large polarization of the confining fields inside a nucleon, even though suggested by some nucleon models [36], would be a very puzzling phenomenon and would once again challenge our picture of the nucleon. Subsection 2.4.1 will discuss in detail the efforts being made at RHIC to address the questions related to Δg , and the prospects for the planned measurements.

What are the patterns of up, down, and strange quark and antiquark polarizations? As is evident from Eq. (1), polarized DIS has given us access to the combinations $\Delta q + \Delta \bar{q}$. We have already discussed one particularly interesting combination of these, $\Delta \Sigma$. To really understand the proton helicity structure in detail, one needs to learn about the various quark and antiquark densities, $\Delta u, \Delta \bar{u}, \Delta d, \Delta \bar{d}, \Delta s, \Delta \bar{s}$, *individually*. This also provides an important additional test of the smallness of the quark spin contribution, independent of the additional input from baryon β -decays necessary so far. It is also important for models of nucleon structure which generally make clear qualitative predictions about, for example, the flavor asymmetry $\Delta \bar{u} - \Delta \bar{d}$ in the proton sea [37]. These predictions are often related to fundamental concepts such as the Pauli principle: since the proton has two valence- u quarks which primarily spin along with the proton spin direction, $u\bar{u}$ pairs in the sea will tend to have the u quark polarized opposite to the proton. Hence, if such pairs are in a spin singlet, one expects $\Delta \bar{u} > 0$ and, by the same reasoning, $\Delta \bar{d} < 0$. Such questions become all the more exciting due to the fact that rather large *unpolarized* asymmetries $\bar{u} - \bar{d} \neq 0$ have been observed in DIS and Drell-Yan measurements [38]. Further fundamental questions concern the strange quark polarization. The polarized DIS measurements point to a sizable negative polarization of strange quarks, in line with other observations of significant strange quark effects in nucleon structure. Recently, in the unpolarized case the asymmetry between strange and antistrange distributions has attracted much attention [39], due to its interest for nucleon models, but also due to its possible implications for an explanation of the $\sim 3\sigma$ “anomaly” in the NuTeV measurement [40] of the Weinberg angle. A measurement of the difference between strange and antistrange polarizations, while probably lying in the future, might give further insights. In subsection ?? we present the possibilities RHIC offers for disentangling

the various flavor polarizations in the nucleon.

What orbital angular momenta do partons carry? Equation (8) shows that quark and gluon orbital angular momenta are the other candidates for the carriers of the proton spin. Consequently, theoretical work focused also on these in the years following the discovery of the “spin crisis”. A conceptual breakthrough was made in the mid 1990s when it was realized [33] that a particular class of “off-forward” nucleon matrix elements, in which the nucleon in the initial and final state has different momentum, measure total parton angular momentum. Put simply, orbital angular momentum is $\vec{r} \times \vec{p}$, with \vec{r} a derivative with respect to momentum transfer in Quantum Mechanics. Thus, in analogy with the measurement of the Pauli form factor it takes a finite momentum transfer on the nucleon to access matrix elements with operators containing a factor \vec{r} . It was also shown how these “off-forward” distributions, really generalizations of the ordinary parton distributions, may be experimentally determined from certain rare exclusive processes in lepton-nucleon scattering, the prime example being “Deeply-Virtual Compton Scattering (DVCS)” $\gamma^* p \rightarrow \gamma p$ [33]. A major emphasis in current and future experimental activities in lepton scattering is on the DVCS and related reactions. There are other observables that are related to orbital angular momentum of nucleon constituents [41]. The Pauli form factor is one of them. Another, accessible in proton-proton scattering, may contribute to spin asymmetries measured with a single transversely polarized proton and an unpolarized one. This brings us to the next compelling question.

What is the role of transverse spin in QCD? So far, we have only considered the helicity structure of the nucleon, that is, the partonic structure we find when we probe the nucleon when its spin is aligned with its momentum. High-energy protons may also be studied when *transversely* polarized, and it has been known for a long time now that very interesting spin effects are associated with this in QCD. Partly this is known from theoretical studies which revealed that besides the helicity distributions Δf discussed above, for transverse polarization there is a new set of parton densities, called “transversity” [42, 43]. They are defined analogously to Eq. (2), but now for transversely polarized partons polarized along or opposite to the transversely polarized proton. Nothing is known so far experimentally about the transversity densities. Their measurement is highly desirable, for a number of reasons. Not only does transversity complete the set of nucleon parton distributions. Differences between the helicity and transversity densities give information about relativistic effects in the nucleon [43]. The transversity densities also give the nucleon tensor charge [43, 44], which is equally fundamental as its axial charge mentioned earlier. Finally, transversity also plays a role in predictions for the neutron electric dipole moment. We will discuss transversity and the prospects for its measurement at RHIC in more detail in section 2.7. The other reason why transverse spin has captured the attention of researchers in QCD for a long time is related to experimental observations of very large single-transverse spin asymmetries in pp scattering [22], where really none were expected. Related azimuthal asymmetries were seen in lepton scattering [45]. Often, when simple expectations are refuted experimentally, new insights emerge, and this has been no different in this case. With time it was realized that single-spin asymmetries may tell us many more things about QCD and the nucleon than anticipated. Particularly interesting examples are parton orbital angular momenta [46] and the color Lorentz force inside a polarized nucleon [47]. We are, however, still far from a complete understanding of all mechanisms that may be involved in single-spin asymmetries. We will show in more detail in section 2.7 that RHIC is poised to provide answers.

We now turn in more detail to the various physics topics relevant at RHIC. We start by a brief

description of the theoretical underpinnings for the theoretical description of “deeply inelastic” hadronic reactions, considering unpolarized scattering for simplicity. We then discuss how polarized pp scattering at RHIC addresses the compelling questions in spin physics raised above. We will also describe other exciting physics opportunities the RHIC spin program offers.

2.3 Unpolarized pp scattering

The basic concept that underlies most of RHIC spin physics is the factorization theorem [16]. It states that large momentum-transfer reactions may be factorized into long and short-distance contributions. The long-distance pieces contain information on the structure of the nucleon in terms of its distributions of constituents, “partons”. The short-distance parts describe the hard interactions of these partons and can be calculated from first principles in QCD perturbation theory. While the parton distributions describe universal properties of the nucleon, that is, are the same in each reaction, the short-distance parts carry the process-dependence and have to be calculated for each reaction considered.

As an explicit example, we consider the cross section for the reaction $pp \rightarrow \pi(p_T)X$, where the pion is at high transverse momentum p_T , ensuring large momentum transfer. X denotes an arbitrary hadronic final state. The statement of the factorization theorem is then:

$$d\sigma = \sum_{a,b,c} \int dx_a \int dx_b \int dz_c f_a(x_a, \mu) f_b(x_b, \mu) D_c^\pi(z_c, \mu) \times d\hat{\sigma}_{ab}^c(x_a P_A, x_b P_B, P_\pi/z_c, \mu) , \quad (10)$$

where the sum is over all contributing partonic channels $a + b \rightarrow c + X$, with $d\hat{\sigma}_{ab}^c$ the associated partonic cross section. The $f_{a,b}$ describe the distributions of partons in the nucleon *. Any factorization of a physical quantity into contributions associated with different length scales will rely on a “factorization” scale that defines the boundary between what is referred to as “short-distance” and “long-distance”. In the present case this scale is represented by μ in Eq. (10). μ is essentially arbitrary, so the dependence of the calculated cross section on μ represents an uncertainty in the theoretical predictions. However, the actual dependence on the value of μ decreases order by order in perturbation theory. This is a reason why knowledge of higher orders in the perturbative expansion of the partonic cross sections is important. We also note that Eq. (10) is of course not an exact statement. There are corrections to Eq. (10) that are down by inverse powers of the momentum transfer, the so-called “power corrections”. These corrections may become relevant towards lower p_T . As we shall see in Figs. 3 and 5 below, comparisons of RHIC data for unpolarized cross sections with theoretical calculations based on Eq. (10) do not suggest that power corrections play a very significant role in the RHIC kinematic regime, even down to fairly low p_T .

Figure 2 offers a graphic illustration of QCD factorization. Thanks to factorization, one can study nucleon structure, represented by the parton densities $f_{a,b}(x, \mu)$, through a measurement of $d\sigma$, hand in hand with a theoretical calculation of $d\hat{\sigma}$. The partonic cross sections may be

*In this particular example, the fact that we are observing a specific hadron in the reaction requires the introduction of additional long-distance functions, the parton-to-pion fragmentation functions D_c^π . These functions have been determined with some accuracy by observing leading pions in e^+e^- collisions and in DIS [48].

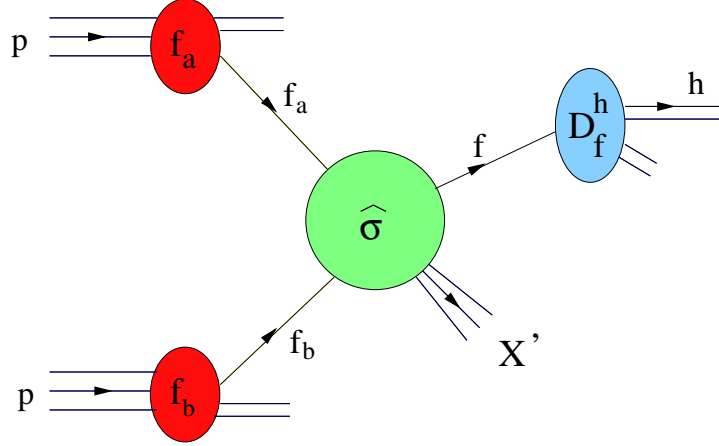


Figure 2: Factorization in terms of parton densities, partonic hard-scattering cross sections, and fragmentation functions.

evaluated in perturbation theory. Schematically, they can be expanded as

$$d\hat{\sigma}_{ab}^c = d\hat{\sigma}_{ab}^{c,(0)} + \frac{\alpha_s}{\pi} d\hat{\sigma}_{ab}^{c,(1)} + \dots \quad (11)$$

$d\hat{\sigma}_{ab}^{c,(0)}$ is the leading-order (LO) approximation to the partonic cross section. The lowest order can generally only serve to give a rough description of the reaction under study. It merely captures the main features, but does not usually provide a quantitative understanding. The first-order (“next-to-leading order” (NLO)) corrections are generally indispensable in order to arrive at a firmer theoretical prediction for hadronic cross sections.

There have already been results from RHIC that demonstrate that the NLO framework outlined above is very successful. Figure 3 shows comparisons of data from PHENIX [49] and STAR [50] for inclusive-pion production $pp \rightarrow \pi^0 X$ with NLO calculations [51, 52, 53]. As can be seen, the agreement is excellent at central and forward rapidities, and down even to p_T values as low as $p_T \gtrsim 1$ GeV. In Fig. 4 we decompose the π^0 cross sections of Fig. 3 into the contributions from the various two-parton initial states [54]. It is evident that in both cases, central and forward, processes with initial gluons dominate by far for the pion transverse momenta accessed so far. This implies that $pp \rightarrow \pi^0 X$ provides an excellent probe of gluons in the nucleon.

A similar comparison is shown for prompt-photon production $pp \rightarrow \gamma X$ in Fig. 5. The left part presents the recent result of a measurement by PHENIX [55], along with the NLO calculation [56, 57]. Again, very good agreement is found. On the right, we show the decomposition of the NLO prompt-photon cross section into the contributions from the initial partonic states. The quark-gluon “Compton process” dominates at a level of 75%.

We note that an agreement between data and NLO calculations like the one seen in Figs. 3 and 5 was not found in previous comparisons made in the fixed-target regime [58]. The good agreement of the pion and photon spectra with NLO QCD at RHIC’s \sqrt{s} , and the good precision of the RHIC data provide a solid basis to extend this type of analysis to polarized reactions. The clear sensitivity to gluons in the initial state further makes the reactions very promising probes of gluon polarization.

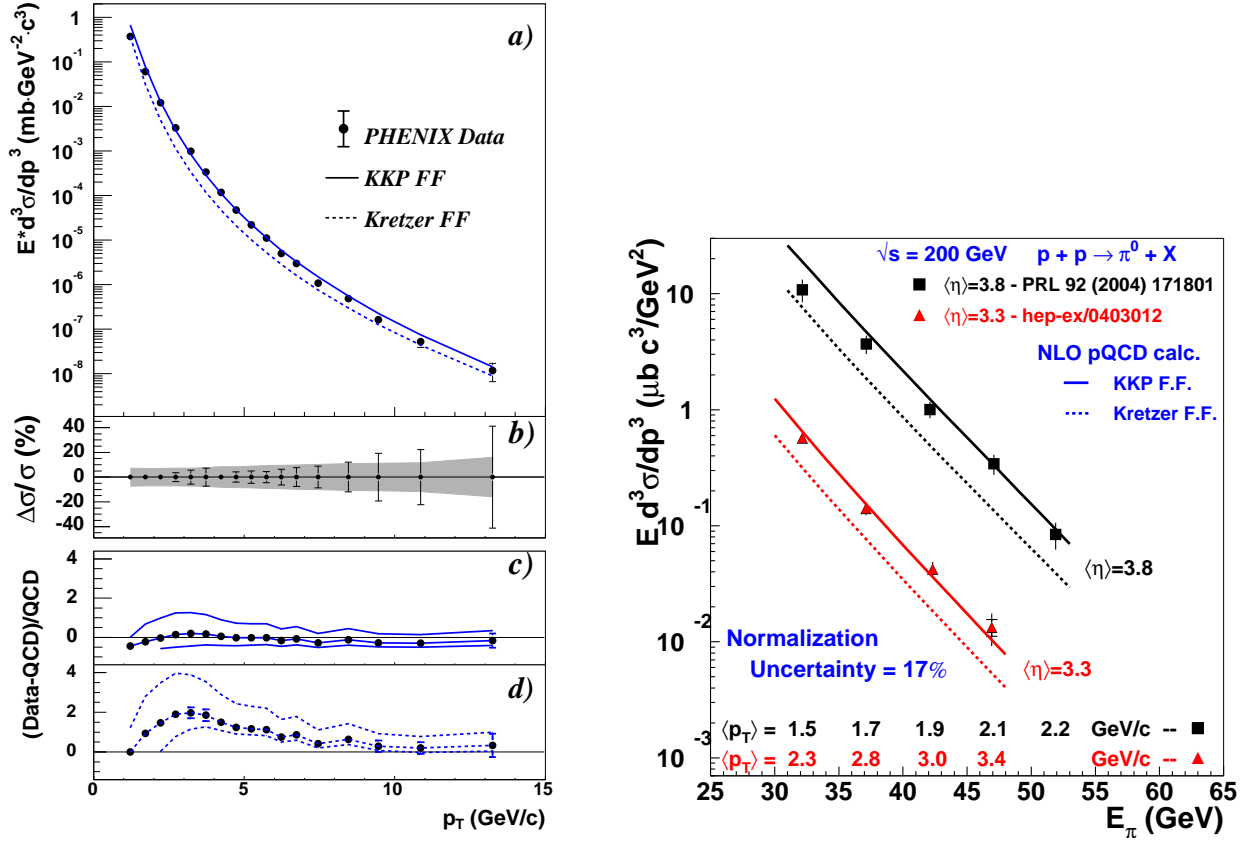


Figure 3: Data from PHENIX (left, [49]) and STAR (right, [50]) for the cross section for inclusive π^0 production $pp \rightarrow \pi^0 X$ at $\sqrt{s} = 200$ GeV. The lines show the results of the next-to-leading order calculation [51, 52, 53] for two different sets of pion fragmentation functions [48].

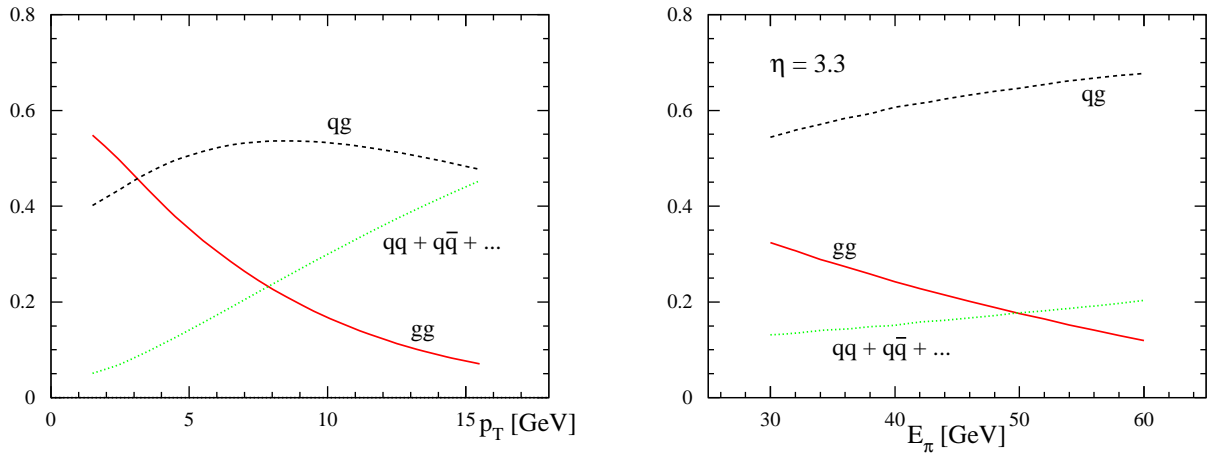


Figure 4: Decomposition of the NLO cross sections for $pp \rightarrow \pi^0 X$ collisions shown in Fig. 3 into the contributions from initial gg , qq , and qg states [54].

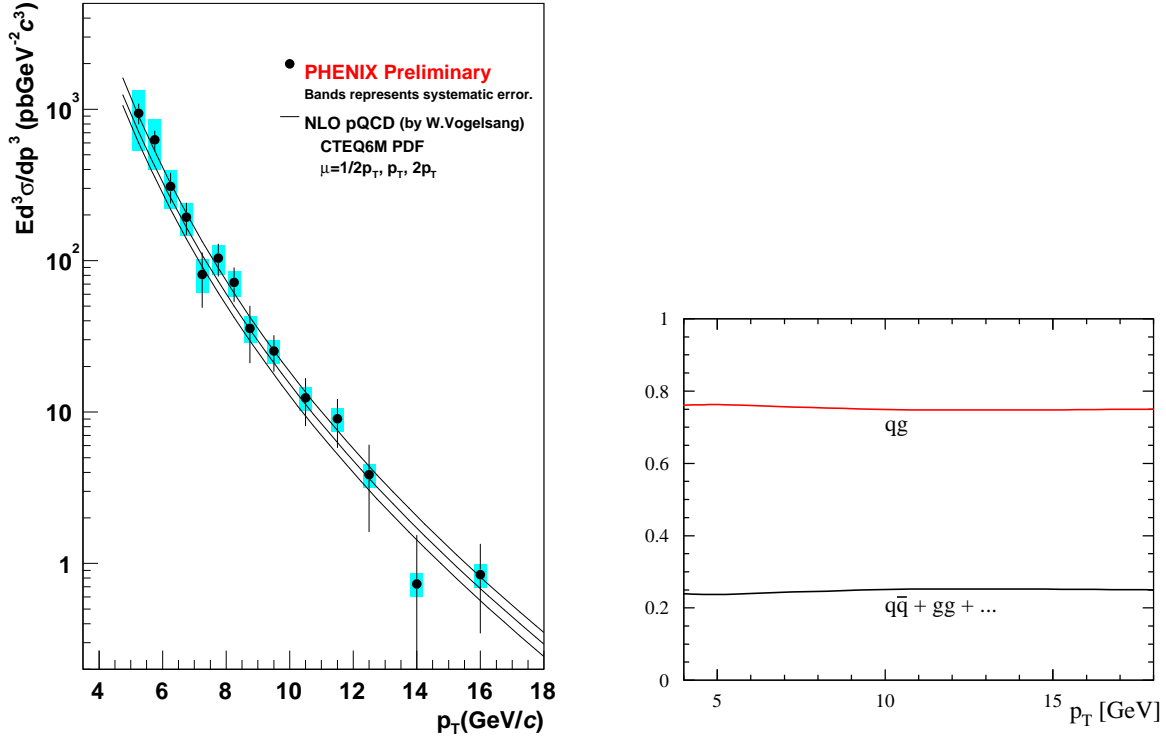


Figure 5: Left: PHENIX data [55] for inclusive prompt-photon production, $pp \rightarrow \gamma X$, compared to the NLO calculation [56, 57]. Right: Decomposition of the NLO cross section into the contributions from initial qg and $q\bar{q}$ +other states.

We have so far only discussed single-inclusive reactions. There are also data from unpolarized proton-proton collisions at RHIC that mark the beginning of promising studies of two-particle correlations in the final state. Figure 6 shows results from STAR for azimuthal correlations of two produced charged hadrons in coincidence [59]. Correlations of the type shown in Fig. 6 are interesting as they are sensitive probes of QCD dynamics. For instance, at lowest order and for initial partons collinear with their parent hadrons, the distribution in Fig. 6 would only have support at $\Delta\phi = \pi$, corresponding to exact back-to-back kinematics. A broad distribution around $\Delta\phi = \pi$ as seen in the figure may result in QCD from gluon radiation, and from “intrinsic” parton transverse momenta. The theoretical analysis of these effects involves *all-order* summations of certain perturbative corrections and is relatively well understood. As we will see in Sec. 2.7, for transversely polarized initial protons there may be interesting spin effects associated with distributions such as the one shown in Fig. 6, that possibly probe parton orbital angular momenta in the proton. We note that also rapidity correlations between two particles in the final state are investigated at RHIC. With improving detector capabilities, these, too, will play an important role in spin physics at RHIC since they allow to pin down the subprocess kinematics to a good degree and hence may contribute to precise mappings of the spin-dependent parton densities of the proton.

We will now turn to polarized pp collisions at RHIC.

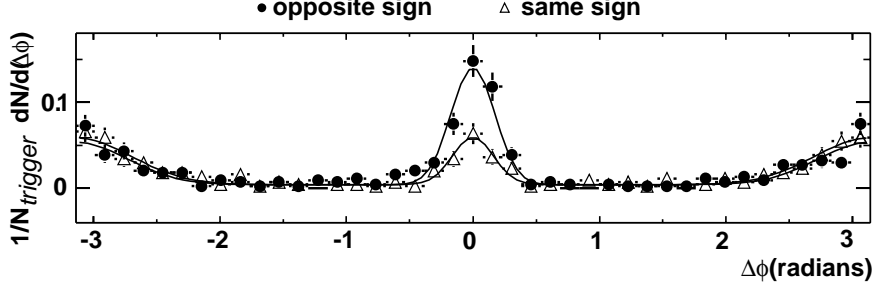


Figure 6: Azimuthal correlations of two charged hadrons produced in coincidence, as seen by STAR [59]. The correlation function is normalized by the number of trigger charged particles having $4 < p_T^{trig} < 6 \text{ GeV/c}$ in the pseudorapidity interval $|\eta| < 0.7$. The associated charged particle has $2 \text{ GeV/c} < p_T < p_T^{trig}$ in the same η interval. Same-sign and opposite-sign charged particles are discriminated. The data are fit by the sum of two Gaussian distributions to represent near-side ($|\Delta\phi| \sim 0$) correlations from di-hadron fragments of the same jet, and opposite-side ($|\Delta\phi| \sim \pi$) hadron fragments from nearly back-to-back jet pairs.

2.4 Probing the spin structure of the nucleon in polarized pp collisions

The measured quantities in spin physics experiments at RHIC are *spin asymmetries*. As an example, for collisions of longitudinally polarized proton beams, one defines a double-spin asymmetry for a given process by

$$A_{LL} = \frac{d\sigma(++) - d\sigma(+-)}{d\sigma(++) + d\sigma(+-)} \equiv \frac{d\Delta\sigma}{d\sigma}, \quad (12)$$

where the signs indicate the helicities of the incident protons. The basic concepts laid out so far for unpolarized inelastic pp scattering carry over to the case of polarized collisions: spin-dependent inelastic pp cross sections factorize into “products” of polarized parton distribution functions of the proton and hard-scattering cross sections describing spin-dependent interactions of partons. As in the unpolarized case, the latter are calculable in QCD perturbation theory since they are characterized by large momentum transfer. Schematically, one has for the numerator of the spin asymmetry:

$$d\Delta\sigma = \sum_{a,b=q,\bar{q},g} \Delta f_a \otimes \Delta f_b \otimes d\Delta\hat{\sigma}_{ab}, \quad (13)$$

where \otimes denotes a convolution and where the sum is over all contributing partonic channels $a + b \rightarrow c + X$ producing the desired high- p_T or large-invariant mass final state. $d\Delta\hat{\sigma}_{ab}$ is the associated perturbative spin-dependent partonic cross section, defined as

$$d\Delta\hat{\sigma}_{ab} = \frac{1}{2} [d\hat{\sigma}_{ab}(++) - d\hat{\sigma}_{ab}(+-)] , \quad (14)$$

the signs denoting the helicities of the initial partons a, b . The sensitivity with which one can probe the polarized parton densities will foremost depend on the weights with which they enter the cross section. Good measures for this are the so-called partonic “analyzing powers”. The latter are just the spin asymmetries

$$\hat{a}_{LL} = \frac{d\hat{\sigma}_{ab}(++) - d\hat{\sigma}_{ab}(+-)}{d\hat{\sigma}_{ab}(++) + d\hat{\sigma}_{ab}(+-)} \quad (15)$$

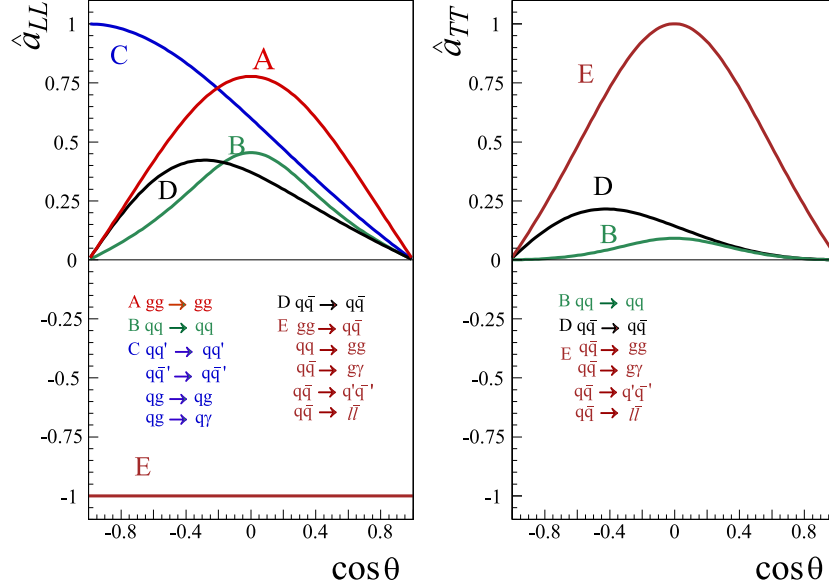


Figure 7: Spin asymmetries for the most important partonic reactions at RHIC at lowest order in QCD. Left: helicity dependence, right: transverse polarization.

for the individual partonic subprocesses. Figure 7 shows these analyzing powers at LO for all partonic reactions. One can see that they are usually very substantial. For future reference, we also give the subprocess asymmetries for *transverse* polarization. Here, Eq. (13) applies as well. The parton densities are then the transversity distributions, to be discussed in more detail below, and the partonic cross sections are defined as in Eq.(14), but for transverse initial polarization. One customarily uses a small δ to designate transversely polarized quantities. In this section, we will focus on the longitudinal case; we will return to transverse polarization in Sec. 2.7.

Since the partonic cross sections are calculable from first principles in QCD, Eq. (13) may be used to determine the polarized parton distribution functions from measurements of the spin-dependent pp cross section on the left-hand side. The crucial point here is, as discussed in the previous section, that the parton distributions are universal. They are the same in all inelastic processes, not only in pp scattering, but also for example in deeply-inelastic lepton nucleon scattering which up to now has mostly been used to learn about nucleon spin structure. This means that inelastic processes with polarization have the very attractive feature that they probe fundamental and universal spin structure of the nucleon. In effect, we are using the asymptotically free regime of QCD to probe the deep structure of the nucleon.

At RHIC, there are a number of sensitive and measurable processes at our disposal. The key ones, some of which will be discussed in detail in the following, are listed in Table 1, where we also give the dominant underlying partonic reactions and the aspect of nucleon spin structure they probe. We emphasize that, even though we have only shown LO results in Fig. 7, the NLO corrections are available for each process relevant for RHIC-Spin, thanks to considerable efforts made over the past decade or so. We give reference to the corresponding work in the first column of Table 1. These calculations bring the theoretical calculations for RHIC-Spin to the same level that has been so successful in the unpolarized case, as demonstrated by Figs. 3 and 5. For each of the processes in Table 1 the parton densities enter with different weights, so that each has its own role in helping to determine the polarized parton distributions. Some will allow a

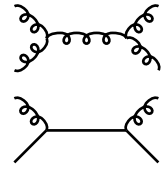
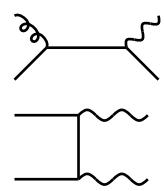
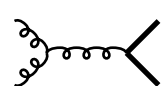
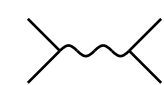
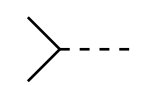
Reaction	Dom. partonic process	probes	LO Feynman diagram
$\vec{p}\vec{p} \rightarrow \pi + X$ [51, 52]	$\vec{g}\vec{g} \rightarrow gg$ $\vec{q}\vec{g} \rightarrow qg$	Δg	
$\vec{p}\vec{p} \rightarrow \text{jet(s)} + X$ [60, 61]	$\vec{g}\vec{g} \rightarrow gg$ $\vec{q}\vec{g} \rightarrow qg$	Δg	(as above)
$\vec{p}\vec{p} \rightarrow \gamma + X$ $\vec{p}\vec{p} \rightarrow \gamma + \text{jet} + X$ $\vec{p}\vec{p} \rightarrow \gamma\gamma + X$ [56, 62, 63, 64, 65]	$\vec{q}\vec{g} \rightarrow \gamma q$ $\vec{q}\vec{g} \rightarrow \gamma q$ $\vec{q}\vec{q} \rightarrow \gamma\gamma$	Δg Δg $\Delta q, \Delta \bar{q}$	
$\vec{p}\vec{p} \rightarrow DX, BX$ [67]	$\vec{g}\vec{g} \rightarrow c\bar{c}, b\bar{b}$	Δg	
$\vec{p}\vec{p} \rightarrow \mu^+\mu^- X$ (Drell-Yan) [68, 69, 70]	$\vec{q}\vec{q} \rightarrow \gamma^* \rightarrow \mu^+\mu^-$	$\Delta q, \Delta \bar{q}$	
$\vec{p}\vec{p} \rightarrow (Z^0, W^\pm)X$ $p\vec{p} \rightarrow (Z^0, W^\pm)X$ [68]	$\vec{q}\vec{q} \rightarrow Z^0$ $\vec{q}'\vec{q} \rightarrow W^\pm, \vec{q}\vec{q}' \rightarrow W^\mp$	$\Delta q, \Delta \bar{q}$	

Table 1: Key processes at RHIC for the determination of the parton distributions of the longitudinally polarized proton, along with the dominant contributing subprocesses, the parton distribution predominantly probed, and representative leading-order Feynman diagrams. The references given in the left column are for the corresponding next-to-leading order calculation.

clean determination of gluon polarizations, others are more sensitive to quarks and antiquarks. Eventually, when data from RHIC will become available for most or all processes, a “global” analysis of the data, along with information from lepton scattering, will be performed which then determines the $\Delta q, \Delta \bar{q}, \Delta g$. For further details, see Sec. 2.5.

We will now address some of the most important processes in more detail, summarizing theoretical predictions and experimental plans and prospects at RHIC. We will start with those that are sensitive to gluon polarization in the proton, and then discuss W production which will give information about the quark and antiquark polarizations.

2.4.1 Exploring the gluon contribution to the proton spin

To learn about the contribution of gluons to the proton spin is the most compelling motivation for doing experiments with polarized protons at RHIC. The importance of measuring the polarized gluon distribution $\Delta g(x, Q^2)$ has been universally recognized ever since the “spin crisis” was

discovered. In fact, besides RHIC, there are – and have been in the past – several other efforts in the world to access Δg . In the early 1990s, the E704 experiment at Fermilab measured [71] the double-spin asymmetry $A_{LL}^{\pi^0}$ in $pp \rightarrow \pi^0 X$ with a polarized proton beam and polarized target at $\sqrt{s} \approx 20$ GeV, accessing pion transverse momenta of $1 \lesssim p_T \lesssim 4$ GeV. As we described in Sec. 2.1, fixed-target inclusive DIS, so far the main tool in QCD spin physics, is not well suited for measuring the gluon density of the probed nucleon, since the photon primarily “sees” the quarks and antiquarks. Nevertheless, certain more exclusive final states such as pairs of heavy flavors or high-transverse momentum hadrons select the photon-gluon fusion process $\gamma^* g \rightarrow q\bar{q}$, resulting in sensitivity to Δg . This is the strategy adopted by the lepton-nucleon scattering experiments SMC [72], HERMES [73] and COMPASS [74]. The ongoing experiments are briefly described in Sec. ??, and their sensitivities to Δg are included in Fig. 16 below.

RHIC unites several features that make it unique to explore the gluon polarization. Its first asset is the high energy, where the all-important theoretical concept of factorization (see Sec. 2.3) is expected to work best. Indeed, several unpolarized pp cross sections for reactions sensitive to gluons have already been measured at RHIC and are described well by perturbative QCD predictions, see Figs. 3 and 5. At collider energies, a range of kinematics opens up, allowing transverse momenta of an observed final state well into the region described by perturbative QCD, for mid and forward rapidities. The span in p_T allows probes of $\Delta g(x)$ over a wide range of x , thus helping to significantly constrain its x -integral. In addition, quark polarization measurements at RHIC will be compared to those from DIS.

The next crucial feature of RHIC is that two different energies, $\sqrt{s} = 200$ and 500 GeV, will be available. It is conceivable that gluons are still rather strongly polarized towards low momentum fractions x , so that a significant contribution to the integral of Δg could come from that region. This becomes evident from the left part of Fig. 8, which shows again results for $\Delta g(x, Q^2 = 5 \text{ GeV}^2)$ from several recent analyses of scaling violations in polarized DIS, along with a typical uncertainty to be assigned to it at present. We plot $x\Delta g(x, Q^2)$ as a function of $\log(x)$, so that any part of the area underneath the curve directly gives the contribution to the integral $\int_0^1 \Delta g(x, Q^2) dx$. The current uncertainties do not rule out sizable contributions from $x \lesssim 0.01$. It is therefore important to have information on Δg at as small momentum fractions as possible. Roughly, at mid rapidity at RHIC, the lowest momentum fraction probed by a high- p_T final state is $x \sim p_T/\sqrt{s}$. Since p_T needs to be high enough for the process to remain amenable to QCD perturbation theory, a key to access smaller x is to *increase* energy. This makes collisions at 500 GeV indispensable. At the same time, measurements at $\sqrt{s} = 200$ GeV better probe larger x , which may also contribute significantly to the integral of Δg . In addition, there is a large overlap between the x regions covered at $\sqrt{s} = 200$ and 500 GeV, respectively. The importance of this is easily overlooked, but hard to overrate. Perturbative QCD makes definite predictions for the energy dependence of the cross section, so the consistency of the $\sqrt{s} = 200$ vs 500 GeV results with the predicted changes will provide an important test of the theoretical framework. QCD spin interactions at such high energies are uncharted territory, becoming only now accessible at RHIC, and their test is of fundamental importance.

What we have described so far would already give RHIC good possibilities to access Δg even if there were only one physics channel to study. The experiments at RHIC will measure a variety of channels, each of them highly sensitive to gluon polarization, and each with its individual strengths. We have listed the key processes in Table 1. PHENIX and STAR differ in their capabilities to detect the various channels. With major upgrades in preparation, coverage

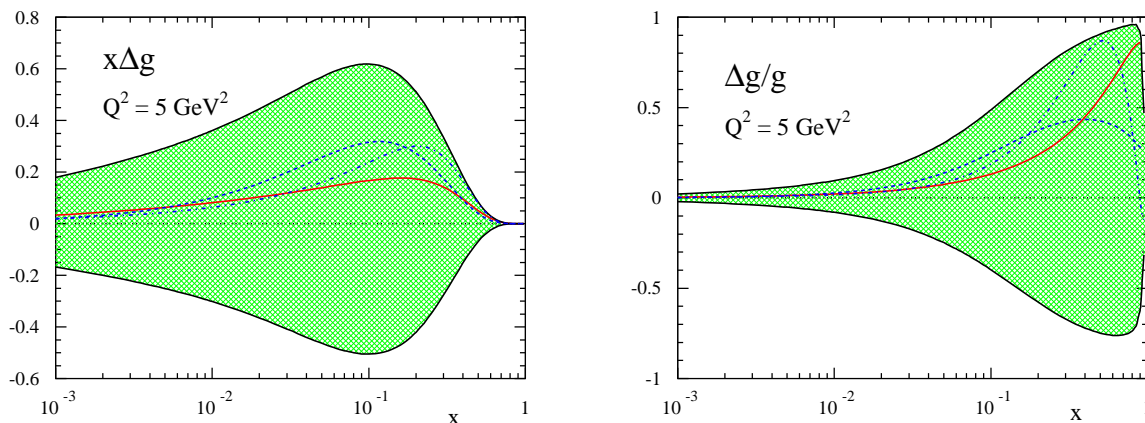


Figure 8: Left: results for $\Delta g(x, Q^2 = 5 \text{ GeV}^2)$ from recent analyses [26, ?, ?] of polarized DIS (as already shown in Fig. 1), along with a typical current uncertainty in Δg . Note that we show $x\Delta g$ as a function of $\log(x)$, in order to display the contributions from various x -regions to the integral of Δg . Right: same as left, but now showing the “net gluon polarization” $\Delta g(x, Q^2)/g(x, Q^2)$ at $Q^2 = 5 \text{ GeV}^2$, where g is the unpolarized gluon distribution of [?].

will be increased with time, allowing ever more detailed studies of the key physics processes, for instance also in terms of coincidences of two particles in the final state. Again, this will be instrumental in precisely mapping the x -dependence of Δg . The main power of RHIC in telling us about $\Delta g(x, Q^2)$ lies in the *combination* of all these planned measurements, which will determine $\Delta g(x, Q^2)$ over a wide range of x and Q^2 as well as test the robustness of the overall approach. We will now address each of the key processes individually. In what follows, we will use the “uncertainty band” from DIS shown in Fig. 8 (and for the quark distributions in Fig. 1) as a guide. Within the framework of the NLO calculation available for each of the key processes, we will translate these bands into a band for the double-spin asymmetries A_{LL} for the various key processes. This band, which we will refer to as “NLO uncertainty band” indicates the uncertainty with which A_{LL} can currently be predicted. We will then confront the result with the expected sensitivities at RHIC and will find that they are much better than the current uncertainties, even for each channel individually. Note that the right part of Fig. 8 displays the typical current uncertainty on $\Delta g/g$. The measured spin asymmetries A_{LL} are effectively proportional to $\Delta g/g$ in linear or quadratic form. The “target” luminosities used for experimental sensitivities aim to improve the statistical precision for Δg over current DIS analyses by at least a factor ≈ 3 , even for the pp channel with the smallest cross section (but cleanest) used to probe Δg .

The abundant probes: high- p_T pions and jets. To match the expected improvements in machine and detector capabilities, STAR and PHENIX will address the gluon polarization with a progression of probes. At the moment, as RHIC is still developing higher luminosity and polarization, measurements exploit the abundant channels for inclusive pion and jet production. Indeed, PHENIX has already published [?] first A_{LL} data for $pp \rightarrow \pi^0$ from the 2003 RHIC run, shown in Fig. 9. Even with an integrated luminosity of only a few hundred nb^{-1} and a beam polarization of $\approx 30\%$ the data are already at the verge of constraining Δg at a level comparable to the information extracted from the polarized-DIS database. This is only the beginning, of course. The improved luminosity and beam polarization anticipated for the 2005 run should

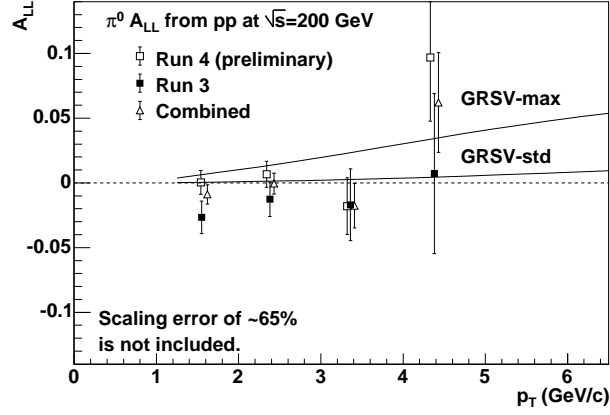


Figure 9: PHENIX measurements [66] for the double-spin asymmetry $A_{LL}^{\pi^0}$ at central rapidities. The curves are NLO calculations based on the parton densities of [26]. “GRSV-std” refers to the central fit to polarized DIS data, “GRSV-max” to the upper end of the bands shown in Fig. 8.

already provide an order of magnitude decrease in statistical uncertainties. With time, this channel will become a high-precision probe. The left part of Fig. 10 shows the “NLO uncertainty band” for $A_{LL}^{\pi^0}$ at $\sqrt{s} = 200$ GeV, along with expected sensitivities for the PHENIX experiment to be achieved for integrated luminosity of 100/pb and polarization 70%. We emphasize that these sensitivities are based on measurements already made, and hence incorporate realistic π^0 triggering and reconstruction efficiencies (40% at high p_T for the PHENIX π^0 measurement). On the right, we show the pion spin asymmetry at very forward rapidities, where STAR is making measurements now, and will have dramatically increased capabilities after planned upgrades that include a forward meson spectrometer. This very forward region is important since it gives information at low gluon momentum fractions $x \sim p_T e^{-\eta}/\sqrt{s}$, as will be discussed in more detail below. Note that we have chosen the kinematic regimes in the two plots in Fig. 10 corresponding to Figs. 3 and 4, that is, where measurements have already been made in the unpolarized case and comparisons to NLO theory have been successful.

Figure 11 shows projections for the spin asymmetry for inclusive-jet production $pp \rightarrow \text{jet} + X$ at STAR, in the rapidity region $-1 \leq \eta^{\text{jet}} \leq 2$. On the left we show calculations for $\sqrt{s} = 200$ GeV. We also show projected uncertainties for the 2005 run. Again, these are based on measurements already made, and hence incorporate realistic jet triggering and reconstruction efficiencies ($\sim xx\%$ and $yy\%$, respectively). Evidently, the data anticipated from the 2005 run alone will significantly reduce the current $\Delta g(x)$ uncertainties. On the right of Fig. 11 we show the prospects for the longer-term future, when collisions at $\sqrt{s} = 500$ GeV will be available. With luminosity 300/pb and polarization of 70%, very precise measurements should emerge. Experimental efficiencies and bandwidth limits are assumed. A systematic uncertainty on the raw asymmetry measurement is assumed at 10^{-3} .

Pion and jet production are very powerful probes of gluon polarization at the single-inclusive level ($pp \rightarrow \pi^0 X$, or $pp \rightarrow \text{jet} X$, respectively). Further important measurements will be A_{LL} for the production of *pairs* of hadrons or jets detected in coincidence, as a function both of transverse momenta $p_{T,1,2}$ and pseudorapidities $\eta_{1,2}$ of the detected particles. For example, in the case of dijet production, if LO $2 \rightarrow 2$ processes and kinematics dominate, one can unambiguously reconstruct the partonic momentum fractions in the incident protons as $x_1 = p_T(e^{-\eta_1} + e^{-\eta_2})/\sqrt{s}$, $x_2 = p_T(e^{\eta_1} + e^{\eta_2})/\sqrt{s}$, where p_T is the transverse momentum of each of the jets. Likewise,

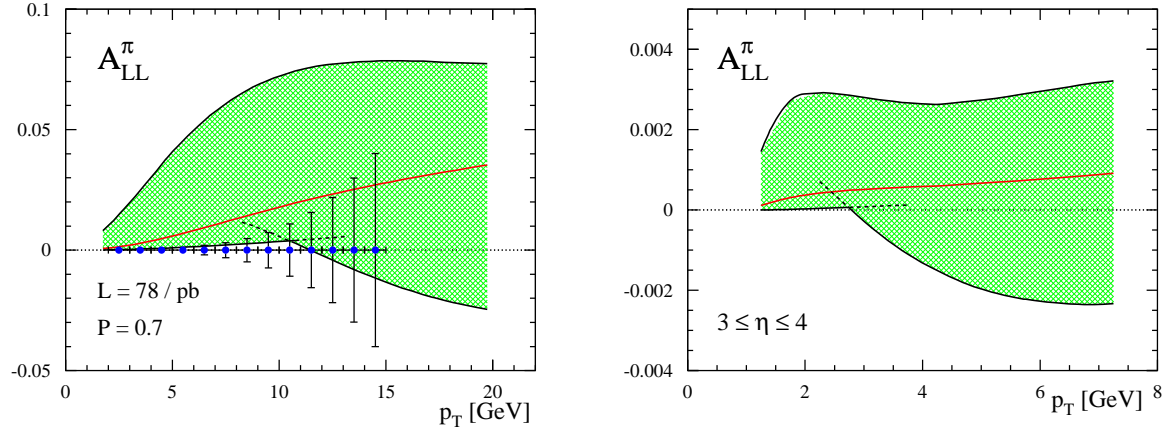


Figure 10: Left: current uncertainty in $A_{LL}^{\pi^0}$ due to Δg (see Fig. 8), and projected sensitivities for measurements by PHENIX at mid rapidities and $\sqrt{s} = 200$ GeV for integrated luminosity of 100/pb and polarization 70%. Note the “cusp” in the uncertainty band near $p_T = 10$ GeV which results from use of a gluon distribution with strong *negative* polarization. The cusp occurs when the process $qg \rightarrow qg$ (which contributes negatively to the spin asymmetry for $\Delta g < 0$) starts to dominate over $gg \rightarrow gg$ (which is always positive). This is emphasized by the dashed lines. Right: same as left, but for very forward pion pseudorapidities $3 \leq \eta \leq 4$. INSERT AND DESCRIBE ERROR BARS ON THE LEFT.

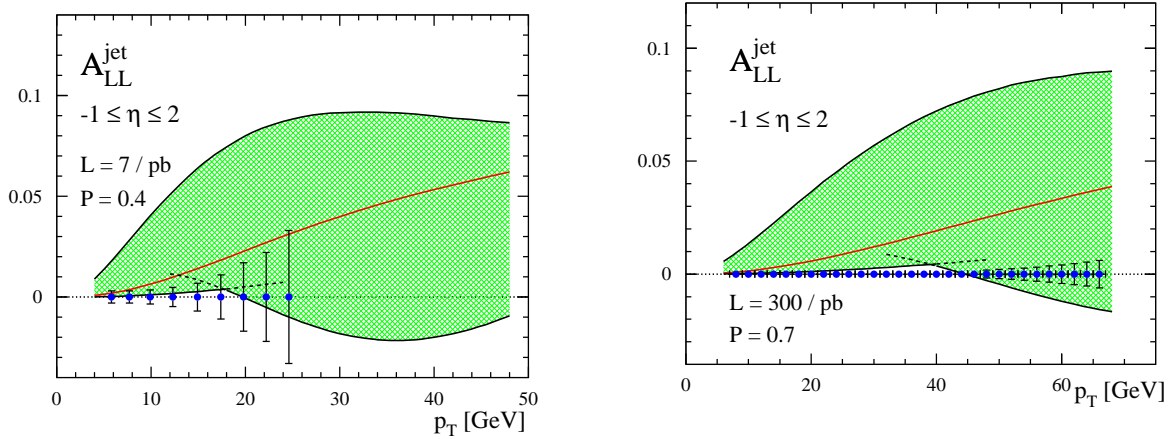


Figure 11: Predicted spin asymmetries for inclusive-jet production at $\sqrt{s} = 200$ and 500 GeV, for rapidities $-1 \leq \eta \leq 2$. As before, the bands illustrate the current uncertainty due to Δg . The errors are projections as described in the text. ERRORS TO BE CHANGED !

strong correlations between measured kinematic variables and partonic momentum fractions are found when the jets are replaced by π^0 's as surrogates. We give an example from a Monte-Carlo simulation in Fig. 12. A neutral “trigger” pion is detected at forward rapidities, $3 \leq \eta_{\pi,1} \leq 4$ and with $p_{T,\pi,1} \geq 2.5$ GeV. Because $\eta_{\pi,1}$ is so large, partonic collisions tend to become

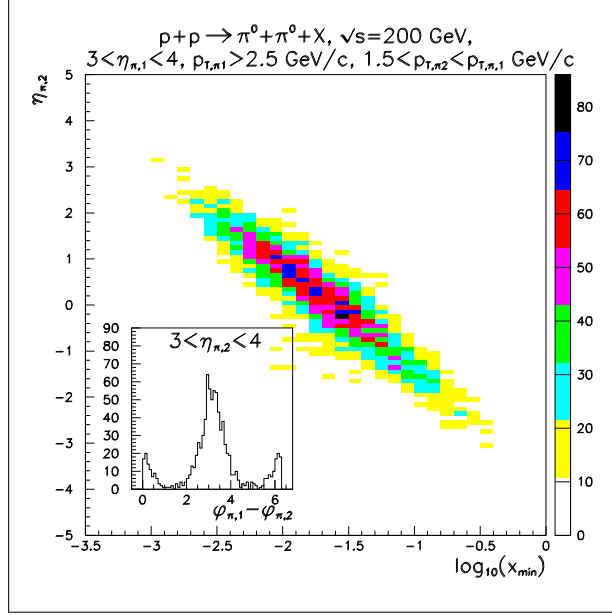


Figure 12: PYTHIA study of $\pi^0\pi^0$ pair production in pp collisions at $\sqrt{s} = 200$ GeV. A trigger pion is detected at forward rapidities. The plot shows correlations between the rapidity of an associated pion and the softer momentum fraction probed in the proton (see text). For the conditions used for this figure, the expected uncertainty for A_{LL} measurements will be at the half-per cent level over the whole range of $\eta_{\pi,2}$.

very asymmetric, with the momentum fraction x_1 associated with the proton moving in forward direction large, and x_2 small. More precisely, selecting a second pion within $1.5 \leq p_{T,\pi 2} \leq p_{T,\pi 1}$, one can now virtually “dial” x_2 , as shown by the correlation in Fig. 12. In particular, choosing $\eta_{\pi,1}$ fairly large as well, the smallest x_2 become accessible. Selective information obtained in this way may be more powerful in providing information on $\Delta g(x, Q^2)$ at lower x than single-inclusive production, and hence be vital in constraining the integral of Δg . As an additional benefit, in this kinematic regime $A_{LL}^{\pi^0}$ is mainly driven by quark-gluon scattering (see the right part of Fig. 4) and hence is mostly sensitive to $\Delta q(x_1)\Delta g(x_2)$. At high x_1 , quark polarization is known to be large from the polarized-DIS measurements. A region of phase space with sizeable p_T and large $\Delta\eta$, where qq scattering is important, can also be exploited to provide quark polarization results that can be compared to DIS.

Direct-photon production. Another important process at RHIC for measuring $\Delta g(x, Q^2)$ is direct-photon production $pp \rightarrow \gamma X$. For RHIC kinematics, this channel is dominated at the $\sim 75\%$ level by quark-gluon scattering, starting with the LO QCD Compton process $qg \rightarrow q\gamma$ (see the right part of fig. 5). The simple pointlike coupling of the photon in this process is well-understood in Quantum Electrodynamics, and the analyzing power is large (see Fig. 7). Thus direct-photon production provides strong and nearly undiluted sensitivity to $\Delta g(x, Q^2)$. In particular, its spin asymmetry is *linear* in Δg , unlike the case of pion and jet production discussed above where at lower p_T terms quadratic in Δg may dominate. An important role of direct-photon production will therefore be to determine the *sign* of gluon polarization. This becomes evident in Fig. 13 which shows the “NLO uncertainty band” for A_{LL}^γ for inclusive photon production at RHIC, for both $\sqrt{s}=200$ and 500 GeV. The calculations include an isolation cut on the photon [63], as will be imposed experimentally to suppress contributions from jets.

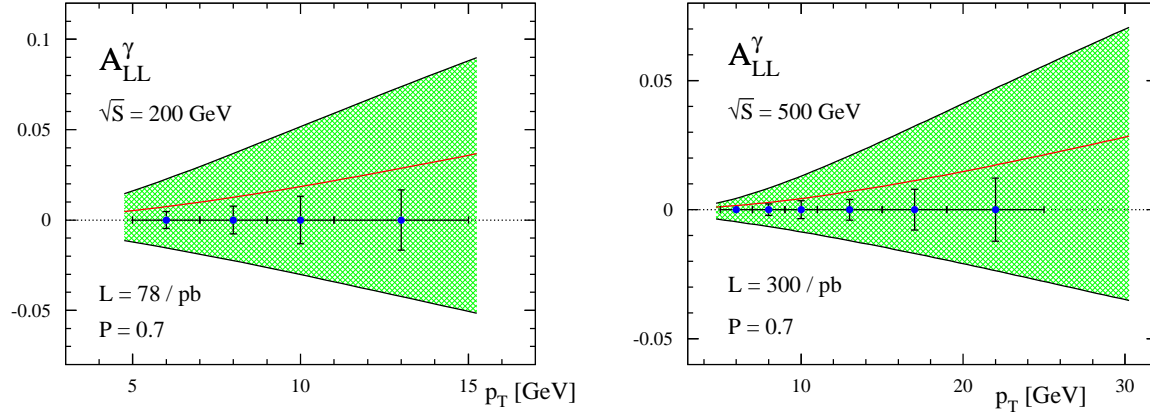


Figure 13: Predicted spin asymmetries for single-inclusive direct-photon production $pp \rightarrow \gamma X$ at $\sqrt{s} = 200$ and 500 GeV, for central rapidities. As before, the bands illustrate the current uncertainty due to Δg . The errors are projections for the PHENIX experiment, as described in the text.

The projected experimental error bars in Fig. 13 represent statistical and background subtraction errors with realistic p_T cuts, achievable at RHIC with beam polarizations of 0.7 and integrated luminosities recorded at PHENIX of 78 pb^{-1} at 200 GeV and 300 pb^{-1} at 500 GeV . Comparison of these error bars with the present theoretical uncertainty band shows that from prompt photon production alone, despite its small cross section, one can substantially improve upon present uncertainties in the gluon polarization. In particular, the sign of $\Delta g(x, Q^2)$ readily translates into the sign of A_{LL}^γ . We emphasize again that there are already measurements of the unpolarized cross section for $pp \rightarrow \gamma X$ at $\sqrt{s} = 200 \text{ GeV}$ from PHENIX which are in good agreement with the NLO theoretical calculations (see Fig. 5).

As for pions and jets, studies of coincidences in the final state, in this case of a photon and a recoiling jet or leading hadron, will prove very useful for providing an experimental map of the so far unconstrained shape of $\Delta g(x, Q^2)$. This again allows event-by-event constraints on the colliding parton kinematics. Measurements of $pp \rightarrow \gamma + \text{jet} + X$ will be a particular emphasis in the STAR experiment. The possibilities are illustrated by STAR simulations in Fig. 14. Here, a LO analysis of the parton kinematics from the detected photon and jet properties has been used to determine the quark and gluon x -values event by event. Conservative cuts requiring $p_T \geq 10 \text{ GeV}$ and the larger of the two momentum fractions to be bigger than 0.2 (the latter to select the quarks with highest polarization) have been imposed on the events included in Fig. 14. The simulations illustrate the statistical uncertainties achievable in $\Delta g(x)$ for three different parameterizations of [27] consistent with the DIS database. The sensitivity to gluon polarization is evident.

Heavy-flavor production. One final set of channels for probing gluon polarization at RHIC that we will discuss here involves the production of hadrons carrying “open” charm or bottom quarks. Planned vertex detectors for both STAR and PHENIX are necessary to select these events based on displaced vertices. The heavy-flavor final states are produced predominantly by gluon-gluon fusion, $g + g \rightarrow Q + \bar{Q}$ (see Table 1), so that the spin asymmetry provides quadratic sensitivity

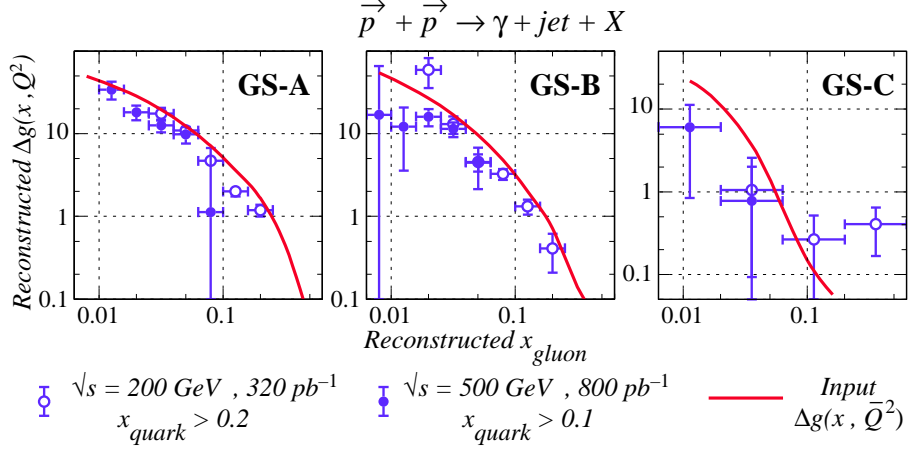


Figure 14: Possibilities for direct reconstruction of $\Delta g(x, Q^2)$ from measurements of the spin asymmetry in $pp \rightarrow \gamma + jet + X$ at $\sqrt{s} = 200$ and 500 GeV (see text). Note that here the statistical uncertainties correspond to integrated luminosities of 320/pb (800/pb) at $\sqrt{s} = 200$ (500) GeV. The curves show three different models for gluon polarization $\Delta g(x, Q^2)$ of [27].

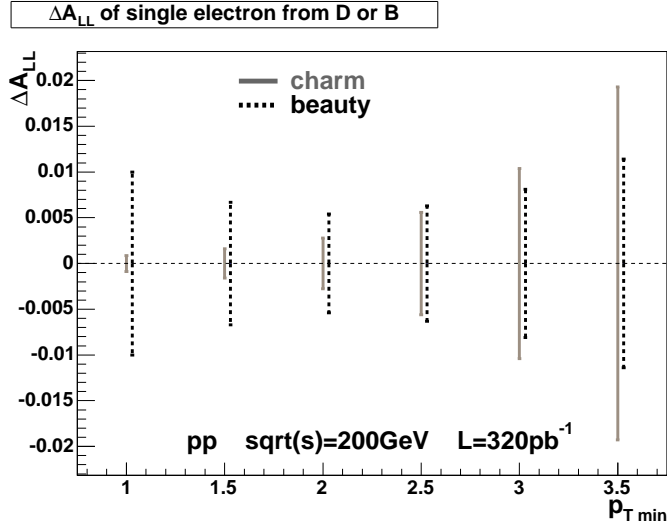


Figure 15: Projected uncertainties in the spin asymmetry for heavy flavor production for PHENIX measurements at $\sqrt{s} = 200$ GeV with a proposed vertex detector.

to Δg . The decay of heavy-flavor mesons dominates the inclusive production of leptons in the $\sim 2 - 10$ GeV/c range, so that the highest statistics measurements of heavy flavor production will be made via inclusive electron or muon spectra. Forward lepton detection would provide access to gluons at low x . Figure 15 shows PHENIX projections of A_{LL} uncertainties attainable via inclusive electron detection at mid-rapidity. RHIC measurements of heavy flavor production, including hidden flavor in J/ψ production, will also help to test the quantitative level of understanding of these channels and the assumption of gluon fusion dominance.

2.5 Global Analysis

Global analysis of hard reaction data with longitudinally polarized targets will provide an optimal framework to combine the various production channels of the previous section into a systematically controlled extraction of $\Delta g(x)$. The technique is to optimize the agreement between measured cross sections σ^{exp} , relative to the data accuracy $\delta\sigma^{\text{exp}}$, and the theoretical σ^{th} by minimizing the χ^2 function

$$\chi^2 = \sum_{\text{measurements}} \left(\frac{\sigma^{\text{th}}(\Delta g, \dots) - \sigma^{\text{exp}}}{\delta\sigma^{\text{exp}}} \right)^2. \quad (16)$$

As demonstrated, RHIC will add to the above sum over “global” measurements new reactions that are mostly sensitive to the gluon polarization $\Delta g(x)$. The advantages of a full fledged global analysis program are manifold:

- The information from the various reaction channels, outlined above, is all combined into a single result for $\Delta g(x)$.
- The global analysis effectively deconvolutes the experimental information, which in its raw form is smeared over the fractional gluon momentum x , and fixes the gluon polarization at definite values of x .
- It provides a framework to determine an error on the gluon polarization.
- Correlations with other experiments, included in the sum (\sum) and sensitive to degrees of freedom different from Δg , are automatically respected.

The above items have been developed very successfully over many years for unpolarized parton densities [?]. The extraction of Δg will benefit from the fact that these techniques can be directly transferred to the analysis of polarized data. To constrain the polarized quark and antiquark distributions, the global analysis will include also the results from polarized DIS, and eventually from W production at RHIC, to which we will turn now.

2.6 W production at RHIC

2.6.1 Introduction

Measurements in polarized DIS [?], when combined with information from baryon octet β -decays [?], show that the total quark-plus-antiquark contribution to the proton’s spin, summed over all flavors, is surprisingly small. In the standard interpretation of the β -decays [?], this finding is equivalent to evidence for a large negative polarization of strange quarks in the proton, which makes it likely that also the $SU(2)$ (u, d) sea is strongly negatively polarized. This view is corroborated by the fact that in this analysis the spin carried, for example, by u quarks comes out much smaller than generally expected in quark models [?], implying that a sizeable negative u -sea polarization partly compensates that of the valence u quarks. Alternative treatments of the information from β -decays [?, ?], when combined with the DIS results, also directly yield large

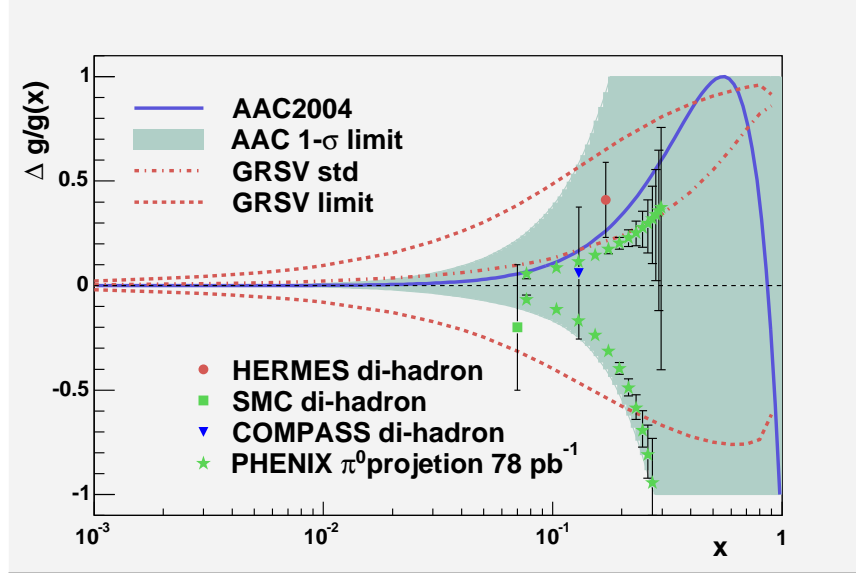


Figure 16: 200 GeV, 78/pb, $P = 0.7$

negative \bar{u} and \bar{d} polarizations. Inclusive DIS (through γ^* exchange) itself is sensitive to the combined contributions of quarks and antiquarks of each flavor but cannot provide information on the polarized quark and antiquark densities separately. Directly measuring the individual polarized antiquark distributions is therefore an exciting task and will also help to clarify the overall picture concerning DIS and the β -decays.

Further motivation for dedicated measurements of antiquark densities comes from unpolarized physics. Experiments in recent years have shown [?, ?, ?] a strong breaking of $SU(2)$ symmetry in the antiquark sea, with the ratio $\bar{d}(x)/\bar{u}(x)$ rising to 1.6 or higher. It is very attractive to learn whether the polarization of \bar{u} and \bar{d} is large and asymmetric as well. Within the chiral quark soliton model based on a $1/N_c$ expansion, it is expected that the polarized flavor asymmetry, $\Delta\bar{u} - \Delta\bar{d}$, is larger than the experimentally established flavor asymmetry in the unpolarized sector [?]. A measurement of the polarized flavor asymmetry will shed light into the underlying mechanism responsible for the expected polarized flavor asymmetry. RHIC experiments will measure the \bar{d}/\bar{u} unpolarized ratio and the \bar{u} and \bar{d} polarizations separately.

Semi-inclusive DIS measurements [?] are one approach to achieving a separation of quark and antiquark densities. This method combines information from proton and neutron (or deuteron) targets and uses correlations in the fragmentation process between the type of leading hadron and the flavor of its parton progenitor, expressed by fragmentation functions. The dependence on the details of the fragmentation process limits the accuracy of this method. At RHIC the polarization of the u , \bar{u} , d , and \bar{d} quarks in the proton will be measured directly and precisely using maximal parity violation for production of W bosons in $u\bar{d} \rightarrow W^+$ and $d\bar{u} \rightarrow W^-$ [?, ?, ?, ?, ?].

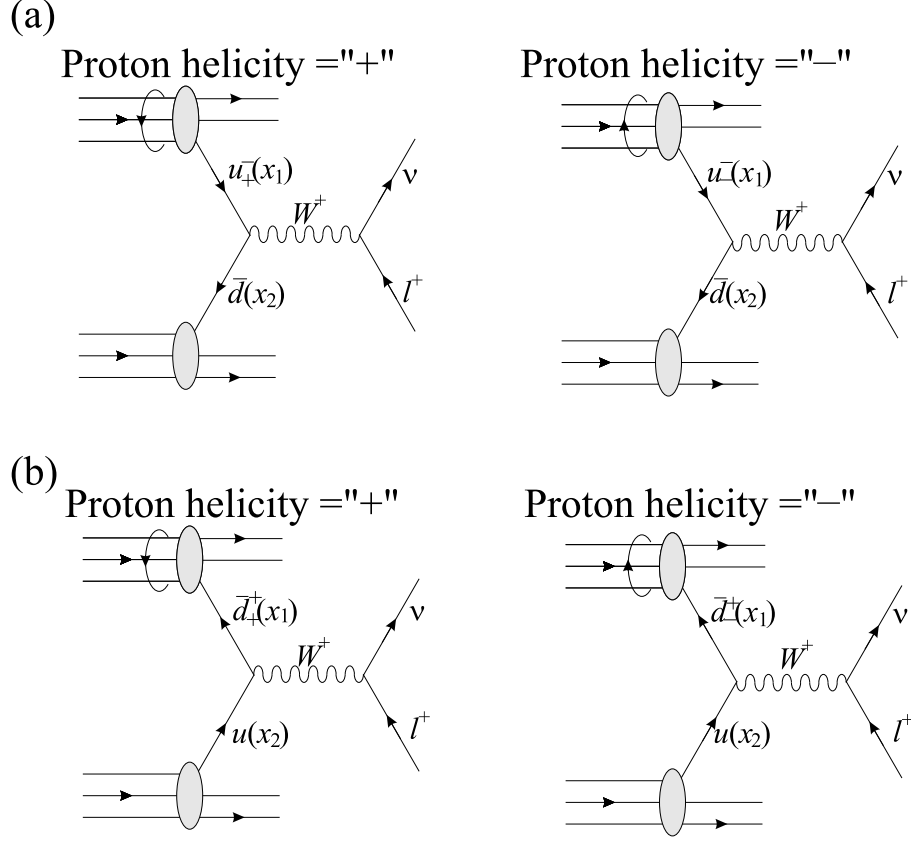


Figure 17: Production of a W^+ in a $\vec{p}p$ collision, at lowest order. (a) Δu is probed in the polarized proton. (b) $\Delta \bar{d}$ is probed.

2.6.2 Weak Boson Production

Within the standard model, W bosons are produced through pure $V-A$ interaction. Thus, the helicity of the participating quark and antiquark are fixed in the reaction. In addition, the W couples to a weak charge that correlates directly to flavors, if we concentrate on one generation. Indeed the production of W s in pp collisions is dominated by u, d, \bar{u} , and \bar{d} , with some contamination from s, c, \bar{s} , and \bar{c} , mostly through quark mixing. Therefore W production is an ideal tool to study the spin-flavor structure of the nucleon.

The leading-order production of W s, $u\bar{d} \rightarrow W^+$, is illustrated in Figure 17. The longitudinally polarized proton at the top of each diagram collides with an unpolarized proton, producing a W^+ . At RHIC the polarized protons will be in bunches, alternately right- (+) and left- (−) handed. The parity-violating asymmetry is the difference of left-handed and right-handed production of W s, divided by the sum and normalized by the beam polarization:

$$A_L^W = \frac{1}{P} \times \frac{N_-(W) - N_+(W)}{N_-(W) + N_+(W)} . \quad (17)$$

We can construct this asymmetry from either polarized beam, and by summing over the helicity states of the other beam. The production of the left-handed weak bosons violates parity maximally. Therefore, if for example the production of the W^+ proceeded only through the diagram in Figure 17a, the parity-violating asymmetry would directly equal the longitudinal polarization

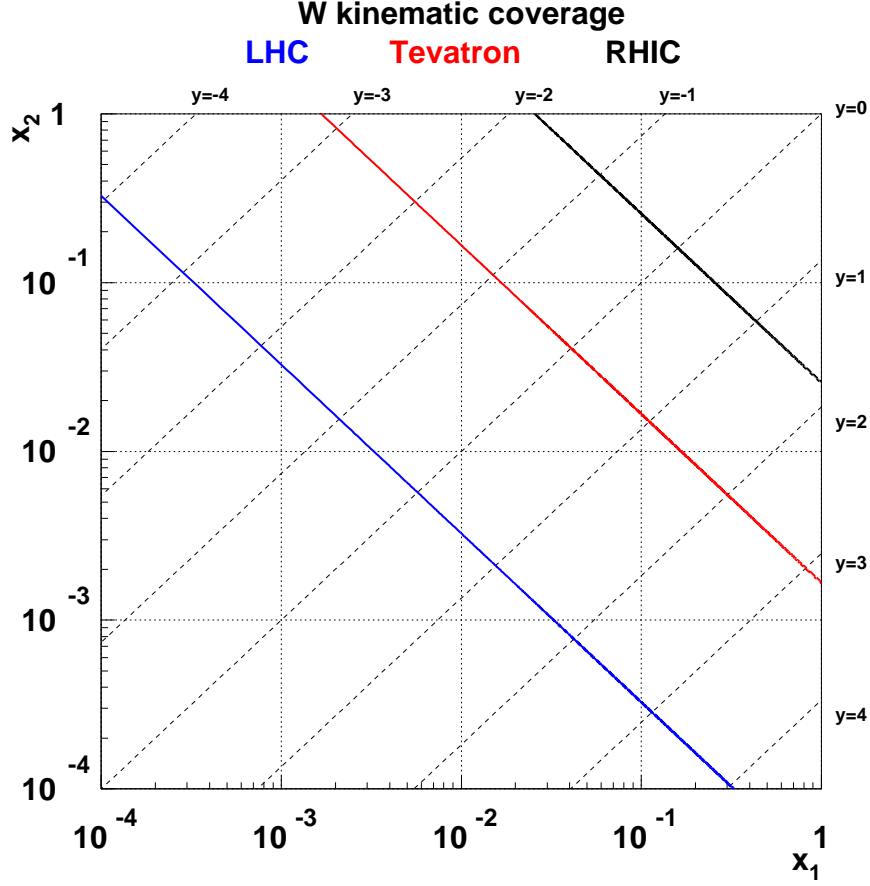


Figure 18:

asymmetry of the u quark in the proton:

$$A_L^{W^+} = \frac{u_-(x_1)\bar{d}(x_2) - u_+(x_1)\bar{d}(x_2)}{u_-(x_1)\bar{d}(x_2) + u_+(x_1)\bar{d}(x_2)} = \frac{\Delta u(x_1)}{u(x_1)}. \quad (18)$$

Similarly, for Figure 17b alone,

$$A_L^{W^+} = \frac{\bar{d}_-^+(x_1)u(x_2) - \bar{d}_+^+(x_1)u(x_2)}{\bar{d}_-^+(x_1)u(x_2) + \bar{d}_+^+(x_1)u(x_2)} = -\frac{\Delta \bar{d}(x_1)}{\bar{d}(x_1)}. \quad (19)$$

In general, the asymmetry is a superposition of the two cases:

$$A_L^{W^+} = \frac{\Delta u(x_1)\bar{d}(x_2) - \Delta \bar{d}(x_1)u(x_2)}{u(x_1)\bar{d}(x_2) + \bar{d}(x_1)u(x_2)}. \quad (20)$$

To obtain the asymmetry for W^- , one interchanges u and d .

For the pp collisions at RHIC with $\sqrt{s} = 500$ GeV, the quark will be predominantly a valence quark. By identifying the rapidity of the W , y_W , relative to the *polarized* proton, we can obtain direct measures of the quark and antiquark polarizations, separated by quark flavor: $A_L^{W^+}$ approaches $\Delta u/u$ in the limit of $y_W \gg 0$, whereas for $y_W \ll 0$ the asymmetry becomes $-\Delta \bar{d}/\bar{d}$. Higher-order corrections change the asymmetries only a little [?, ?].

The kinematics of W production and Drell-Yan production of lepton pairs is the same. The momentum fraction carried by the quarks and antiquarks, x_1 and x_2 (without yet assigning which is which), can be determined from y_W ,

$$x_1 = \frac{M_W}{\sqrt{s}} e^{y_W}, \quad x_2 = \frac{M_W}{\sqrt{s}} e^{-y_W}. \quad (21)$$

Note that this picture is valid for the predominant production of W s at $p_T \sim 0$. The experimental difficulty is that the W is observed through its leptonic decay $W \rightarrow l\nu$, and only the charged lepton is observed. We therefore need to relate the lepton kinematics to y_W , so that we can assign the probability that the polarized proton provided the quark or antiquark. Only then will we be able to translate the measured parity-violating asymmetry into a determination of the quark or antiquark polarization in the proton.

Missing: Discussion of x_1, x_2 and helicity plot

The rapidity of the W is related to the lepton rapidity in the W rest frame (y_l^*) and in the lab frame (y_l^{lab}) by

$$y_l^{\text{lab}} = y_l^* + y_W, \quad \text{where } y_l^* = \frac{1}{2} \ln \left[\frac{1 + \cos\theta^*}{1 - \cos\theta^*} \right]. \quad (22)$$

Here θ^* is the decay angle of the lepton in the W rest frame, and $\cos\theta^*$ can be determined from the transverse momentum (p_T) of the lepton with an irreducible uncertainty of the sign [?], since

$$p_T^{\text{lepton}} = p_T^* = \frac{M_W}{2} \sin\theta^*. \quad (23)$$

In this reconstruction, the p_T of the W is neglected. In reality, it has a p_T , resulting for example from higher-order contributions such as $gu \rightarrow W^+d$ and $u\bar{d} \rightarrow W^+g$, or from primordial p_T of the initial partons.

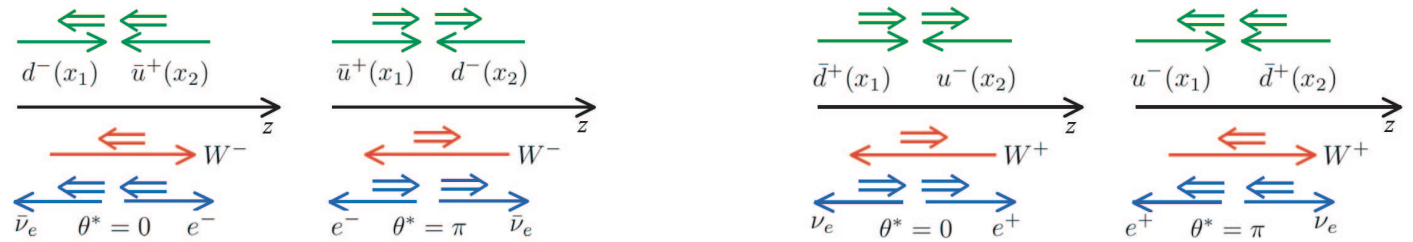


Figure 19:

Usually W production is identified by requiring charged leptons with large p_T and large missing transverse energy, due to the undetected neutrino. Since none of the detectors at RHIC is hermetic, measurement of missing p_T is not available, which leads to some background. Possible sources of leptons with high p_T include charm, bottom, and vector boson production. Above $p_T \geq 20$ GeV/c, leptons from W decay dominate, with a smaller contribution from Z^0 production. Both PHENIX and STAR can estimate the single-lepton Z^0 background from measured Z^0 production. The additional background from misidentified hadrons is expected to be small.

Missing: RHICBOS YIELD for STAR

Expected yields were estimated with PYTHIA [?] and RESBOS [?]. The cross section at RHIC for W^+ (W^-) production is about 1.3 nb (0.4 nb). These estimates vary by 5–10% according to the choice of the parton distribution set. For 800 pb^{-1} and $p_T \geq 20 \text{ GeV}/c$, PHENIX expects about 8000 W^+ s and 8000 W^- s in the muon arms (that the numbers are equal is due to the decay angle distribution and acceptance), as well as 15,000 W^+ and 2500 W^- electron decays in the central arms. STAR, with its large acceptance for electrons, expects 72,000 W^+ s and 21,000 W^- s. Using Equation 21 to reconstruct x , Figure 20 shows the expected sensitivity for $\Delta f(x)/f(x)$, with $f = u, d, \bar{u}, \bar{d}$, for the PHENIX muon data.

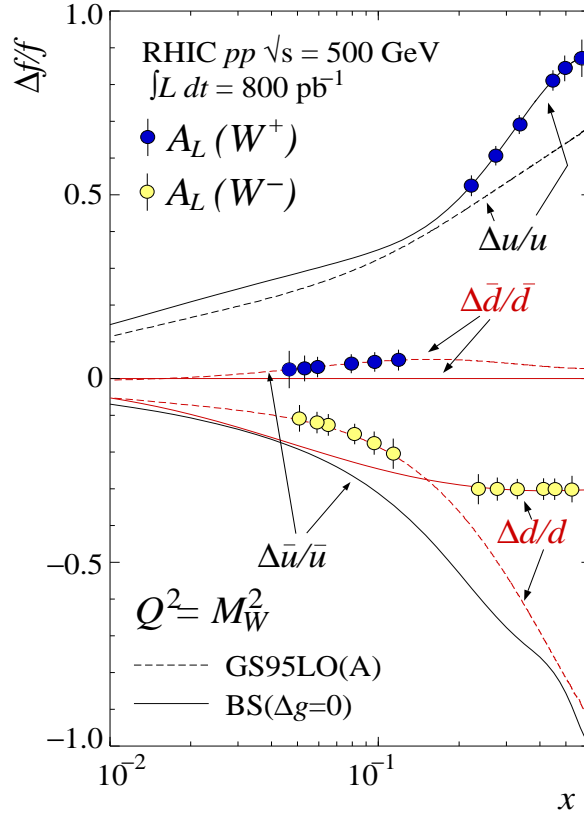


Figure 20: Expected sensitivity for the flavor-decomposed quark and antiquark polarization overlayed on the parton densities of Reference [?] (BS) and of Reference [?] [GS95LO(A)]. Darker points and error bars refer to the sensitivity from $A_L(W^+)$ measurements, and lighter ones correspond to $A_L(W^-)$.

Missing: RHICBOS

RHIC will also significantly contribute to our knowledge about the unpolarized parton densities of the proton, since it will have the highest-energy pp collisions. $p\bar{p}$ production of W s has a much stronger valence component in the determined [?] $u(x)/d(x)$ ratio. Isospin dependence in Drell-Yan production of muon pairs in pp, pd scattering [?], violation of the Gottfried sum rule [?, ?], and recent semi-inclusive DIS measurements [?] have shown that the unpolarized sea is not $SU(2)$ symmetric. At RHIC, the ratio of unpolarized W^+ and W^- cross sections will directly probe the \bar{d}/\bar{u} ratio, as shown in Figure 23.

RHICBOS W simulation (P=0.7 L=400pb⁻¹)

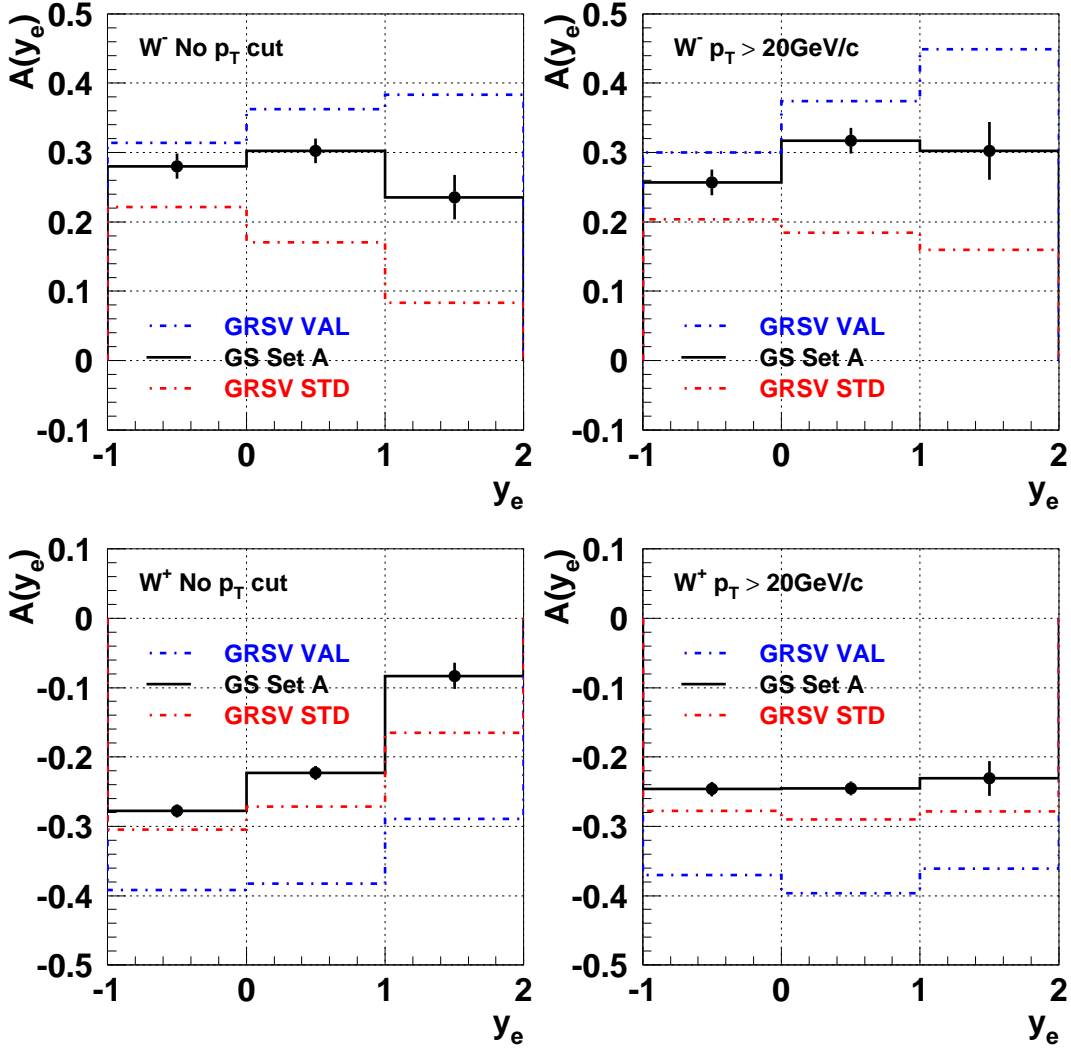


Figure 21:

2.7 Transverse spin structure

With the proton spin transversely polarized with respect to its momentum or the collision axis, a novel helicity flip *chiral-odd* twist-2 quark distribution - known as the transversity distribution, $\delta q(x)$, is theoretically allowed [42, 43]. The relation between the leading twist quark distribution, $q(x)$, helicity distribution, $\Delta q(x)$, and the transversity distribution, $\delta q(x)$, is best represented in terms of a (double) density matrix notation [75]

$$\mathcal{F}(x) = \frac{1}{2}q(x) I \otimes I + \frac{1}{2}\Delta q(x) \sigma_3 \otimes \sigma_3 + \frac{1}{2}\delta q(x) (\sigma_+ \otimes \sigma_- + \sigma_- \otimes \sigma_+) , \quad (24)$$

for the quark distribution function $\mathcal{F}(x)$ of a nucleon. In Eq. (24), the first matrix in the direct product is in the nucleon helicity space and the second in the quark helicity space. The transversity distribution $\delta q(x)$ is as fundamental as $q(x)$ and $\Delta q(x)$ in QCD, and has its unique

RHIC polarized pp W production at 500GeV cms energy

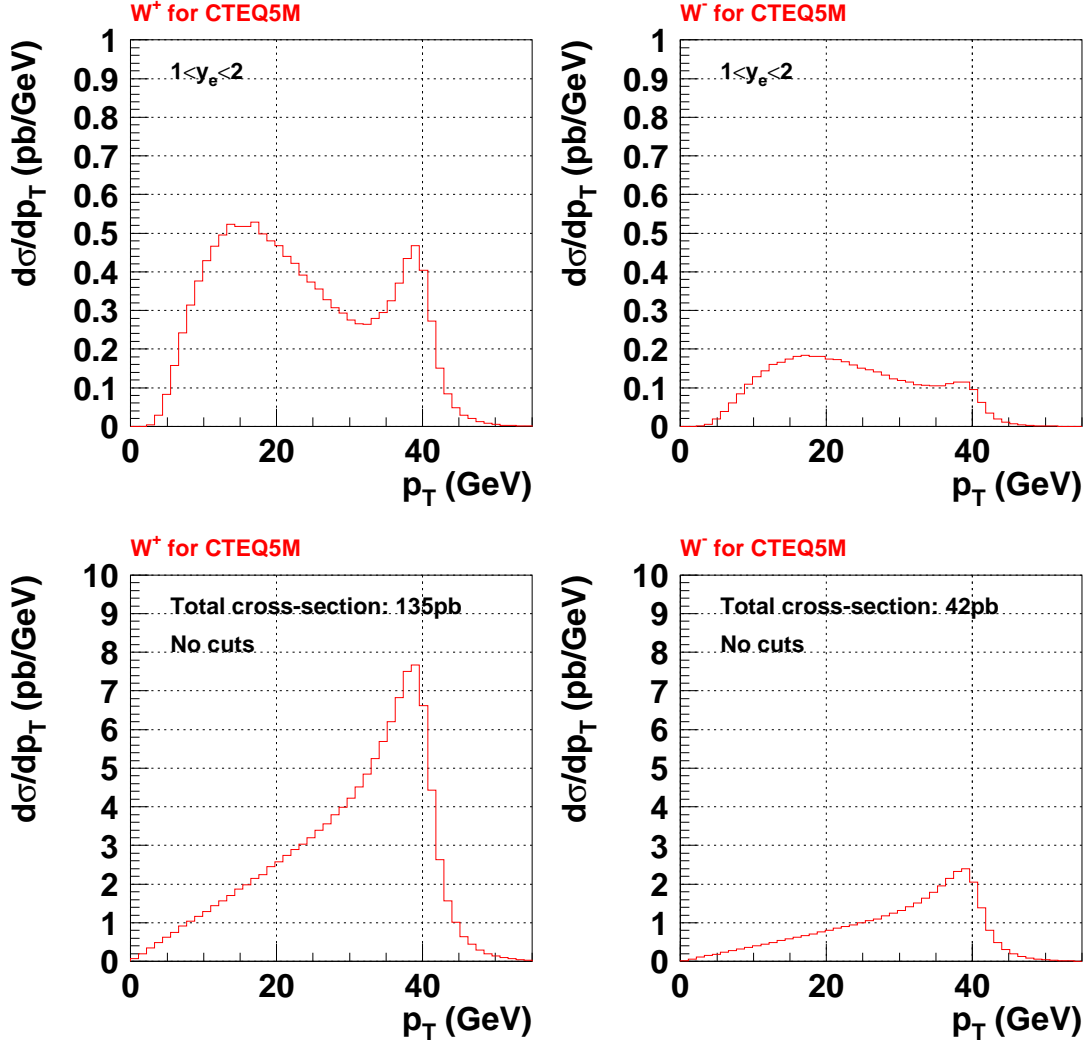


Figure 22:

factorization scale dependence [76, 77, 78] and transverse spin sum rule [79]

$$\frac{1}{2} = \frac{1}{2} \sum_{q, \bar{q}} \int dx \delta q^a(x) + \sum_{q, \bar{q}, G} \langle L_{\mathbf{s}_T} \rangle^a$$

where $L_{\mathbf{s}_T}$ is the component of orbital angular momentum along the transverse spin of nucleon \mathbf{s}_T . But, it is not completely independent because of Soffer's Inequality [80],

$$|2 \delta q(x)| \leq q(x) + \Delta q(x),$$

which is valid for each quark flavor q . Independent measurements of $q(x)$, $\Delta q(x)$, and $\delta q(x)$ and their factorization scale dependence can provide a direct test of QCD dynamics. Measurements of $\delta q(x)$ are important to compare with lattice QCD results for the tensor charge of the nucleon [81]. Furthermore, $\delta q(x)$ and $\Delta q(x)$ are not trivially related because boosts and rotations do not

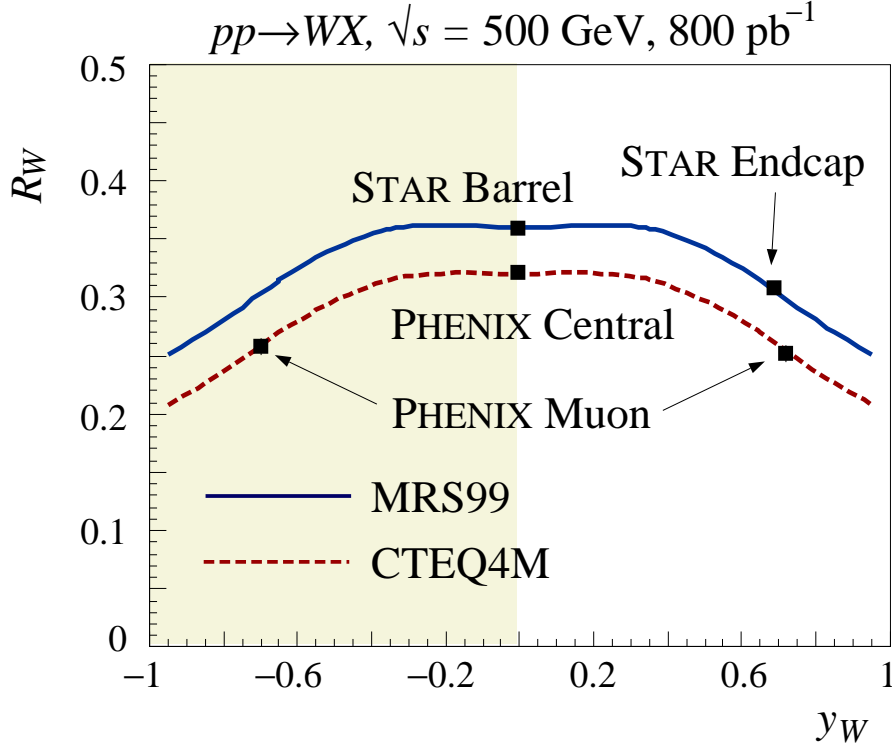


Figure 23: The ratio $R_W = (d\sigma(W^-)/dy)/(d\sigma(W^+)/dy)$ for unpolarized pp collisions at RHIC. The shaded region indicates that unpolarized pp collisions are symmetric in y_W . To illustrate the sensitivity of the measurement, we show an earlier set of parton densities (CTEQ4M [?]) and a set (MRS99 [?]) that includes the latest information from Drell-Yan data [?]. Both curves include an asymmetric sea with \bar{d}/\bar{u} rising to 1.6 for increasing antiquark momentum fraction $x_{\bar{q}}$, but the latter also includes a drop-off in the ratio for higher $x_{\bar{q}}$.

commute. The difference between these two distributions carries important information about the non-perturbative structure of the proton.

It is very important to point out that there is no leading twist gluon transversity in comparison to the longitudinal case where a role of gluon spin may be responsible for violation of Ellis-Jaffe sum rule. This is because gluon transversity would require two units of helicity flip that the nucleon density matrix cannot provide.

$\delta q(x)$ is difference between probability finding quark with parallel and anti-parallel transverse spin in transversely polarized nucleon. But unlike for $q(x)$ and $\Delta q(x)$, measurement requires helicity flip amplitudes. Since perturbative hard processes conserve helicity, chiral-odd distributions must appear in pairs. Although transversity distribution can be in principle extracted from the measurements of *double* transverse spin asymmetries, $A_{TT} \propto \delta q(x) \otimes \delta q'(x')$, the asymmetries are often too small to be useful because of the dominance of gluonic contribution to the unpolarized cross sections. Therefore, $\delta q(x)$ is better determined from observables dominated by quark-initiated partonic processes, like A_{TT} of Drell-Yan [42, 43], or *single* transverse spin asymmetries (SSA), $A_N \propto \delta q(x)$ with an aid of another chiral-odd unpolarized non-perturbative functions, like the Collins function [82] described below.

Single longitudinal-spin asymmetries for single particle inclusive production vanish due to

parity and time-reversal invariance. Transverse single spin asymmetries are not forbidden by these basic symmetries. Because of Lorentz invariance of QCD, we need at least four vectors including the spin vector to construct a physically observed SSA. With a proton spin vector S not parallel to its momentum, a hadron level SSA can be constructed to be proportional to $\epsilon_{\mu\nu\alpha\beta} S^\mu P_A^\nu P_B^\alpha p^\beta$ with beam momenta, P_A and P_B , and observed particle momentum p in single hadron inclusive production, $P_A(\mathbf{S}) + P_B \rightarrow h(p_T) + X$. Significant single transverse-spin asymmetries, A_N , of ten or more percent of the unpolarized cross sections, were recorded by Fermilab E704 experiment in the beam fragmentation region of hadronic π production at p_T as large as GeV [83]. Since then, nonvanishing single transverse-spin asymmetries have been observed in lower energy hadronic collisions [84] and semi-inclusive lepton-hadron deeply inelastic scattering (SIDIS) [45, 85], as well as, in much higher energy pp collisions at RHIC [86].

However, theoretically, it was pointed out a long time ago [87] that perturbative QCD calculation at leading power in collinear factorization formalism predicts vanishing single transverse-spin asymmetries, $A_N \propto \alpha_s m_q / p_T$, in inclusive single hadron production at large p_T . This is because the SSA requires a hadron-level helicity flip and is proportional to a T-odd combination of the vectors, $A_N \propto \mathbf{S} \cdot (\mathbf{P}_A(\text{or } \mathbf{P}_B) \times \mathbf{p}_T)$. Within the leading twist collinear factorization formalism, the hadron helicity flip requires a quark helicity flip, which leads to the m_q dependence, while the suppression of a power of α_s is due to the T-odd combination which requires a phase (or an imaginary part) from one of the two amplitudes.

Major theoretical progress has been made in last decade in understanding the “unexpected”, but, observed large SSA. It is believed that SSA is a unique and excellent probe for studying parton’s transverse motion and strength of color Lorentz force inside a bound nucleon.

Within the collinear factorization formalism, the parton-level helicity flip can be achieved from the interference between an amplitude of a spin (1/2) quark state and a spin (-1/2) quark-gluon composite state without requiring a quark helicity flip and the quark mass [88, 89]. The required phase for the SSA can be generated from a partonic pole, which corresponds to the non-local feature of the composite quark-gluon state and leads to a natural growth of the SSA into the fragmentation region [89, 90]. Because of the coherent interference of a single parton state and a two-parton composite state, corresponding nonperturbative matrix elements, known as high twist matrix elements, do not have probability interpretations. Since parton’s transverse momentum k_T is integrated over in defining the matrix elements, quark-gluon correlation functions in collinear factorization formalism provide *averaged* (or integrated) information on partons’ transverse motion, like averaged color Lorentz force experienced by the quarks [89]. These correlation functions are new physical observables for probing nonperturbative QCD dynamics.

If we go beyond the collinear factorization formalism, the necessary hadron-level helicity flip for SSA can be achieved without requiring a parton-level helicity flip because of parton’s transverse motion or orbital angular momentum. A k_T -dependent (or un-integrated) parton distribution provides a direct information on both longitudinal and transverse motion of a parton inside a bound nucleon. As pointed out by Sivers [?], a k_T -dependent quark distribution of a transversely polarized nucleon, could have both symmetric and antisymmetric terms when nucleon spin $\mathbf{S} \rightarrow -\mathbf{S}$. Sivers function is an initial state correlation between transverse spin of nucleon \mathbf{S}_T and the parton transverse momentum in the nucleon \mathbf{k}_T of the form $\mathbf{S}_T \cdot (\mathbf{P} \times \mathbf{k}_T)$, where \mathbf{P} is the nucleon momentum. The antisymmetric term, known as the Sivers function, could be a source of nonvanishing single transverse spin asymmetries. Similarly, the Sivers mechanism can

apply to the un-integrated gluon distribution to define a gluonic Sivvers function [91]. The Sivvers functions should be directly related to parton's transverse motion and orbital angular momentum. Understanding this connection is an active area of theoretical research [92].

It is also possible to generate a SSA by combining a non-zero quark transversity distribution together with the Collins-Heppelmann effect[93]. Since transversity is a chiral-odd function, measurements of transversity require another chiral-odd function. This can be a polarized chiral-odd fragmentation function (FF) which acts as an analyzer of the transversely polarized quark. The Collins FF [82] displays itself as a correlation of the form $\mathbf{S}_T \cdot (\mathbf{P}_{jet} \times \mathbf{k}_T)$ where \mathbf{S}_T is the transverse spin vector, \mathbf{P}_{jet} is the jet momentum and \mathbf{k}_T is the transverse momentum of the hadron with respect to the fragmenting quark. This FF depends not only on the momentum fraction of the hadron with respect to the parton $z = E_h/E_{parton}$, but also on k_T . Early on it was realized that the measurement of two hadrons (h_1 and h_2) within a jet would also be sensitive to the Collins effect via the correlation $\mathbf{S}_T \cdot (\mathbf{P}_{h_1} \times \mathbf{P}_{h_2})$ [93]. But, a recent theoretical study[96] suggests the Collins effect may be suppressed for pion production in p+p collisions due to cancellations induced by quantum phases. It was also pointed out that interference between s-wave and p-wave mesons can analyze the transverse polarization of a quark[94]. This so-called interference FF has an advantage in that the theory can work within a collinear approximation, free from gluon radiation effects which can modify the asymmetry. The Interference FF depends on the invariant mass of the two hadrons and total momentum fraction $z = z_{h_1} + z_{h_2}$. A model calculation[94] suggests that there will be sign change around the ρ mass for fragmentation into $\pi^+ + \pi^-$. Recent data from the HERMES collaboration provides some indication of the existence of this change in sign [95].

Much of the predictive power of QCD is provided by the universality of nonperturbative parton distributions and/or fragmentation functions in factorization theorems for hadronic processes. In order to quantify and measure the parton's transverse motion inside a polarized nucleon, it is necessary to have a gauge invariant definition of k_T -dependent parton distributions and/or fragmentation functions, from which Sivvers functions and/or Collins functions can be better defined [97, 98, 99, 100]. To extract the k_T -dependent distributions, such as the Sivvers and Collins functions, from physical observables, a factorization formalism in terms of these universal distributions is required. Since the Sivvers functions and Collins functions are sensitive to the transverse motion of partons at relatively low parton transverse momentum k_T , which is nonperturbative, another large physical scale, $Q \gg k_T$, is needed to ensure the k_T -dependent factorization [101]. For example, Drell-Yan pair of invariant mass Q at low transverse momentum $q_T \sim k_T$ is a good probe of k_T -dependent quark distribution while the large Q ensures the factorization [101]. Sivvers functions could also be measured in terms of asymmetric di-jet correlations at RHIC [102].

The k_T -factorization formalism was used to calculate the SSA in hadronic pion production, A_N , in terms of contributions of Sivvers functions [103] and/or a combination of transversity and Collins functions [104]. Although the k_T -factorization for the inclusive pion production in hadronic collisions has not been formally proved, the calculated A_N , with x_F dependence mainly determined by the extracted Sivvers and/or Collins functions from fitting low energy data, are consistent with new RHIC data at $\sqrt{S} = 200$ GeV [105]. When pion's p_T is much larger than the typical k_T of parton's transverse motion in a nucleon, $p_T \gg \langle k_T \rangle$, the k_T factorization formalism used in these calculations is not expected to be valid. On the other hand, the proved collinear factorization formalism at twist-3 for hadronic A_N should be a good approximation [89]. Measurements of the SSA and its transition from the low p_T region where $A_N \propto p_T/\langle k_T \rangle \sim$

$k_T/\langle k_T \rangle$, which is sensitive to the parton's transverse motion k_T , to the high p_T region where $A_N \propto \langle k_T \rangle/p_T$, which probes to the *averaged* transverse momentum or *averaged* color Lorentz force [89], should provide excellent information on parton's transverse motion.

With the ability to measure correlations of two hadrons (or jets) at RHIC, it is possible to separate Sivers effect, parton's transverse motion in initial-state hadron wave functions, from Collins effect, parton's spin effect in final-state hadronization. Sivers functions, Collins functions, transversity distribution and quark-gluon correlations all provide information on nucleon's spin structure that cannot be reached by measurements of longitudinal spin asymmetries. Single transverse spin asymmetry, A_N , opens a unique window to see partons' transverse motion inside a polarized nucleon.

There have been several measurements of transverse single spin asymmetries for pion production initiated by collisions of polarized proton. The E704 collaboration observed [83] that the analyzing power (A_N) for pions produced in polarized proton collisions at $\sqrt{s}=20$ GeV at large Feynman x had magnitudes up to $\sim 30\%$ that increased with increasing x_F . A_N was found to be positive for π^+ production, negative for π^- production and positive for π^0 production, although smaller in magnitude than for π^+ production. Similar trends for the x_F dependence of A_N for pion production are also observed for pion production at lower \sqrt{s} [107]. More recently, new polarized p+p experiments[84] as well as semi-inclusive deep-inelastic lepton scattering (SIDIS) experiments[45, 85] have reported measurements of transverse single-spin asymmetries (SSA) which are significantly different from zero. The intrigue surrounding the results has sparked a wave of experiments to explore the unexpectedly large signals. For example, the HERMES collaboration has devoted significant effort to measure SIDIS with a transversely polarized target [108]. At RHIC, A_N for pion production has been measured at $\sqrt{s} = 200$ GeV over a large rapidity range[86, 122] as shown in figure24 and figure25. At forward rapidity, the large value of A_N was found to be large for positive x_F , even at collision energies a factor of 10 larger than earlier experiments. The STAR experiment has more recently reported preliminary A_N results at $x_F < 0$ which are consistent with zero within the current experimental accuracy[122]. According to a recent study[96], the large negative x_F region is dominated by gluon Sivers effects. While the preliminary STAR data may already exclude a saturated gluon Sivers function, more data are required to establish its size.

There has been tremendous theoretical progress towards understanding of transverse spin effects within QCD, as reviewed at the most recent International Symposium on Spin Physics [106]. In part, the significant theoretical progress has been stimulated by new experimental results for transverse SSA results from SIDIS experiments and from polarized proton collisions at RHIC. There is promise that this progress can be sustained since we also know that the unpolarized yield for large x_F π^0 production in p+p collisions at $\sqrt{s} = 200$ GeV is well described with fixed order pQCD calculations, unlike at lower \sqrt{s} [124].

Higher precision measurements of A_N over a wide range of x_F and p_T at $\sqrt{s} \geq 200$ GeV will provide the opportunity to understand the physics behind these large spin effects. With a data sample 3 pb^{-1} and $P = 0.5$, the statistical precision for the negative x_F analyzing power for π^0 production [122] will be improved by a factor of 5. Measuring the p_T dependence of A_N can also be accomplished with such a data sample and will provide a crucial test of these pQCD based models, as they all predict a $1/p_T$ dependence. Figure 26 shows the statistical error projection for A_N as function of p_T at fixed $x_F = 0.5$ for π^0 production at STAR. Obtaining precise data

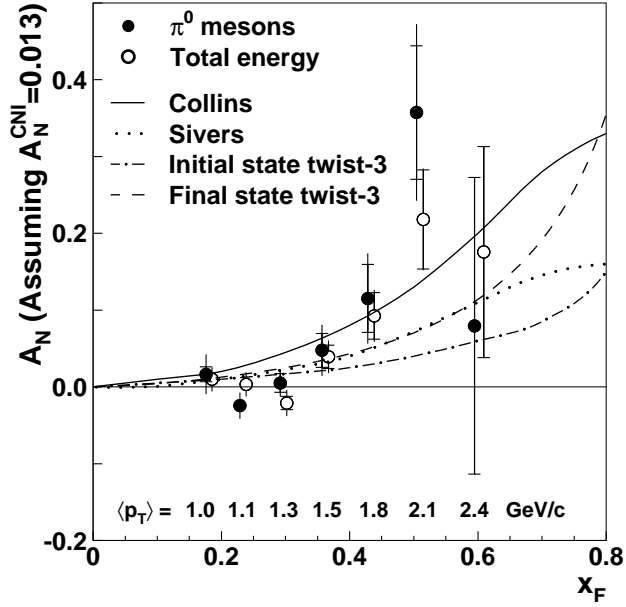


Figure 24: The A_N for π^0 produced in the interval $3.3 < \eta < 4.1$ for polarized proton collisions at $\sqrt{s}=200$ GeV from STAR[86]. The solid points are for identified π^0 mesons. The open points are measurements of the total energy shifted by 0.01 in x_F . The curves are predictions from pQCD models evaluated at $p_T=1.5$ GeV/c.

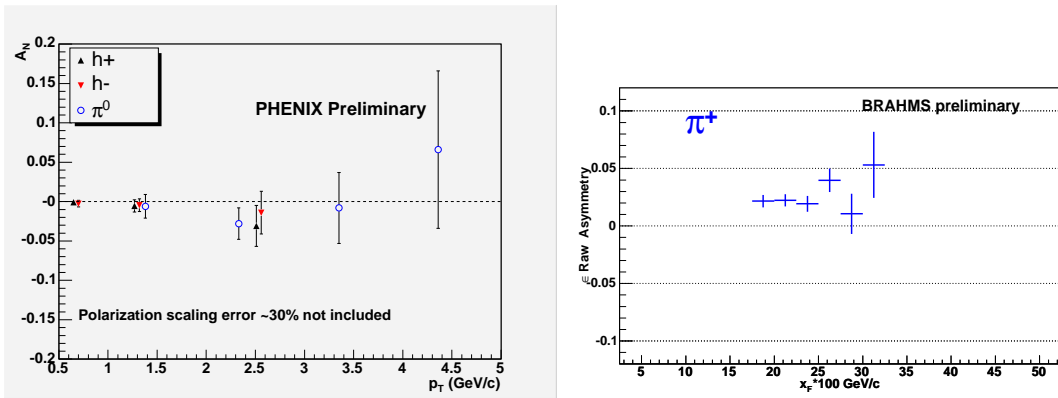


Figure 25: (Left) A_N for $\pi^{0,\pm}$ production at polarized proton collisions at $\sqrt{s}=200$ GeV as a function of p_T at mid-rapidity from PHENIX[123]. (Right) Preliminary A_N for π^+ production in the interval $2.3^\circ < 3.5^\circ$ as function of x_F from BRAHMS.

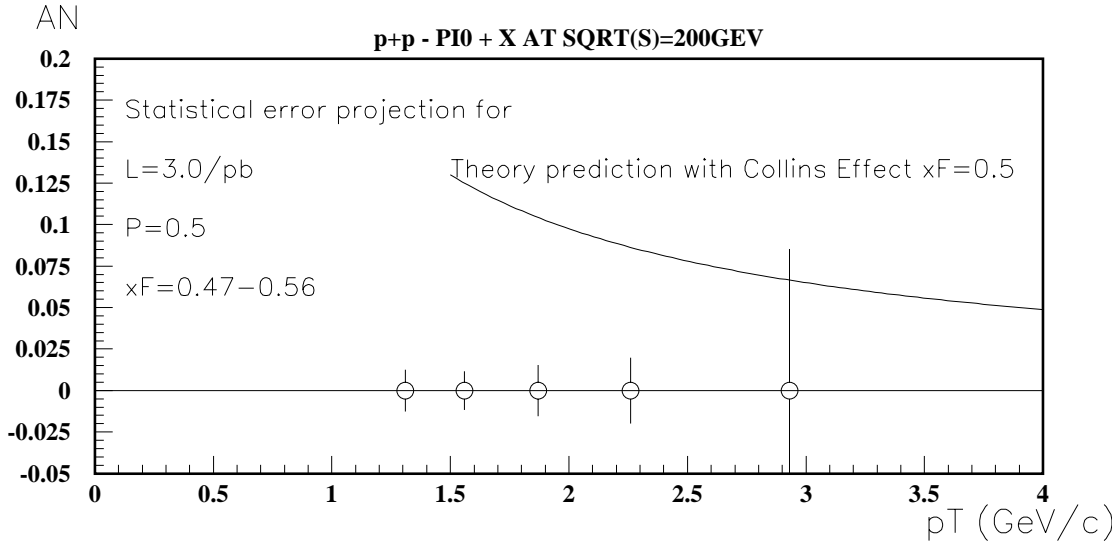


Figure 26: Statistical error projection for A_N as function of p_T for π^0 production at STAR. A theory prediction for Collins effect for $x_F = 0.5$ is also shown.

for at very large x_F is interesting in order to test the expectation of A_N to decrease[112], rather than continuing to increase with increasing x_F , due to the Soffer inequality [80] bounding the transversity distribution function. It is also expected that comparable precision will be achieved for similar sized data samples for π^\pm production at $\sqrt{s}=200$ GeV by the BRAHMS collaboration.

To disentangle these different mechanisms behind A_N , it is required to go beyond inclusive measurements and measure either 2 final state particles or jet. There is a class of observables for which the k_T -dependent distributions or fragmentation functions appears at leading-power, and are directly sensitive to a small measured transverse momentum[82, 93, 113]. A goal of the transverse spin program at RHIC is to disentangle the potential contributions to A_N by directly measuring the intrinsic k_T of partons in nucleon or transverse momentum in fragmentation process.

Back-to-back di-jet production in unpolarized and transversely polarized pp collisions can be used to measure the Sivers function[113]. The deviation of the azimuthal angle difference $\delta\phi = \phi_{jet1} - \phi_{jet2} + \pi$ from zero directly measure the k_T of the partons. When the jet production plane is close to the spin of the nucleon, the k_T imbalance of the parton will affect the $\delta\phi$ distribution and can give spin asymmetry $A_N(\delta\phi) = \frac{\sigma^\uparrow(\delta\phi) - \sigma^\downarrow(\delta\phi)}{\sigma^\uparrow(\delta\phi) + \sigma^\downarrow(\delta\phi)}$. At mid rapidity, this asymmetry may reach a few percent or more and could provide access to the gluon Sivers function, as shown in figure 27. The error bars are estimated statistical uncertainties for di-jet measurement at STAR with two luminosity and polarization assumptions based on existing data[114] and expected increase in the acceptance. Improvements of trigger to enrich di-jet events is currently investigated, and it is possible to surpass this estimation.

Measurements at RHIC are planned to study the Collins FF for jets and the Interference FF for two particles within a jet with unpolarized and transversely polarized pp collisions. Figure 28 shows projections for the analyzing power for $\pi^+ - \pi^-$ coincident pairs produced at midrapidity in p+p collisions at $\sqrt{s}=200$ GeV after simulating the acceptance and resolution of PHENIX. The projections assume that transversity saturates the Soffer bound and assume the interference FF from Ref. [94].

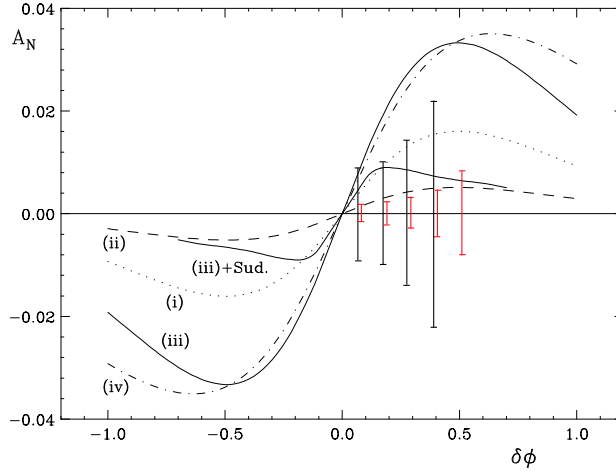


Figure 27: Predictions for the spin asymmetry A_N for back-to-back di-jet production at $\sqrt{s} = 200$ GeV, for various different models for the gluon Sivers function. The solid line marked as “(iii)+Sud” shows the impact of leading logarithmic Sudakov effects on the asymmetry for model (iii)[113]. The error bars are estimated statistical uncertainties for di-jet measurement at STAR with 4 pb^{-1} , beam polarization of 40% (black) and 30 pb^{-1} , 70% (red).

By increasing the coverage of electromagnetic calorimetry in the forward direction, spin-dependent particle correlation studies are enabled in the rapidity range where large transverse single spin effects have already been observed. With calorimetry that spans $2.5 < \eta < 4.0$, coincident $\pi^0 - \pi^0$ pairs can be observed at large rapidity. Near-side correlations ($|\Delta\phi| = |\phi_{\pi,1} - \phi_{\pi,2} + \pi| \approx 0$) can be used to identify di-hadron fragments from a jet produced at large rapidity. In general, di-hadron fragmentation functions are not well constrained. Nonetheless, their rate can be estimated by the Lund string model, as implemented in PYTHIA 6.222 [129].

There are two interesting scenarios to consider, each involving azimuthally correlated $\pi^0 - \pi^0$ pairs indicative of di-hadron fragmentation of a forward jet.

In the Sivers picture A_N should be associated with the forward jet. We should therefore expect that the large A_N observed for a single π^0 with $x_{F1} > 0.4$ produced at $\eta_1 \approx 3.8$ would also be present for forward jets that fragment into forward $\pi^0 - \pi^0$ pairs having $x_{F,1} + x_{F,2} > 0.4$. From PYTHIA, we expect $\approx 1.5 \times 10^4$ $\pi^0 - \pi^0$ events with $x_{F1} > 0.25$ and $x_{F2} > 0.15$ in the near-side jet-like peak in $\Delta\phi$ for 1 pb^{-1} of integrated luminosity for p+p collisions at $\sqrt{s}=200$ GeV. PYTHIA predicts the near-side $\Delta\phi$ peak sits atop a uniform background. The signal to background ratio is approximately 1 : 1. With corrections for background, an accuracy of $\delta A_N \approx 0.01$ could be achieved with 3 pb^{-1} of integrated luminosity and 70% beam polarization.

The Collins mechanism attributes A_N to the correlation $\mathbf{S}_T \cdot (\mathbf{P}_{h_1} \times \mathbf{P}_{h_2})$ involving the momenta of two hadronic fragments of a jet and the proton spin vector. The transverse momentum associated with jet fragmentation that produces a π^0 with $x_{F1} > 0.4$ and $3 < \eta_1 < 4$ can be determined by detecting a second forward π^0 with $x_{F2} > 0.15$ and by requiring the $\pi^0 - \pi^0$ pair have $|\eta_1 - \eta_2| < 0.5$. Again, PYTHIA predicts the $\Delta\phi$ correlation has a jet-like near-side correlation peak sitting atop a uniform background. For these kinematics, we expect 2×10^3 $\pi^0 - \pi^0$ pairs in the near-side correlation peak for 1 pb^{-1} of integrated luminosity. If a non-zero Collins effect is observed, then larger integrated luminosity samples would be required to map out the x dependence of the transversity structure function.

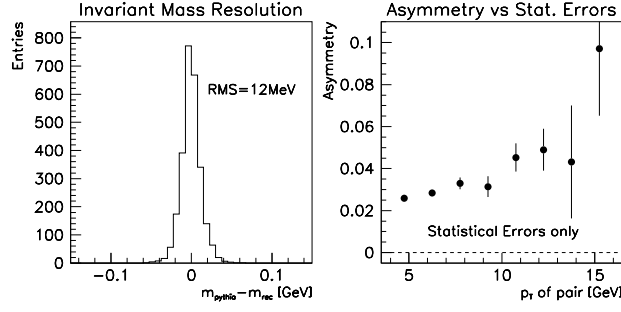


Figure 28: A statistical error projection for Collins asymmetry measurement at PHENIX with Soffer bound saturated transversity and interference FF model[94] at 30 pb^{-1} and $P = 0.5$

In addition, Collins and Interference FF's are being measured at $e^+ + e^-$ colliders[115] such as Belle, where analysis is currently ongoing[116]. Once these FF's are known, measurements at RHIC can be used to access the transversity distribution function.

As stated earlier, transversity measurement always need the two chiral-odd functions, and both can be transversity. Transverse double-spin asymmetries A_{TT} for jet or high- p_T particle production will be sensitive to the product of 2 transversity functions. The advantage of this method is that it does not require any knowledge of polarized fragmentation functions. On the other hand, this asymmetry is highly suppressed because there is no gluon transversity and also because double chirality flip process are color suppressed. It is important that A_{TT} measurements are done in a kinematical region where gluon contributions are small, since there is no gluon transversity. Such measurements require the full luminosity and polarization of RHIC. There are other ideas to access transversity via the production of J/ψ [119] and D mesons[120]. Perhaps the cleanest way to measure transversity is through the Drell-Yan process[118], but this likely requires a RHIC luminosity upgrade to have sufficient sensitivity.

Table 2: Physics cases with transverse spin and its luminosity and polarization requirements

Channel	Luminosity [pb^{-1}]	Polarization
Inclusive A_N	0.4	0.2
Mapping A_N in x_F and p_T space	3	0.4
Sivers from di-jet	10	0.5
Transversity from di-hadron correlations within a jet	30	0.5
A_{TT} of jet or high- p_T particle	100	0.7
DY	1000	0.7

Table 2 is a summary of physics channels with transverse spin and its luminosity and polarization requirement. There are achievable goals with a very low luminosity and polarization requirements. The first polarized proton collisions at RHIC have already produced exciting transverse spin asymmetry results. As the luminosity and polarization increase, more channels will become accessible. It is reasonable to expect that measurements of both the Sivers function and transversity can be achieved with $\sim 20 \text{ pb}^{-1}$ of integrated luminosity with polarization of 0.5. To achieve this goal, and minimize the impact on the longitudinal spin physics cases, we will devote about 1/3 to 1/4 of beam time for transverse spin physics. It is important to note that PHENIX and STAR have separate spin rotators; the two experiments can choose longitudinal and transverse spin independently with minimum overhead to re-tune the beam with difference configuration.

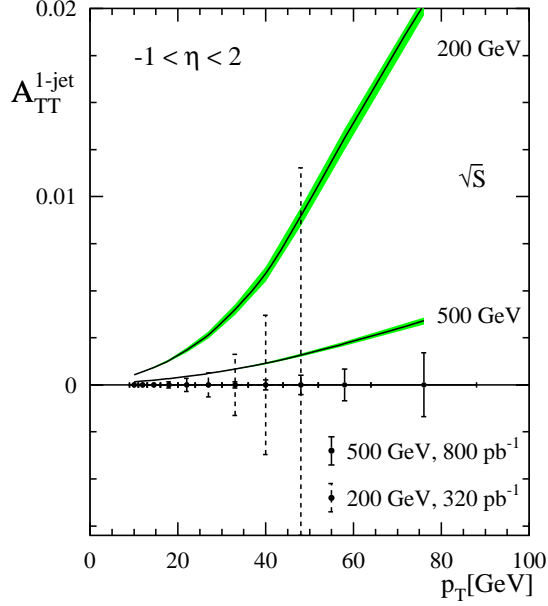


Figure 29: Maximally possible A_{TT} for single-inclusive jet production at $\sqrt{s} = 200$ and 500 GeV as a function of p_T . Jet rapidities are integrated over $-1 < \eta < 2$. The shaded bands represent the theoretical uncertainty in A_{TT} estimated by varying scale by factor 2. Also indicated as error bars is the expected statistical accuracy with design luminosity of the RHIC [117].

Most of transverse spin measurements prefers the $\sqrt{s} = 200$ GeV for larger asymmetries.

2.8 What else is going on around the world?

The fundamental nature of the nucleon spin puzzle and the high interest in understanding its spin structure as has lead to experimental programs around the world, to study this problem. In addition to RHIC spin at BNL, they are the COMPASS experiment at CERN in Geneva Switzerland, the HERMES experiment at DESY and the Jefferson Laboratory's Continuous Electron Beam Accelerator Facility (CEBAF) in Newport News, VA in the US. In this section we emphasize the essential features of their physics program.

2.8.1 COMPASS experiment at CERN

The COMPASS (Common Muon and Proton Apparatus for Structure and Spectroscopy) collaboration at CERN has built a dual purpose experiment, aiming to study hadron spectroscopy and hadron structure. The main spin objectives are to measure: 1) the polarization of gluons within polarized nucleons, 2) light quark helicity distributions by flavor, 3) lambda and anti-lambda polarization, 4) transverse spin distributions. Leptoproduction of open charm and of hadron pairs with large transverse momenta are considered to be the most promising options to measure $\Delta g/g$.

The experiment is located in the CERN-SPS M2 beamline of longitudinally polarized 160

GeV muons, which was used also by the former EMC and SMC experiments. The kinematic coverage is thus similar. A solid polarized target filled with ^6LiD provides both longitudinal and transverse target polarizations. A newly designed large-angle spectrometer is used to reconstruct the scattered muons and the produced hadrons in wide momentum and angular ranges. Muon, pion, kaon, and proton identification are among its distinctive features.

The first data, collected in 2002 and 2003 using predominantly longitudinal configurations of the beam and target spins, correspond to integrated luminosities of about 600 and 900 inverse picobarn, respectively. A much smaller sample with transverse target polarizations has been collected. The recently released results on the inclusive structure function g_1^d are in agreement with those from previous experiments and provide a factor two improvement in statistical uncertainty in the region $0.004 < x < 0.03$ [127]. Similar improvements in the semi-inclusive measurements of quark helicities by flavor may be anticipated.

While the improved measurements of the inclusive structure function may lead to some refinement in the indirect extraction of gluon polarization via NLO pQCD analysis, the main COMPASS objective is of course in the direct determination of $\Delta g/g$ via the production of open charm and of high p_T hadron pairs. A first preliminary result on $\Delta g/g$ from high p_T hadron pairs [?] has a precision similar to those from the preceding HERMES and SMC analyses. The first open charm measurements are statistics limited.

The COMPASS detector will not take data in 2005 but the spin program is expected to resume in 2006, sharing its beam time with the spectroscopy program. The spectrometer acceptance will benefit greatly from an upgrade to the target magnet planned to be installed in 2005.

2.8.2 HERMES experiment at DESY

The HERMES experiment at DESY studies the spin structure of the nucleon by scattering the lepton beam in HERA off pure internal targets that can be polarized longitudinally and transversely. The large acceptance of the detector allows for inclusive and semi-inclusive measurements and, like the COMPASS spectrometer, has particle identification. The longitudinal program has delivered measurements of: 1) the inclusive structure function g_1 of the proton and neutron with similar precision and kinematic coverage as preceding measurements at SLAC, 2) detailed semi-inclusive measurements which allow the extraction of quark helicity distributions functions with different model assumptions than preceding work [128], and 3) the first analysis of spin effects in lepton production of high p_T hadron pairs, a process which is sensitive to Δg .

The present focus is on transverse and nuclear spin effects. HERMES has made the first measurement of the single-spin asymmetries for semi-inclusive electroproduction of charged pions in deep-inelastic scattering of positrons with a transversely polarized hydrogen target [128]. The asymmetry depends on the azimuthal angles of the pion and the target spin axis about the virtual photon direction and relative to the positron scattering plane. Transverse quark polarization in the target nucleon and a correlation between the intrinsic transverse quark momentum and the transverse target polarization give rise to signal of this nature, and can be distinguished in experiment with this technique. The latter might provide an indication for non-zero orbital momentum of quarks in the polarized nucleon. The HERMES program is foreseen to continue to the summer of 2007.

2.8.3 Nucleon Spin related experiments at Jefferson Laboratory

Unlike the above two experimental detectors, the CEBAF accelerator at Jefferson laboratory is a superconducting radio frequency electron accelerator facility that was commissioned during early 1990s. The accelerator uses state of the art photocathode gun that is capable of delivering beams of high polarization and high currents ($\sim 200 \mu\text{A}$) to two of its three experimental halls while maintaining high polarizations but low current ($\sim \text{few nA}$) beam to the third. Polarizations of the order of 85% are routinely achieved, and maximum beam energies are 5.5 GeV. The main machine consists of a "racetrack" layout of circulating beam line with two linear accelerators joined by two 180° arcs. The accelerator design also allows experiments with brehmsstrahlung photons off of the continuous electron beams in experimental areas with fluxes ranging from $10^7\text{--}8/\text{sec}$.

The relatively low beam energies (compared to COMPASS, HERMES and in general RHIC) available at CEBAF and the fixed target experimental halls A, B, and C [?] makes the facility ideally suited to study the nucleon structure (including spin) in the transition region from non-perturbative to perturbative QCD. This is also the region where higher twist effects play an important role and hence can be studied at JLab. The high intensity electron beams coupled with solid or gaseous targets allows exclusive measurements using the high acceptance detectors in the halls[?]. With these characteristics of the experimental conditions, Jlab experiments explore regions of $Q^2 \rightarrow 0$ with high statistical accuracies. To date some of the most accurate tests of the GDH spin sum rule [?], high precision high x spin structure functions of the nucleon[?], azimuthal single spin asymmetries relating to transverse dynamics of the partonic spin structure and measurements of a class of reactions such as DVCS (deeply virtual compton scattering) which are proposed to lead us to the measurement of the orbital angular momentum of the quarks[?], are results of experiments performed at Jefferson Lab. It is the proposed measurement of the orbital angular momentum of the partons inside the nucleon, those of transversity related structure functions, and their role in understanding the overall spin structure of the nucleon makes Jlab measurements complementary to those of Δg and transversity with RHIC Spin.

2.9 Elastic Scattering of polarized high energy protons

The previous sections discuss the physics of hard scattering at RHIC with polarized protons, which can be understood as collisions of polarized quarks and gluons. The scattering is so energetic that we can use perturbative QCD to describe the interactions of the quarks and gluons, and, thus, probe the spin structure of the proton at very small distances. For example, scattering at $Q^2=(80 \text{ GeV})^2$ probes wave lengths of 0.003 fermi. Small-angle scattering, from total cross section to $t = -1 \text{ (GeV}/c)^2$, probes the static proton properties and constituent quark structure of the proton, covering distances from 4 fermi [$-t = 0.003 \text{ (GeV}/c)^2$ in the Coulomb nuclear interference (CNI) region] to a distance of ≈ 0.2 fermi. Unpolarized scattering shows striking behavior in this region, from the surprise that total cross sections rise at high energy, to observed dips in elastic cross sections around $-t = 1 \text{ (GeV}/c)^2$. Both polarimetry at RHIC and the PP2PP experiment explore this region for spin-dependent cross sections, for $\sqrt{s}=7\text{--}500 \text{ GeV}$, for the first time.

In the very forward region, the nuclear and electromagnetic amplitudes are of comparable

magnitude, and the interference between them results in a small but significant maximum in the single transverse spin asymmetry A_N making elastic scattering in the CNI region useful for polarimetry [?]. Important results have already been obtained in the RHIC spin program in this region. There have been measurements made near RHIC injection energy (24 GeV/c) using a carbon micro-ribbon target both at the AGS (E950) and in RHIC; there have been measurements made with a 100 GeV/c beam on a carbon target and independently on a gas jet target. In addition, a measurement in the colliding beam mode has been carried out (pp2pp).

CNI scattering produces an asymmetry from the scattering of an unpolarized particle, a proton in one RHIC beam or a carbon nucleus in a fixed target, from the anomalous magnetic moment of a polarized proton, with a maximum of $A_N = 0.04$ at $-t = 0.003$ (GeV/c)². However, a hadronic spin-flip term can also contribute to the maximum, and this term is sensitive to the static constituent quark structure of the proton. The authors of Reference [?] remark that the helicity flip probes the shortest interquark distance in the proton, and that the helicity nonflip is sensitive to the largest quark separation in the proton due to color screening. The helicity-flip term, if present, can indicate an isoscalar anomalous magnetic moment of the nucleons [?], an anomalous color-magnetic moment causing helicity nonconservation at the constituent quark-gluon vertex [?], and/or a compact quark pair in the proton [?, ?].

The results from scattering 100 GeV/c polarized protons on carbon [?] are shown on the left side of Fig. 2.8.1. An asymmetry reaching 0.02 is observed, with a t -dependence quite different from pure electro-magnetic spin flip (the top curve in the figure). The curves use a standard CNI form [?], with the lower curve including hadronic spin flip. Results from scattering 100 GeV/c protons on a highly polarized hydrogen jet target [?] are shown on the right side of Fig. 2.8.1. In this figure, the solid line is the prediction with only electromagnetic spin flip. For pp scattering at this energy ($\sqrt{s}=14$ GeV), no hadronic spin flip is observed. Further, the preliminary measurement from the pp2pp experiment with colliding beams at RHIC is also shown on the right side of the figure. For $\sqrt{s}=200$ GeV, the asymmetry is somewhat larger than the CNI curve without hadronic spin flip, but, including present uncertainty on the beam polarization (the measurement was made in 2003 before the jet provided more precise polarization results), the result is consistent also with no hadronic spin flip for pp scattering.

Small-angle scattering at high energy is presently understood in a Regge picture as being dominated by Pomeron exchange [?]. The Pomeron, which has the vacuum quantum numbers with charge-conjugation $C = +1$, can be interpreted as a two-gluon exchange. These results for A_N for carbon and proton targets imply that the isospin 0 Reggeons (which include the Pomeron) have a significant spin flip coupling for the carbon target. The $I = 1$ Regge poles for the proton target scattering must be sufficiently strong to nearly cancel the $I = 0$ contribution at this energy. The couplings required have been determined and indicate that asymptotically the Pomeron will contribute about 10% spin-flip; i.e. the cancellation leading to no spin-flip in pp at 100 GeV/c will go away as the energy increases [?].

Many other measurements can be made at RHIC, for elastic scattering and also for diffractive scattering with rapidity gap measurements. The two-spin transverse asymmetry in the CNI region, for example, is sensitive to a C-parity odd exchange, referred to as the Odderon [?]. For larger t , a steep exponential fall with momentum transfer, characteristic of pomeron exchange matches on to an approximate t^{-8} dependence at larger $-t$ in the unpolarized cross sections. The latter has a natural interpretation in terms of three vector exchanges between pairs of valence

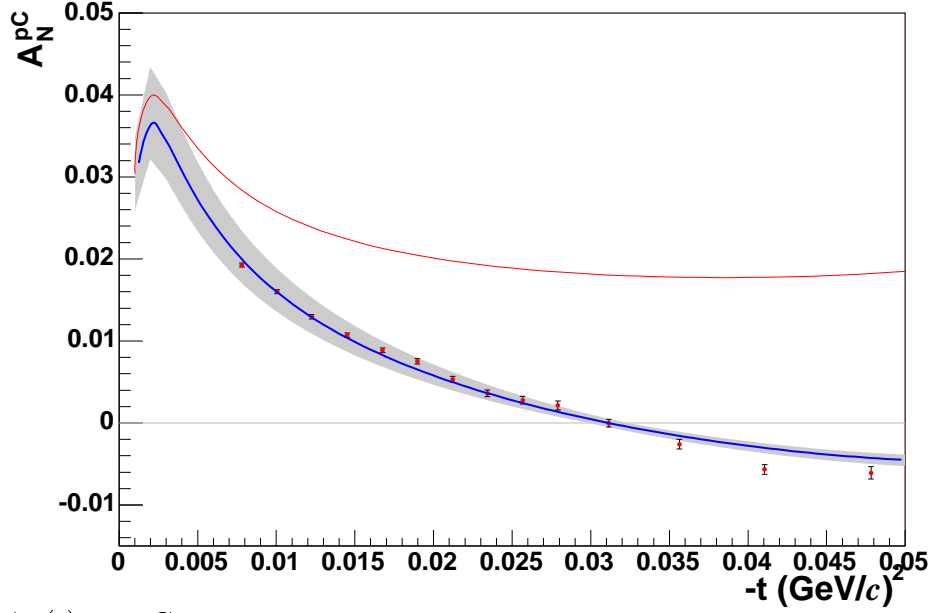


Figure 30: $A_N(t)$ for pC elastic scattering at 100 GeV, left side [?]. The shaded band represents the systematic uncertainties of the measurement. The solid line in the band is a fit to the data including a significant hadronic spin-flip contribution (see text). The result is significantly different from the no hadronic spin-flip prediction (top curve). The right side shows preliminary measurements (solid points) of proton-proton elastic scattering with 100 GeV protons incident on a highly polarized atomic hydrogen jet target [?]. The open squares are data from E704 at Fermilab [?]. In this case, the curve with no hadronic spin flip describes the data well. The closed box is the preliminary result from the colliding beam experiment pp2pp, $\sqrt{s}=200$ GeV [?]. A 20% systematic uncertainty for the beam polarization is not shown. The dashed curve is the pure CNI prediction for this energy. Again, no hadronic spin flip appears to be required.

quarks. Whether these individual scatterings should be thought of as single gluons, or as (at least in part) perturbative exchanges in color-singlet configurations remains to be seen. This profile is fairly stable with energy, even as the details of its shape change. The observation of a stable profile in polarized elastic scattering at RHIC would surely initiate a new class of theoretical investigations. Lastly, the dramatic spin dependence of proton-proton elastic scattering at moderate $-t$ observed in the Argonne and BNL experiments of twenty years ago remains an outstanding puzzle [?]. This could also be explored at RHIC.

At higher energy, such as at the LHC, the CNI region becomes inaccessible. The minimum $-t$ reachable with colliding beams depends on scattering the protons out of the beams. For fixed $-t$, the scattering angle falls as $1/p_{\text{beam}}$, whereas the beam size falls more slowly as $1/\sqrt{p_{\text{beam}}}$. Roughly, this limits an experiment at the LHC to $-t > 0.01$ (GeV/c) 2 .

2.10 Physics beyond the Standard Model

Single beam helicity asymmetries violate parity and give access to searches for potential physics beyond the Standard Model, for example for quark substructure, new neutral gauge bosons present in some supersymmetric models, and supersymmetric particle production. In general, parity violation searches can compete with the sensitivity of much higher energy unpolarized

colliders, and a parity violation signal beyond known electro-weak effects would be a decisive signal for new physics. Furthermore, if a new interaction is discovered for example at the LHC, and lower-lying masses are accessible at RHIC, RHIC will be able to explore the chiral structure of the new interaction. We discuss the potential new particles and mass ranges where RHIC can contribute in this section. For sensitivities we generally consider the target luminosity for $\sqrt{s}=500$ GeV of $\sim 1 \text{ fb}^{-1}$ and a large acceptance detector, for example STAR with high luminosity capability. However, a suggested new detector[?] and an order of magnitude higher luminosity from a RHIC II can also be considered. Fig. 31 shows the single beam helicity asymmetry A_L

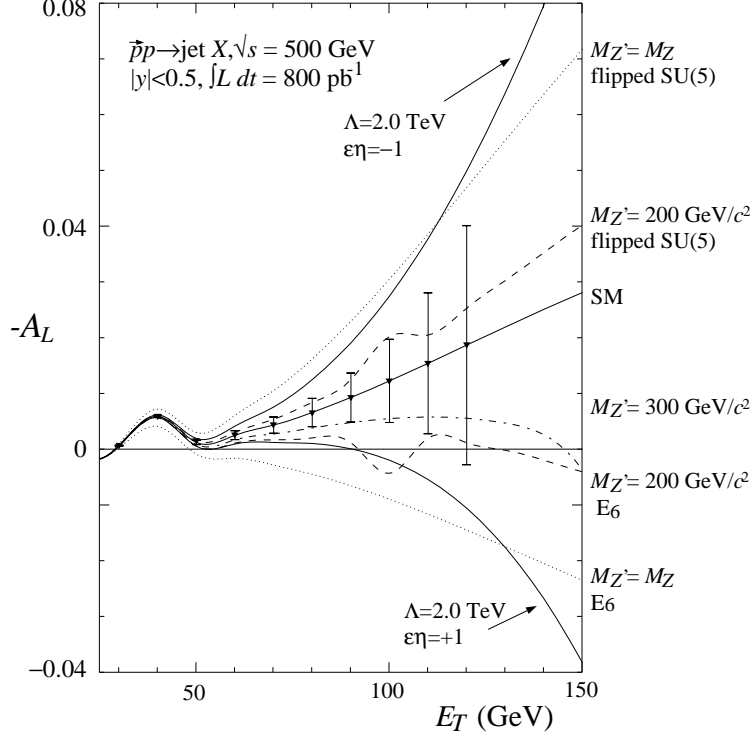


Figure 31: A_L , for one-jet inclusive production in $\vec{p}p$ collisions versus transverse energy, for $\sqrt{s} = 500$ GeV. The solid curve with error bars represents the SM expectations. The error bars show the sensitivity at RHIC for 800 pb^{-1} , for the STAR detector. The other solid curves, labeled by the product of $\epsilon\eta$, correspond to the contact interaction at $\Lambda = 2 \text{ TeV}$ [135]. The dashed and dotted curves correspond to different leptophobic Z' models. The calculations are at leading order.

for single jet production versus the transverse energy of the jet, for the Standard Model and for extensions of the Standard Model (SM). The SM predicts a parity violating asymmetry, around $E_T = M_W/2$ for $W \rightarrow 2 \text{ jet}$ production, and from QCD-electroweak interference. Indeed, the data in the peak provides a calibration for the SM effects. The figure uses the top RHIC energy, our target luminosity for this energy, and the acceptance of the STAR detector [135]. The sensitivities shown do not include detector efficiency. Two extensions of the SM are considered: quark compositeness, and new neutral gauge bosons that appear in several string-derived models [140](non-supersymmetric models may be also constructed [141]).

The quark compositeness curves in the figure use a compositeness scale $\Lambda=2 \text{ GeV}$ and maximal parity violation $\eta = \pm 1$. Composite models of quarks and leptons [131] generally violate

parity, since the scale of compositeness $\Lambda_c \gg M_W$. The present limit from unpolarized collisions, at the Tevatron, is [132] $\Lambda \cong 1.6$ TeV. Due to the direct sensitivity to parity violation, RHIC can compete with the much higher energy Tevatron. Further, if an anomalous parity violation signal is observed, it would be a definitive observation of new physics. The limits of sensitivity for Λ in the contact model of quark compositeness [135] are tabulated in Table 3 for RHIC, the Tevatron and the LHC, with $L \sim 1 \text{ fb}^{-1}$ integrated luminosity the RHIC spin target for $\sqrt{s}=500$ GeV, and $L \sim 10 \text{ fb}^{-1}$ a potential target for RHIC II.

Collider	$\sqrt{s}(\text{TeV}) L(\text{fb}^{-1})$	Λ (TeV)	
RHIC $p \uparrow p$	0.5	1	3.3
	0.5	10	5.5
Tevatron limit	1.8	0.5	1.6
Tevatron $p\bar{p}$	2.0	30	5
LHC pp	12	100	30

Table 3: Limits on quark compositeness (Λ) at 95% CL for different colliders. RHIC uses $P=0.7$, STAR detector acceptance, single jet production, parity violation signal with maximal parity violation for compositeness. The Tevatron and LHC use the deviation from the Standard Model for the jet inclusive cross section versus p_T [135]. **What systematic errors are assumed??**

Also shown in Fig. 31 are Z' curves. In the framework of supersymmetric models with an additional Abelian $U(1)'$ gauge, it has been shown [142] that the Z' boson could appear with a relatively low mass ($M_Z \leq M_{Z'} \leq 1$ TeV) and a mixing angle with the standard Z close to zero. The effects of different representative models are also shown in Fig. 31 (see Ref. [136] for details). RHIC covers some regions of parameters space of the different models that are unconstrained by present and forthcoming experiments, and RHIC would also uniquely obtain information on the chiral structure of the new interaction. In Ref. [137], it has been suggested to extend this study to the collisions of polarized neutrons, which could be performed with colliding at RHIC polarized ^3He nuclei [143]. The authors argue that, in case of a discovery, a compilation of the information coming from both polarized $\vec{p}\vec{p}$ and $\vec{n}\vec{n}$ collisions should constrain the number of Higgs doublets and the presence or absence of trilinear fermion mass terms in the underlying model of new physics.

There are also a number of other examples. Using the RHIC target luminosity (1 fb^{-1}) at $\sqrt{s}=500$ GeV, parity violating asymmetry in the production of XX from supersymmetric sparticles with mass to 75 GeV would be observable [144]. Another example is to search for a transverse asymmetry for W or Z boson production, where the SM prediction is very small [150]). W^\pm and Z^0 production in $p^\uparrow p$ collisions at RHIC is expected to have good sensitivity on $\mathcal{L}_{SM}-\delta\mathcal{L}$ interference at the parton level due to strong correlation between the proton spin and the polarization of high- x valence quarks, that participate in the gauge boson production [154]. This asymmetry could arise from anomalous electroweak dipole moments of quarks [150, 151, 152]. The RHIC sensitivity can improve present experimental limits by a factor ~ 5 -10 [152]. These limits do not approach the SM expectations, and a signal would be direct indication of new physics.

2.11 Connection to eRHIC

The addition of a high energy, high polarization lepton (electron/positron) beam facility to the existing RHIC Complex, able to collide with its hadron beam, would dramatically increase RHIC's capability to do precision QCD physics. Such a facility with 10 GeV/c polarized electrons/positrons has been proposed and is called eRHIC. There are many connections between the RHIC spin program and eRHIC. We categorize them in two groups:

- *Direct connections to RHIC Spin:* In these, the physics observables measured by the existing RHIC spin physics program will be measured in complementary kinematic regions, or in some cases augmented to complete the understanding of nucleon spin.
- *Indirect Connections to RHIC Spin:* These include measurements not possible with RHIC Spin, but of significance to understanding QCD with spin in general or nucleon spin in particular.

2.11.1 Direct Connections

Direct connections between RHIC Spin and eRHIC are made on three principal topics : the measurement of the polarized gluon distribution, the measurement of polarized quark-anti-quark distributions, and on transverse physics measurements.

For polarized gluon distribution measurements, eRHIC enables an increase in the kinematic range and precision, particularly at low x . At eRHIC the polarized gluon distribution will be measured using a) the scaling violations of spin structure functions $g_1^{p/n}$ and b) di-jet and high p_T di-hadron production in the photon gluon fusion process [?]. The RHIC spin measurements discussed earlier in this document will be most significant in the medium-high x range, $x > 10^{-2}$, while eRHIC will complement them with precision on low x , ($x < 10^{-2}$) all the way to $x \sim 10^{-4}$.

RHIC Spin will be the first to measure in a model independent way the polarized quark and anti-quark distributions using single spin longitudinal asymmetry measurements in pp scattering at $\sqrt{s}=500$ GeV/c via (W^\pm) production. Analysis of these asymmetries will give us $\Delta u, \Delta \bar{u}, \Delta d, \Delta \bar{d}$???. The quark-anti-quark separation in such a way is not possible in fixed target DIS where the virtual γ is the propagator of the force which cannot differentiate between quarks and anti-quarks. However, at high enough energy, in DIS at eRHIC, virtual W^\pm also get exchanged. If $\Delta q = u, \bar{u}, d, \bar{d}$ are known by early next decade from RHIC Spin, eRHIC will be able to continue this program by exploring the heavy quarks, i.e. identify the spin contributions from $\Delta c/\bar{c}$ and $\Delta s/\bar{s}$. Of course, traditional methods to get quark flavor distributions (quark-antiquark unseparated) using semi-inclusive DIS measurements of charged and neutral pions and kaons will also continue, with access to flavor separation at lower x than is possible in current fixed target DIS experiments.

Transversity is the last as yet unmeasured spin structure function, discussed in detail in ???. The measurements at RHIC with pp scattering will be made using measurements of Collins Fragmentation Function (CFF), Interference Fragmentation Functions (IFF) and if very large luminosities are achieved, also with Drell Yan (DY) processes [?]. These measurements will be made

in the center of mass energy range from 200 to 500 GeV. The eRHIC will make a complementary set of measurements, with high precision using CFF and IFF measurements, not unlike those made by the HERMES collaboration currently.

2.11.2 Indirect Spin Connections

In addition to the measurements eRHIC will do that will extend or complement the investigation of nucleon spin with RHIC Spin, there is another class of nucleon spin and other helicity related measurements that could also be made with eRHIC. A partial list includes:

- Measurement of spin structure functions g_1 of the proton and neutron and the difference between them that tests the Bjorken spin sum rule. eRHIC will do this with accuracies that will for the first time start competing and challenging the experimental systematic uncertainties at the level of 1- 2%. Low x phenomena have been some of the most exciting aspects of the physics coming from unpolarized DIS measurements in the last decade, and eRHIC will probe low x kinematics for the first time with polarized beams
- eRHIC will be the only possible facility in the foreseeable future at which QCD spin structure of the virtual photon could be explored. The process employed for this investigation is that of photon-gluon fusion [?].
- Deeply virtual Compton scattering (DVCS) for final state photons as well as other vector mesons measured using almost complete acceptance (4π) detectors has been suggested as a preliminary requirement towards the measurement of the Generalized Parton Distributions (GPDs). A series of different GPD measurements may be required eventually to extract the orbital angular momentum of the partons. This is the last part of the nucleon spin puzzle which we may have to address after the spin of the gluon is understood. Although the theoretical formulation is not yet ready, it is expected that by the time eRHIC comes on line, there will be a formalism available to take the measured GPDs and determine the orbital angular momentum of partons. These measurements at eRHIC will be complementary, at much higher energy scales, to those being planned at Jefferson Laboratory after its 12 GeV upgrade plan.
- Drell Hearn Gerasimov spin rule measurements presently underway at Jefferson laboratory [?] and at MAMI [?] are mostly at low values of ν [?]. While the significance of the contribution the spin sum rule from high ν is small, absolutely no measurements exist beyond the value of $\nu > \approx 1$ GeV. eRHIC will extend direct measurements of the high ν component up to 500 GeV.
- Precision measurements of spin structure functions in very high $x \sim 0.9$ region could be part of the eRHIC physics program with specially designed detectors as has been discussed in [?].

The physics programs with polarized proton beams at RHIC and eRHIC have much in the way of complementarity of physics measurements. It's also clear that success at eRHIC passes through a period of successful measurements and collider development by the RHIC spin program not only at $\sqrt{s}=200$ GeV/ c but also at $\sqrt{s}=500$ GeV/ c .

3 Accelerator performance

As of 2004, polarized proton beams have been accelerated, stored and collided in RHIC at a proton energy of 100 GeV. The average store luminosity reached $4 \times 10^{30} \text{cm}^{-2} \text{s}^{-1}$, and the average store polarization 40% (see Tab. 4). Over the next 4 years we aim to reach the Enhanced Luminosity goal for polarized protons, consisting of an average store luminosity of

- $60 \times 10^{30} \text{cm}^{-2} \text{s}^{-1}$ for 100 GeV proton energy, and
- $150 \times 10^{30} \text{cm}^{-2} \text{s}^{-1}$ for 250 GeV proton energy,

both with an **average store polarization of 70%**. Tab. 4 gives a projection of the luminosity and polarization evolution through FY2008. Luminosity numbers are given for 100 GeV proton energy and one of two interaction point. For operation with more than two experiments, the luminosity per interaction point is reduced due to an increased beam-beam interaction. For each year the maximum achievable luminosity and polarization is projected. Projections over several years are not very reliable and should only be seen as guidance for the average annual machine improvements needed to reach the goal. We assume that 10 weeks of physics running are scheduled every year to allow for commissioning of the improvements and development of the machine performance.

Table 4: Maximum projected RHIC polarized proton luminosities through FY2008. Delivered luminosity numbers are given for 100 GeV proton energy and one of two interaction points. 10 weeks of physics operation per year are assumed. The designation 2002A refers to achieved, and 2005E refers to expected.

Fiscal year		2002A	2003A	2004A	2005E	2006E	2007E	2008E
No of bunches	...	55	55	56	79	79	100	112
Protons/bunch, initial	10^{11}	0.7	0.7	0.7	1.0	1.4	2.0	2.0
β^*	m	3	1	1	1	1	1	1
Peak luminosity	$10^{30} \text{cm}^{-2} \text{s}^{-1}$	2	6	6	16	31	80	89
Average luminosity	$10^{30} \text{cm}^{-2} \text{s}^{-1}$	1.5	3	4	9	21	53	60
Time in store	%	30	41	41	50	53	56	60
Max luminosity/week	pb^{-1}	0.2	0.6	0.9	2.8	6.6	18.0	21.6
Max integrated luminosity	pb^{-1}	0.5	1.6	3	20	46	126	151
Average store polarization	%	15	30	40	45	65	70	70
Max LP ⁴ /week	nb^{-1}	0.1	5	23	120	1180	4330	5190

In Fig. 32 the integrated luminosity delivered to one experiment is shown through FY2012 for two scenarios: 10 weeks of physics operation per year, and 10 weeks of physics operation every other year. For every projected year shown in Fig. 32 the weekly luminosity starts at 25% of the final value, and increases linearly in time to the final value in 8 weeks. During the remaining weeks the weekly luminosity is assumed to be constant. For the maximum projection the values in Tab. 4 are used as final values until FY2008. For later years the FY2008 values are assumed with no further improvement. The minimum projection is one third of the maximum projection, based on past experience in projecting heavy ion luminosities [155].

For the scenario with 10 weeks of physics operation per year, the assumed proton energy is 100 GeV until FY2009, and 250 GeV thereafter. For the scenario with 10 weeks every other year, the assumed proton energy is 100 GeV throughout the entire period to reach the physics goal.

For the scenario with 10 weeks of physics operation every other year, the final values are not increased in years without proton operation, since no time is available to develop the machine performance. Thus in our projections we reach the Enhanced Luminosity goal in FY2008 with 10 week physics operation per year, but need until FY2011 with 10 weeks of physics operation every other year. For operation at 250 GeV proton energy, the luminosity projections in Tab. 4 need to be multiplied by 2.5. We expect no significant reduction in the averages store polarization after full commissioning of polarized proton ramps to 250 GeV.

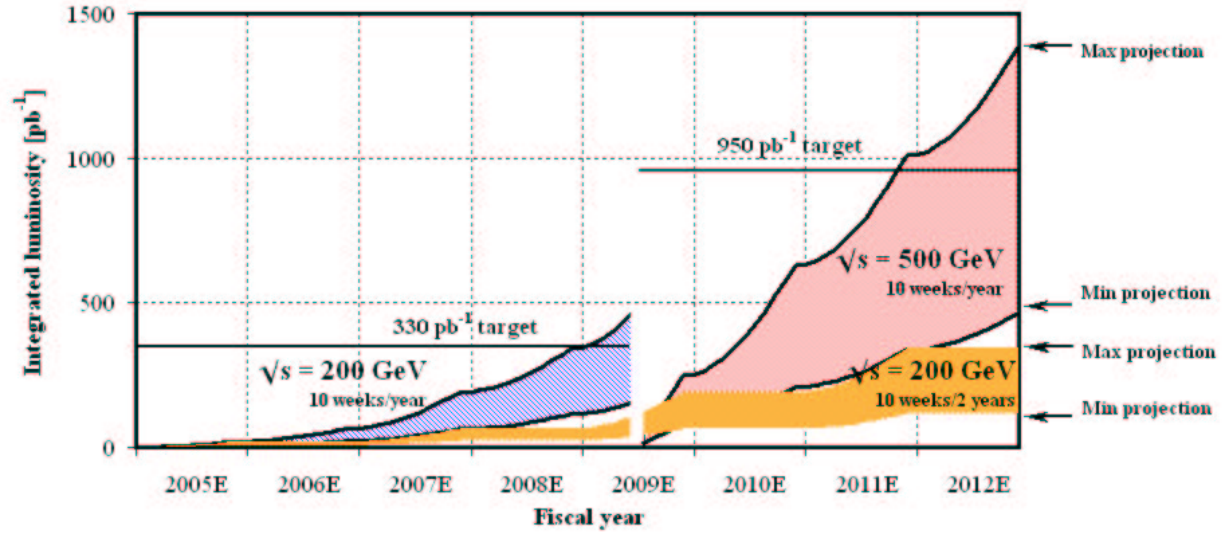


Figure 32: Minimum and Maximum projected integrated luminosity through FY2012. Delivered luminosity numbers are given for one of two interaction points. For the scenario with 10 weeks of physics operation per year, the assumed proton energy is 100 GeV until FY2009, and 250 GeV thereafter. For the scenario with 10 weeks every other year, the assumed proton energy is 100 GeV throughout the entire period.

3.1 Polarization limitations

The RHIC beam polarization at 100 GeV is currently limited by the AGS beam polarization transmission efficiency of about 70%, and the source polarization. With the installation of a new solenoid in FY2005, the source polarization is expected to increase from 80% to 85%. The existing AGS polarized proton setup includes a 5% warm helical snake for overcoming imperfection spin depolarizing resonances and an RF dipole for overcoming 4 strong intrinsic spin resonances. This setup has two drawbacks:

1. All the weak intrinsic spin resonances are crossed with no correction and result in a total depolarization of about 16%.
2. Operation with the RF dipole still leads to about 15% depolarization.

In addition, the AGS has shown a dependence of the beam polarization on the bunch intensity. These shortcomings can be overcome with the installation of a new AGS cold snake, to be initially commissioned in 2005. With a scheme that combines the AGS cold snake of 15%, and the AGS warm snake of 5%, depolarizations at all imperfection and all intrinsic spin resonances should be

eliminated, making the AGS spin transparent with the exception of some mismatch at injection and extraction [156].

Obtaining 70% beam polarization in RHIC at 250 GeV is challenging because of strong intrinsic and imperfection resonances beyond 100 GeV. Betatron tunes and orbit distortions have to be controlled precisely to avoid depolarization due to snake resonances. Simulations show that orbit distortions have to be corrected to less than 0.3 mm rms. Orbit errors are introduced due to misalignments and remain if the orbit cannot be corrected completely. A realignment of the entire ring is scheduled for the 2005 summer shutdown. Efforts continue to improve the existing beam position monitor system, and the orbit correction techniques. A beam-based alignment technique is under development. With the existing hardware and software, orbit distortions of 1 mm rms were achieved, as measured by the beam position monitors. Acceleration of polarized proton beams beyond 100 GeV is planned in 2005. The result of this machine development effort will provide guidance for the tolerable levels of machine misalignments and orbit errors.

3.2 Luminosity limitations

A number of effects limit the achievable luminosity. Currently the bunch intensity is limited to about 1×10^{11} to maintain maximum polarization in the AGS. This restriction should be removed with the AGS cold snake. With intense bunches the beam-beam interaction will limit the luminosity lifetime. With bunches of 2×10^{11} protons and 2 interaction points, the total beam-beam induced tune spread will reach 0.015. Operation with more than two collision will significantly reduce the luminosity lifetime. RHIC is also the first hadron collider to operate in a strong-strong beam-beam regime. High intensity beams also lead to a vacuum breakdown, caused by electron clouds. In the warm sections, NEG coated beam pipes are installed, that have a lower secondary electron yield, and provide linear pumping. In the cold regions, additional pumps are installed to improve the vacuum to an average value of 10^{-5} Torr before the cool-down starts. With the PHENIX and STAR detector upgrades, the vacuum system in the experimental regions will also be improved.

Time in store can be gained through faster machine set-up, a reduction in system failures, and the injection of multiple bunches in each AGS cycle. We project that the time in store can be increased to about 100 hours per week, or 60% of calendar time.

3.3 Polarimetry

Beam polarization measurements in RHIC provide immediate information for performance monitoring, and absolute polarization to normalize the experimental asymmetry results. Two types of polarimeters are used. Both are based on small angle elastic scattering, where the sensitivity to the proton beam polarization comes from the interference between the electromagnetic spin-flip amplitude that generates the proton anomalous magnetic moment and the hadronic spin non-flip amplitude, and possibly a hadronic spin-flip term.

One type of polarimeter uses a micro-ribbon carbon target, and provides fast relative polarization measurements. The other type uses a polarized atomic hydrogen gas target, and provides

slow absolute polarization measurements. In addition, both PHENIX and STAR have developed local polarimeters that measure the residual transverse polarization at their interaction points. These polarimeters are used to tune and monitor the spin rotators that provide longitudinal polarization for the experiments. They polarimeters are discussed in the Experiments section.

The fast proton-carbon polarimeter was first developed at the AGS [157]. It measures the polarization in RHIC to $\Delta P = \pm 0.02$ in 30 seconds. Measurements taken during a typical store in 2004 are shown in Fig. 33. A carbon ribbon target is introduced into the beam, and the left-right scattering asymmetry of recoil carbon ions is observed with silicon detectors inside the vacuum. The silicon detectors observe the energy and time of flight of the recoil particles near 90° [158]. The detector selects carbon ions with a momentum transfer in the coulomb-nuclear interference (CNI) region, $-t = 0.005 - 0.02 \text{ (GeV/c)}^2$. In this region, the interference of the electromagnetic spin flip amplitude and the hadronic non-flip amplitude produces a calculable t -dependent asymmetry of 0.03 to 0.02. The cross section is large, so that the sensitivity to polarization is large. A term from a hadronic spin flip amplitude is also possible and is reported in Ref. [157]. This contribution is not calculable, so that this polarimeter must be calibrated using a beam of a known polarization.

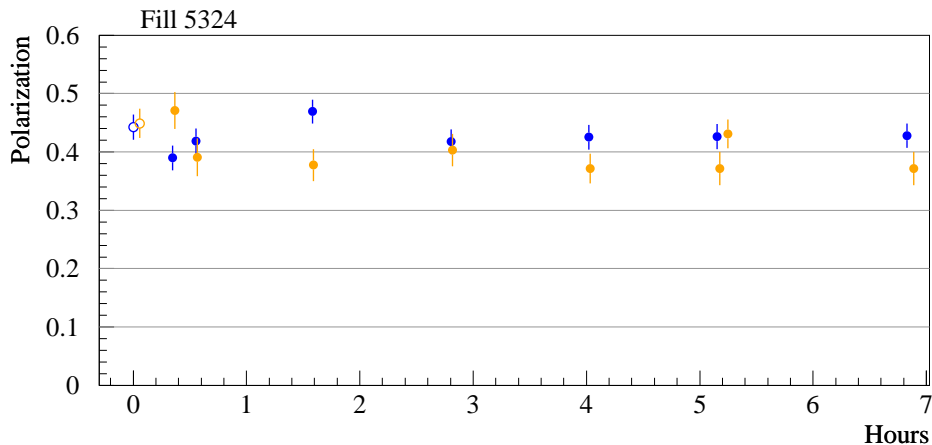


Figure 33: Measured polarization during one store of RHIC in 2004.

A polarized atomic hydrogen gas jet target was used for the first time in RHIC in 2004 [159]. The atoms are polarized with the Stern-Gehrlach process to select one electron polarization state, with rf transition to select proton polarization. The atoms are focused in the RHIC beam region to 6 mm FWHM using the atomic hydrogen magnetic moment. A Breit-Rabi polarimeter after the RHIC beam measures the polarization by cycling through rf transition states. The polarization was determined to be 0.92 ± 0.02 , including correction for the measured 2% molecular fraction (4% nuclear fraction) that is unpolarized. The online target polarization measurements are shown in Fig. ???. The target polarization was reversed roughly every 8 minutes by changing rf transitions. Silicon detectors measure a left-right asymmetry for proton-proton elastic scattering in the CNI region, similar to the p-carbon polarimeters. By measuring the asymmetry with respect to the target polarization sign, we measure the analyzing power for proton-proton elastic scattering, as shown in the elastic scattering subsection. By then measuring the left-right asymmetry with respect to the beam polarization sign, flipping each bunch (every 200 ns), we obtain the absolute beam polarization. The absolute beam polarization was measured to better than $\Delta P/P = 7\%$ in 2004 (preliminary).

A remaining issue is whether the carbon polarimeter calibration can be used for different

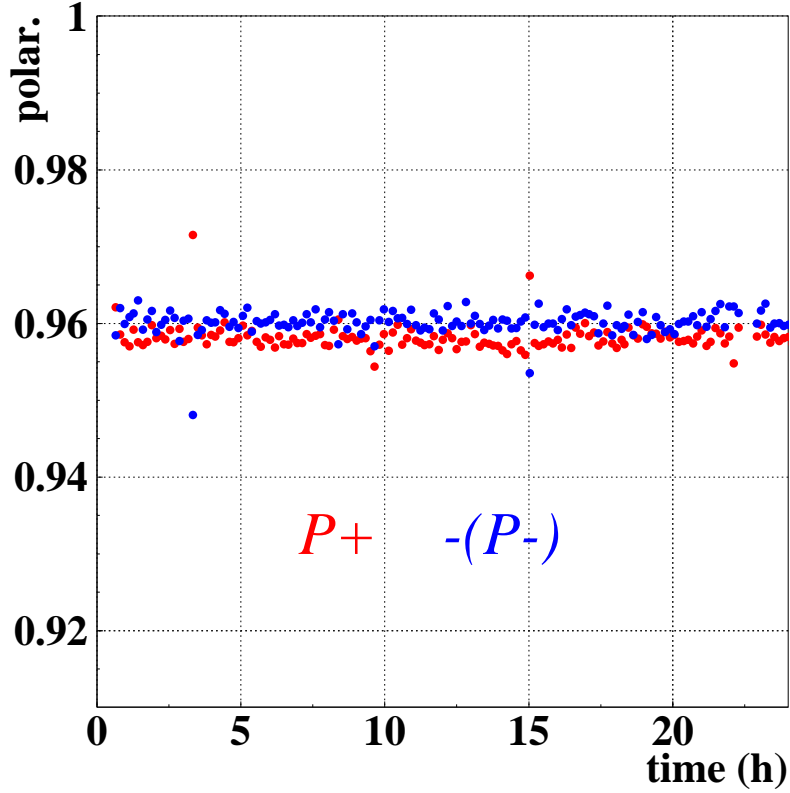


Figure 34: Online polarization measurements for the polarized atomic hydrogen jet target in RHIC during the 2004 run. No correction is included for molecular hydrogen contamination (about 3%).

detectors, from year to year, or whether it will be necessary to recalibrate each year using the jet target. We can also choose to use the jet target as the RHIC polarimeter, with the carbon polarimeter used for corrections, for example for different polarization of the bunches and for a polarization profile of the beams. It will also be necessary to improve the lifetime of the silicon detectors from radiation to avoid changing detectors mid-run, which worsens the RHIC vacuum and is not expected to be compatible with high luminosity running. A related issue is development to be able to bake out the polarimeter region.

3.4 Long-term perspective

A number of ideas are pursued for long-term improvements of the machine performance. RHIC II aims at increasing the heavy ion luminosity by an order of magnitude through electron cooling. For protons, cooling at store is not practical but pre-cooling at injection might be beneficial. A further reduction of β^* , especially at 250 GeV proton energy appears possible. Some benefits may also come from stochastic cooling, currently developed for heavy ions. We expect a luminosity improvement of a factor 2-5 for polarized protons for RHIC II.

With a new interaction region design, the final focusing quadrupoles can be moved closer to the interaction point, thus allowing to squeeze β^* further. This, however, makes some space unavailable for the detectors. Additional increases in the luminosity may come from a further increase in the number of bunches, to close to 360, as is planned for eRHIC, or operation with very long bunches. The latter requires a substantial R&D effort, as well as a new timing system for the detectors.

4 Experiments (Brahms, PHENIX, and STAR)

4.1 PHENIX

RHIC has made great strides towards providing high luminosity beams of highly polarized protons. To make statistically sensitive asymmetry measurements with low systematics requires well understood detectors; clean, highly selective triggers, reliable measurements of beam luminosity and polarization, and the ability to take and analyze data at high rates. In this section we discuss the current and proposed capabilities of the PHENIX detector in the context of meeting the challenges of the spin program.

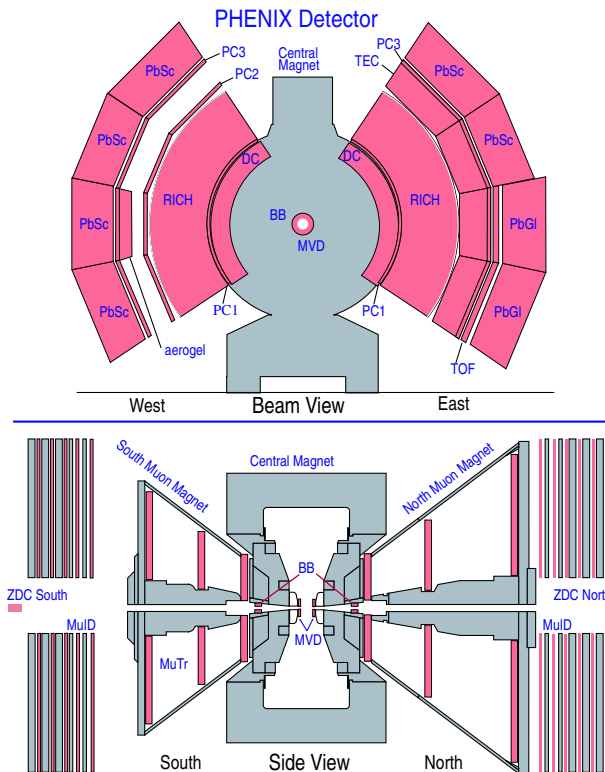


Figure 35: Plan view and side views of the PHENIX detector.

As shown in Fig. 35, the PHENIX detector comprises four instrumented spectrometers (arms) and three global detectors. The east and west central arms are located at central rapidity and instrumented to detect electrons, photons, and charged hadrons. The north and south forward

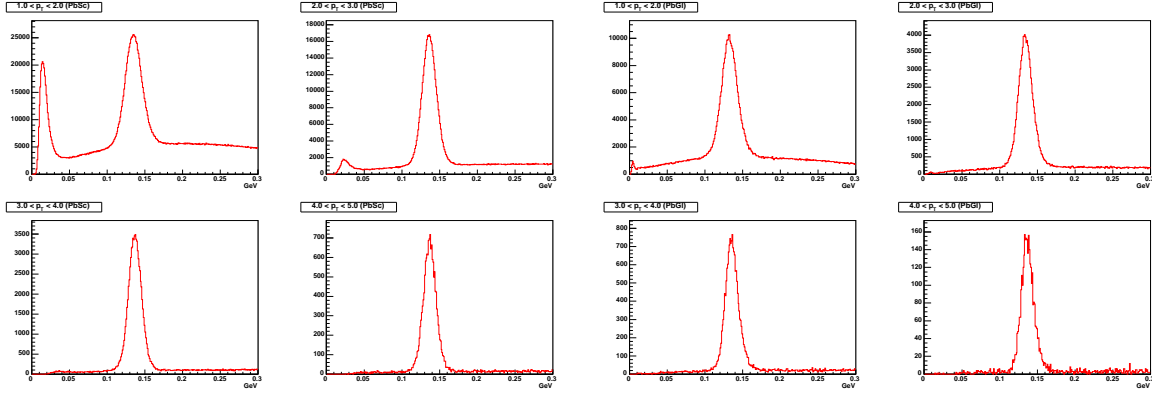


Figure 36: The π^0 mass spectrum in p_T bins 1-2 GeV/c, 2-3 GeV/c, 3-4 GeV/c and 4-5 GeV/c. Events are collected using the high p_T photon trigger of the EMCal. The combinatorial background is $\sim 30\%$ for a p_T range from 1 to 2 GeV/c, and $\sim 5\%$ for $p_T > 5$ GeV/c. The two-photon invariant mass resolution is 8.5% in the first p_T bin and 6.4% in the last p_T bin.

arms have full azimuthal coverage to detect muons. In addition, the zero degree calorimeters (ZDCs) and beam-beam counters (BBCs) measure the time and position of the collision vertex.

4.1.1 PHENIX Central Arms

The PHENIX central arms consist of tracking systems for charged particles and electromagnetic calorimetry. We require a calorimeter with the ability to distinguish isolated photons from those from π^0 decays over a large p_T range. A thorough understanding of the calorimeter and associated triggers is vital for these measurements.

The calorimeter (EMCal) is the outermost subsystem of the central arms, located at a radial distance of ~ 5 m from the beam line. Each arm covers a pseudorapidity range of $|\eta| < 0.35$ and an azimuthal angle interval of $\Delta\phi \approx 90^\circ$, and is divided into sectors containing a lead scintillator (PbSc) calorimeter or lead glass (PbGl) calorimeter. Each calorimeter tower subtends a solid angle $\Delta\phi \times \Delta\eta \sim 0.01 \times 0.01$, ensuring the two photons from π^0 decay are resolved up to a p_T of 12 GeV/c. Shower profile analysis can extend this p_T range beyond 20 GeV/c. The energy calibration used the position of the two photon invariant mass peak from π^0 decay, the energy deposit from minimum ionizing charged particles traversing the EMCal (PbSc), and the momentum determined by the tracking detectors of electrons and positrons identified by the ring-imaging Čerenkov detector. These studies show that the accuracy of the energy measurement was within 1.5%. At a p_T of ~ 11 GeV/c, this uncertainty translates into a systematic error on the π^0 cross-section of $\sim 12\%$. The effective energy resolution was deduced by comparing the measured energy and momentum for identified electrons and positrons and from the widths of the π^0 invariant mass peaks as shown in Fig.36. Their widths varied with p_T from 7% to 10% (PbSc) and 12% to 13% (PbGl).

The number of recorded high- p_T π^0 's is enhanced by a high- p_T trigger which uses threshold discrimination applied to sums of the analog signals from 2×2 groupings (tiles) of adjacent EM-Cal towers. A plot of the trigger efficiency as a function of deposited energy is shown in Fig. 38.

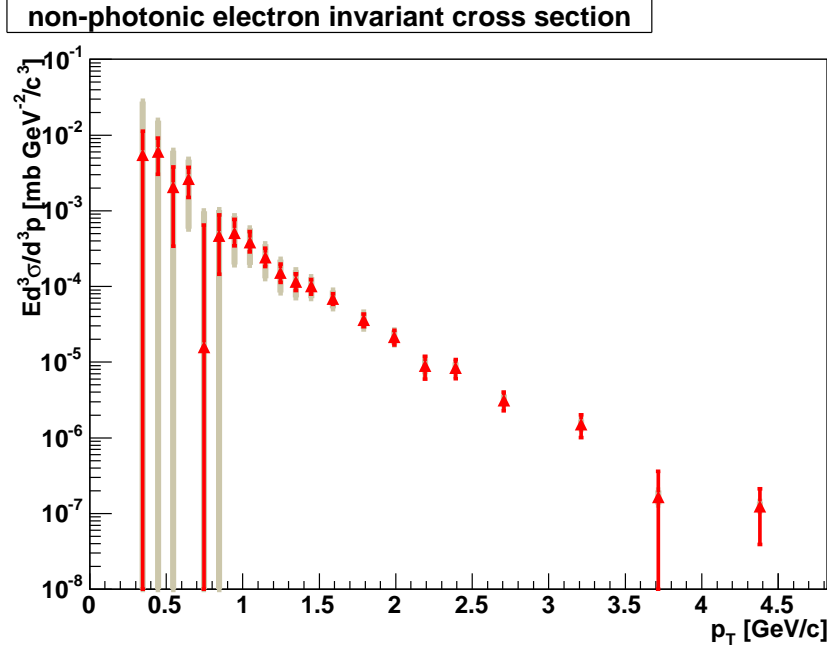


Figure 37: The decomposed nonphotonic spectrum extracted using the converter method with statistical uncertainties corresponding to $\sim 20 \text{ pb}^{-1}$ and one central arm.

The efficiency reached a plateau of 0.78 ± 0.03 at $\sim 3 \text{ GeV}$, which is consistent with the geometrical acceptance of the active trigger tiles, and was reproduced by Monte Carlo calculations. Charged particle contamination in the photon sample was minimized by using information from the PHENIX ring-imaging Čerenkov and tracking detectors [?].

The calorimeter and trigger performance have enabled PHENIX to make many significant measurements within the first few years of running. Measurements have been made of the cross-section and double-helicity asymmetry for π^0 production (see Figures. ??, ??). The prompt photon production cross section in pp collisions has also been measured ([168], [169]) and the NLO pQCD calculation is in good agreement with the data ?? In Ref. [169], a photon isolation cut was applied as a first step towards a spin asymmetry measurement. The cut reduces the level of background photons diluting the analyzing power. With increased luminosity we expect improved precision on these measurements, and the first measurements of the double spin asymmetry $A_{LL}^{pp \rightarrow \gamma X}$.

The EMCal will also be used for measurements of inclusive electron asymmetries from semi-leptonic decays of charm and beauty produced mainly by gluon-gluon fusion in pp scattering. Electrons in the central arms are identified by the RICH detector (Čerenkov threshold for $\pi^\pm \approx 4.9 \text{ GeV}/c$) and the EMCal. The yield of electrons can be categorized into nonphotonic electrons mainly from semi-leptonic decays of charm and beauty (see Fig.37), and photonic electrons mainly from gamma conversion and Dalitz decays of neutral mesons such as π^0 and η [162].

The transverse single-spin asymmetries of neutral pions and non-identified charged hadrons, shown in Fig.21 (page 28 of the combined report), have been measured at mid-rapidity ($x_F \approx 0$). The asymmetries seen in this previously unexplored kinematic region are consistent with zero within statistical errors of a few percent. The dominant partonic processes in both π^0 and

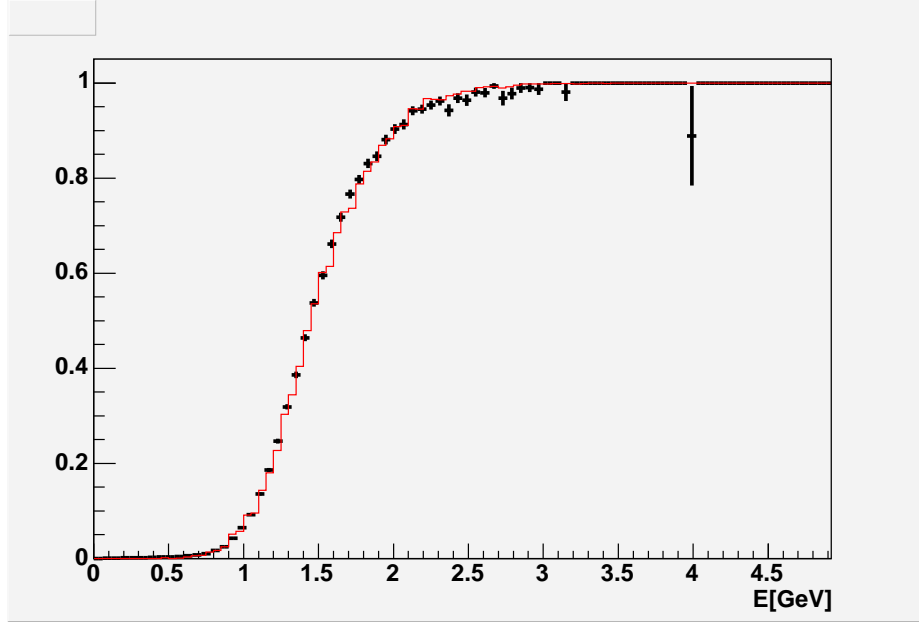


Figure 38: The normalized efficiency of the high- p_T trigger a function of energy. The red line shows the predictions of a Monte Carlo simulation.

charged hadron production in the current kinematic region are gluon-gluon and gluon-quark scattering. This gluon dominance suggests that any potential contribution from mechanisms involving transversity [161] should be significantly suppressed, and the current measurement is probing gluon-sensitive models such as that given in [?, ?].

Ref. to Fig.4, page 13 of the combined report.

4.1.2 Muon Arms

The systematic study of J/ψ production at Relativistic Heavy Ion Collider (RHIC) energies with wide p_T and rapidity coverage should provide crucial tests of J/ψ production models. In addition, the RHIC proton-proton results provide a baseline for studying cold and hot nuclear matter in proton-nucleus and nucleus-nucleus collisions using J/ψ yields as a probe. PHENIX has two forward spectrometers devoted to the characterization of single and di-muon events in the forward rapidity regions. The central magnet poles act as a hadron absorber in front of a radial field magnet with acceptance from $1.2 < \eta < 2.2$ (2.4) in the South (North) spectrometer. Inside the magnetic field, there are three high resolution cathode strip tracking chambers capable of determining space point position to < 100 microns. Downstream of each spectrometer magnet is a Muon Identifier (MuId) which covers the same rapidity region. They consist of five layers of steel absorber sandwiching both horizontal and vertical proportional tubes, with the total thickness of the absorber material of 60cm. The minimum muon momentum able to penetrate all 5 gaps is 2.7 GeV/c, and the pion rejection factor at 3 GeV/c is 2.5×10^{-3} . The MuId is also used as trigger counter as well as identifying the muon. It uses full hit information of the detector, 9 cm in horizontal and vertical direction for each gap, and determines whether a muon candidate road exists for each beam crossing. Triggers on single and double tracks with sufficient depth in the MuId

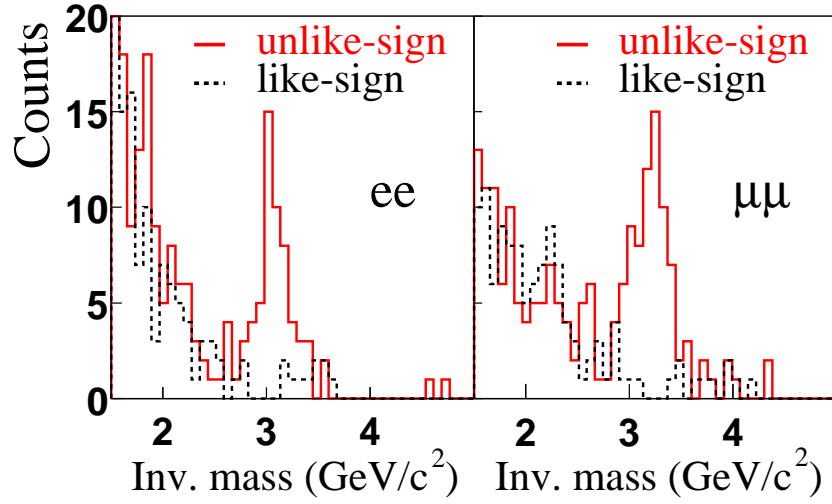


Figure 39:

are passed to the global level-1 trigger of PHENIX. J/ψ yields in the muon arm were obtained by reconstructing $\mu^+ - \mu^-$ pairs. Muon tracks were reconstructed by finding a track seed in the MuID and matching it to clusters of hits in each of the three MuTr stations. The momentum was determined by fitting, with a correction for energy loss, the MuID and MuTr hit positions and the vertex position. J/ψ mass resolution of $160 \text{ MeV}/c^2$ in $p - p$ collisions has been achieved as shown in Fig.39. Together with the di-electron measurement in the central arm, PHENIX has published the p_T , rapidity and total cross-sections for J/ψ at $\sqrt{s} = 200 \text{ GeV}/c$ shown in Fig.40.

4.1.3 PHENIX Local Polarimetry and Relative Luminosity Detectors

Local polarimeters, sensitive to the transverse polarization at collision, were used to set up the spin rotators, and to monitor the beam polarization direction at the PHENIX experiment. The local polarimeters utilized a transverse single spin asymmetry in neutron production in $p-p$ collisions at $\sqrt{s} = 200 \text{ GeV}$ [163]. For vertically polarized beam a left-right asymmetry is observed for neutrons produced at very forward angles, with no asymmetry for production at very backward angles. A fully longitudinally polarized beam produces no asymmetry.

Neutrons with $E_n > 20 \text{ GeV}$ and production angle $0.3 < \theta_n < 2.5 \text{ mrad}$ were observed by two hadronic calorimeters located ± 18 meters from the interaction point (ZDC or Zero Degree Calorimeter [164]). Scintillator hodoscopes at 1.7 interaction length provided the neutron position at the ZDC, and thus the neutron production angle and azimuthal angle $\phi = \arctan(x/y)$ with \hat{y} vertically upward. The \hat{x} axis forms a right handed coordinate system with the \hat{z} axis defined by the beam direction for forward production. Figure 41 shows the observed asymmetry, for the spin rotators off and on, for the blue and yellow beams. With the spin rotators off, a left-right asymmetry is observed from the vertically polarized beam. With the spin rotators on, the measured transverse polarization, averaged over the run, was $\langle P_{Bx} \rangle = 0.033 \pm 0.019$, $\langle P_{By} \rangle = 0.008 \pm 0.020$, $\langle P_{Yx} \rangle = -0.020 \pm 0.013$ and $\langle P_{Yy} \rangle = 0.054 \pm 0.017$, out of $\langle P \rangle = 0.27$. The double spin transverse polarization was $\langle P_{Bx} P_{Yx} \rangle = (0.4 \pm 1.1) \cdot 10^{-3}$ and $\langle P_{By} P_{Yy} \rangle =$

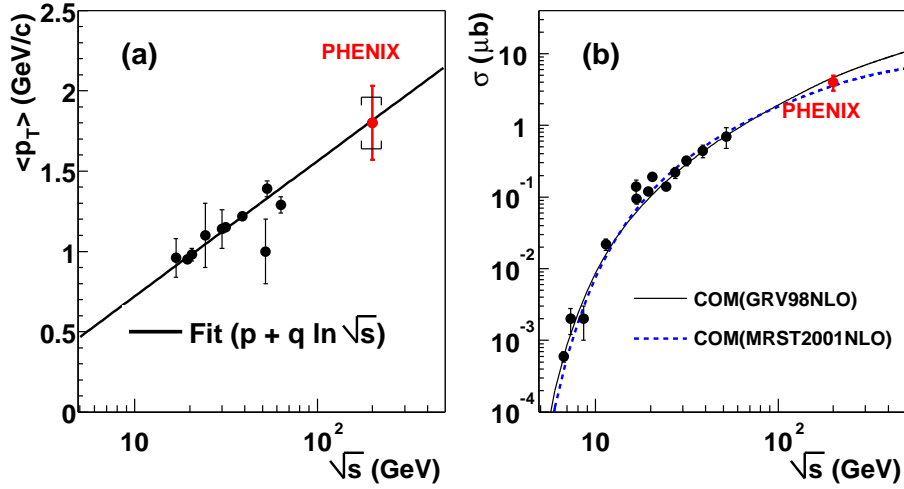


Figure 40:

$(-0.2 \pm 0.8) \cdot 10^{-3}$, compared to $\langle P_B P_Y \rangle = 0.07$. Therefore, with the spin rotators on, the transverse asymmetry is greatly reduced, indicating a high degree of longitudinal polarization: the longitudinal fraction of the beam polarization was 0.99 and 0.98 for the blue and yellow beams, respectively. A separate run with the spin rotators set to give radial polarization confirmed the direction of the polarization for each beam.

In addition to the polarization of the beams, we require information on the intensities and possibly the profiles of the colliding bunches. This knowledge is necessary because the spin structure information we seek is manifested in correlations between the rate or angular distribution of specific final states and the spin direction(s) of the colliding protons. However, spin correlated differences in the production rate of particular final states may appear even in the absence of a physics asymmetry simply because the luminosities of the colliding bunches are different. This necessitates measurements of the relative luminosities of the colliding bunches which are insensitive to the beam polarization to greater precision than the smallest physics asymmetry to be measured. PHENIX is currently equipped with two forward detectors, the BBCs and ZDCs [164, 165], which are primarily sensitive to the pp inelastic and double-diffractive cross-section respectively. They have demonstrated an insensitivity to the colliding bunch polarization at the level of 2×10^{-3} and have a high rejection of backgrounds. Scalars attached to these detectors which can count a single event per crossing form the successful foundation of the relative luminosity measurements suitable for the current pp luminosity. As RHIC approaches its design luminosity, $2 \times 10^{32} \text{ cm}^{-2}\text{s}^{-1}$, there will be multiple pp interactions per bunch crossing, and ambiguities in determining whether the event vertices lay within the PHENIX acceptance. We believe these complications may be overcome by incorporating additional information from the detectors which is linear in the rate such as phototube charge and multiplicity. Bunch profile information from existing RHIC instrumentation, new analysis techniques and trigger logic will be used to minimize the effects of vertex ambiguities on the determination of relative luminosity. Also, a dedicated low cost, small acceptance detector is being considered to address these issues.

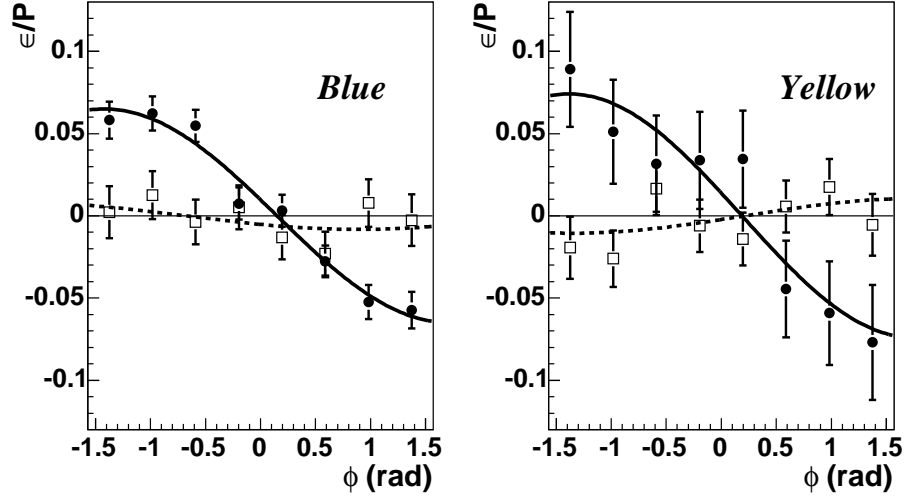


Figure 41: The raw asymmetry normalized by the beam polarization, ϵ/P , as a function of azimuthal angle ϕ , for forward neutron production. The solid points and curve correspond to the spin rotators off (transverse polarization) and the open points and dashed curve correspond to the spin rotators on (longitudinal polarization). Curves are sine function fits to the data, representing possible transverse polarization. The data are for special runs used to set up the spin rotators, where the blue (yellow) polarization was 0.24 and 0.33 (0.08 and 0.28), for spin rotators off and on, correspondingly.

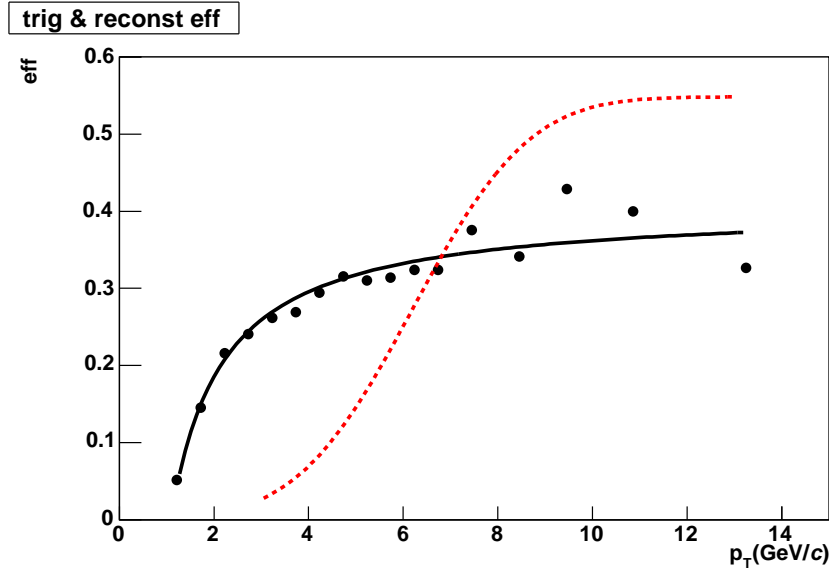


Figure 42: Experimental efficiency for π^0 detection at PHENIX. It includes BBC and EMCal trigger efficiencies and offline data-selection and reconstruction efficiencies. Black points are efficiencies calculated from run2-pp data and the black line is a smoothed curve for eye guidance. It saturates around 35-40%. The dotted red line shows expected efficiency in the future. We can achieve this by removing the BBC coincidence from the trigger and set the EMCal trigger threshold energy higher.

4.1.4 PHENIX DAQ and Computing

The PHENIX DAQ has been designed from inception as a parallel pipelined buffered system capable of very high rates of nearly deadtime-less data-taking. In this kind of design, data is sent from each detector element in multiple parallel streams, buffered at each stage in the chain to smooth out fluctuations in rate, and then uses simultaneous read and write for the highest possible throughput with existing technology. At present, PHENIX has achieved data rates of over 4 kHz, and with improvements in noise reduction PHENIX should be able to approach the peak design interaction trigger rate of 12.5 kHz. Such high DAQ rates are crucial for providing the capability and flexibility to record the many different kinds of interesting rare events at RHIC.

In addition to the RHIC Computing Facility (RCF) at BNL, PHENIX has regional computing centers in the world. The biggest one is the RIKEN CC-J, Computing Center in Japan. The CC-J has comparable computing resources to the PHENIX part of the RCF, CPU power and data storage capability with the High Performance Storage System (HPSS). The main missions of the CC-J for the PHENIX experiment are the primary simulation center, an Asian regional computing center, and a computing center for the spin physics. In run5, it will be used for the reconstructed data production of all the polarized proton collision data. Raw data are sent from PHENIX to the CC-J as much as possible with using WAN connection in parallel to the RCF HPSS. The sustained transfer rate of WAN connection between BNL and RIKEN was confirmed about 10 MB/s and recently is reaching up to 60 MB/s while sustained PHENIX DAQ rate is expected to be 60 - 100 MB/s. Remaining data are sent after the DAQ period through the WAN, and with tape transfer on plains as a backup plan. A bottleneck of the WAN was upgraded very recently, so the transfer rate may reach up to 100 MB/s.

5 Detector upgrades

physics goals and table of physics for upgrades to be inserted ...

5.0.5 Muon trigger

The flavor separation of quark and anti-quark polarizations for up- and down-quarks, requires separate high statistics measurements of inclusive lepton counting rate asymmetries: $A_L^{W^+ \rightarrow \mu^+}(p_T)$ and $A_L^{W^- \rightarrow \mu^-}(p_T)$. These measurements translate into the following experimental requirements for the PHENIX muon arms: (a) superior event selection capability in order to reduce the 10MHz collision rate to the data archiving bandwidth available in PHENIX (b) the ability to assign the correct proton polarization (that is bunch crossing number) to a given W-event candidate, (c) tracking resolution to correctly determine the lepton charge sign and (d) good signal to background ratios in the off-line analysis.

Extensive Monte Carlo simulations including a full GEANT simulation of the muon arms show that the existing muon spectrometers are capable of defining a clean sample of W events for the off-line analysis: a requirement of $p_T > 25$ GeV on the transverse momentum of the final state decay muon will remove most of the collision and beam related backgrounds; we expect a

signal to background ratio of about 2:1.

A new first level muon trigger is required to improve the online performance of the present first level muon trigger. Rejection factors achieved in the present PHENIX muon trigger, based on information from the existing muon identifier system are about $R=250$. Measurements at the luminosities needed for the W-physics program will require rejection factors of $R>5000$.

The new first level muon trigger in PHENIX will be based on three fast trigger stations which will be added to each of the PHENIX muon spectrometers. The triggers stations will use Resistive Plate Chamber (RPC) technology developed for the muon trigger in the CMS experiment at LHC. In addition new front end electronics for the existing muon tracker chambers will make it possible to introduce muon arm tracking information in the future muon trigger. Information from the three RPC stations and the muon tracker will be processed in standard PHENIX first level trigger processor boards. The boards carry large Xilinx Fast Programmable Gate Arrays (FPGAs) and can carry out a fast online momentum measurement based on the tracking information from the RPCs and the muon trackers. In addition a timing cut will be applied in the RPC front-end electronics to remove beam backgrounds.

Detailed beam background measurements and beam-loss Monte Carlo simulations were carried out to confirm that the substantial steel absorbers in PHENIX introduce a large asymmetry in the incoming and outgoing beam backgrounds which can be exploited for triggering purposes. Overall the expected rejection power of the new muon trigger is expected to be well above the required level of $R>10000$.

5.0.6 Silicon VTX

PHENIX Collaboration proposes to construct a Silicon Vertex Tracker (VTX) in the next few years. The VTX will substantially enhance the physics capabilities of the PHENIX central arm spectrometer. Our prime motivation is to provide precision measurements of heavy-quark production (charm and beauty) in A+A and p(d)+A collisions, and polarized p+p collisions. In addition, addition of a large acceptance central detector capable of monitoring multiplicities of hadronic final states along with their directions, will enhance PHENIX's capability to study Jets in hadron-hadron collisions. These are key measurements for future RHIC program, both for heavy ion physics which intends to study the properties of dense nuclear medium created in their collisions, and for the exploration of the nucleon spin-structure through polarized pp collisions. While the detailed list of physics measurements possible with the VTX detector are discussed elsewhere [1], the principle measurements associated with polarized pp program are:

- $\Delta G/G$ from charm and beauty production in polarized pp scattering
- x dependence of $\Delta G/G$ from γ -Jet correlations

The PHENIX detector has been described before in Section 2.2. The heavy quark production has been measured by PHENIX only indirectly through the observation of inclusive (decay) electrons. These measurements are limited in accuracy by the systematic uncertainties resulting from possible large electron backgrounds originating from Dalitz Decays and photon conversions. The measurements are statistical in nature, and one uses different models to distinguish between charm

and beauty contributions. The VTX detector will provide tracking with a resolution of $< 50\mu\text{m}$ over a large coverage both in rapidity $|\eta| < 1.2$ and in azimuthal angle ($\Delta\phi \sim 2\pi$). A significantly improved measurement of heavy quarks in pp collisions is deemed possible over a wide kinematic range with the VTX.

The proposed VTX detector will have four tracking layers. For the inner two layers we propose to use silicon pixel devices with $50 \times 425\ \mu\text{m}$ channels that were developed for the ALICE experiment at CERN/LHC. Our preferred technology for outer two detector layers is a silicon strip detector developed by the BNL instrumentation division which consists of $80\mu\text{m} \times 3\text{ cm}$ strips layed to achieve an effective pixel size of $80 \times 1000\ \mu\text{m}$. We plan to use SVX4 readout chip developed at FNAL for the strip readout. The main aim in using ready technology has been to shorten the cost and the time for the R&D that for such a project could ordinarily be rather high and long, respectively.

PHENIX anticipates that the project will be mainly funded by two agencies: the DOE Office of Nuclear Physics and RIKEN Institute of Japan. While RIKEN funding has been available since 2002, it is anticipated that the construction funds for the Strip Layers will be available from DOE starting FY06 and the project will get constructed and commissioned in the RHIC - Run 8 which presently is expected to be a long Au+Au run.

5.0.7 Nosecone calorimeter

The PHENIX detector at RHIC has been designed to study hadronic and leptonic signatures of the new states of matter in heavy ion collisions and polarized proton collisions. The baseline detector measures muons in two muon spectrometers located forward and backward of mid-rapidity, and measures hadrons, electrons, and photons in two central spectrometer arms each of which covers 90 degrees in azimuth and ± 0.35 unit of rapidity where the existing electromagnetic calorimeters are installed. Further progress requires extending rapidity coverage for hadronic and electromagnetic signatures beyond the limits set by already built central spectrometer, in particular upgrading the functionality of the PHENIX forward spectrometers to include photon and jet measurement capabilities.

A forward spectrometer upgrade is being proposed with the objective of introducing new detector capabilities and greatly enhancing present capabilities for PHENIX in the forward direction. When completed it will result in a nearly ten-fold increase to its rapidity coverage for photons and to some extent for hadrons, jet detection and better triggering capabilities. Newly acquired access to forward production of inclusive jets, direct photons or Drell-Yan pairs at large x_F in nucleon-ion collisions at RHIC will provide a new window for the observation of saturation phenomena expected at high parton number densities and its importance in the evolution of the partonic distribution functions.

The PHENIX Forward Upgrade is constrained by the existing muon spectrometer configuration including its wire chambers, hadron absorber walls and magnet yokes. The core element of the proposed upgrade are compact tungsten calorimeters with silicon pixel readout and fine transverse and longitudinal segmentation built to identify and measure forward electromagnetic activity and provide jet identification and coarse jet energy measurements. The principal performance aspect of the NCC is its ability to run in the unassisted mode (without upstream tracking).

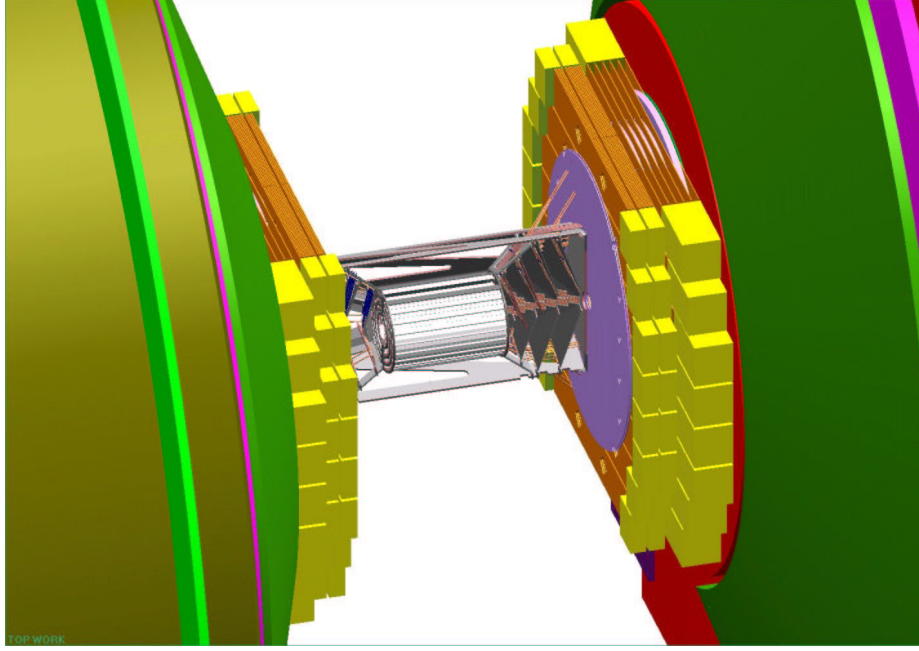


Figure 43: Longitudinal structure of a single NCC tower.

The NCC is extremely dense sampling calorimeter using tungsten absorber interleaved with silicon readout layers. Geometrical constraints of the PHENIX detector call for a very nontraditional and challenging design. Tungsten has a very short radiation length and a large absorption/radiation length ratio. The electromagnetic shower develops in the very first few cm of the calorimeter depth thus allowing the implementation of both electromagnetic and shallow hadronic compartments within the available 20 cm of space. To further suppress the hadronic contribution to the energy seen in the electromagnetic compartment the electromagnetic calorimeter will be only ~ 10 radiation length deep. The tungsten in the first 16 layers has a thickness of 2.5 mm ($\sim 0.7 X_0$), and 1.6 cm in the remaining 6 layers serving as hadron rejecter - shower tail catcher (Fig.43).

Limiting the depth of the electromagnetic compartment to ~ 10 Lrad serves to improve the em/hadron shower discrimination but will also result in the loss of resolution for high energy showers. To minimize the potential negative impact of this decision on the performance of the energy measurements the same lateral granularity will be retained in both electromagnetic and hadronic compartments ($1.5 \times 1.5 \text{ cm}^2$ pixels). Measurements in both compartments will combine to recover the resolution.

plan for high luminosity (relative luminosity, trigger rates, bandwidth, backgrounds, crossing ambiguities, vertex, ... to be inserted ...

text for funding issue to be inserted ...

5.0.8 Projection for future measurements

In the PHENIX, single electrons from semi-leptonic decays from charm and beauty meson are measured in mid-rapidity region. These 2 sources will be decomposed by VTX upgrade (looking at DCA distribution). This picture shows the ALL error of electrons from charm and beauty. Each points shows ALL error of the pT region above pTmin.

Ref. to Fig.8, page 19 of the combined report.

Ref.[167]

5.1 STAR

5.1.1 Recent spin-related upgrades

A cross-sectional view of the STAR detector, emphasizing the subsystems that have been added to the baseline detector [?] with the spin program as a primary driver, is shown in Fig. ???. The relative luminosity monitoring critical for the measurement of spin asymmetries is provided by Beam-Beam Counters (BBC) that have been added just to the east and west of the STAR magnet, at a distance of 3.5 m from the beam intersection point. The BBC's are plastic scintillating tile detectors for charged particles over the pseudorapidity range $3.3 < \eta < 5.0$. An east-west prompt coincidence, sensitive to nearly 90% of the total non-singly diffractive pp cross section at $\sqrt{s} = 200$ GeV, is used to discriminate beam collisions from beam-gas interaction background. The azimuthal segmentation of the scintillating tiles permits the measurement of left-right and up-down single-spin asymmetries. Comparison of the BBC asymmetries measured at STAR for hits in the inner BBC tiles with those measured simultaneously in the RHIC CNI polarimeters has revealed a small analyzing power ($A_N \approx 0.006$) suitable for use of the BBC's as a local polarimeter when the beam spin at STAR is transverse to its momentum. In particular, this functionality is important for tuning the STAR spin rotators: as shown in Fig. ??, one can adjust the rotator magnet currents to give longitudinal polarization at the IR by arranging for both (left-right and up-down) BBC transverse asymmetries to vanish.

The major STAR upgrades already installed for the spin program represent additions of electromagnetic calorimetry (EMC) for the detection of high-energy photons, electrons and π^0 over a broad range of pseudorapidity. The Forward Pion Detectors (FPD) are small lead-glass-based calorimeters placed to the left and right, and above and below, the beam lines 7.5 m to the east and west of the center of STAR. The FPD provides π^0 detection and identification at high $x_{Feynman}$ and forward rapidity ($3 < \eta < 4$), where large single-spin transverse asymmetries have been observed (see Fig. ??). (The measurements reported in [?] were made with a precursor of the present FPD's that was a prototype section of the STAR Endcap EMC.) The FPD's will continue to be used to probe the origins of these large single-spin effects, and also to investigate gluon polarization via di-hadron coincidence measurements of the type discussed in Sec. ??. In addition, the FPD provides access in d+Au collisions at STAR to the low Bjorken- x regime where gluon saturation models predict, and the RHIC experiments have observed, the onset of significant suppression of moderate- p_T hadron production.

The largest additions to STAR relevant to the spin program are the Barrel (BEMC [?]) and Endcap (EEMC [?]) calorimeters, funded by DOE and NSF, respectively. Each of these subsystems is a multi-layer sandwich of Pb radiator sheets and plastic scintillator with light collection via optical fibers. Each contains a fine-grained Shower-Maximum Detector (SMD – gaseous for BEMC, plastic scintillator for EEMC) for discriminating single photons from π^0 daughter photon pairs, by means of the transverse shape of the electromagnetic showers produced. Each also contains pre-shower layers, and the endcap adds a post-shower layer as well, to improve electron/hadron discrimination. As shown in Figs. ?? and ??, the fabrication and installation of both EMC's has been completed during 2004, although installation of final readout electronics for the east half of the BEMC is still anticipated to occur during February-March 2005. The EMC's provide critical detection and triggering capability for STAR studies of jets, photons, π^0 , W -bosons and J/ψ (as well as heavier quarkonium species), all of which play significant roles in the spin program described in earlier sections of this document. The broad pseudorapidity coverage ($-1 \leq \eta \leq 2$) permits, for example, study of γ -jet coincidences spanning a broad range of x -values for the participating gluons, while still maintaining large transverse momenta ($p_T > 5$ GeV/c) for the partonic scattering.

Both EMC's have performed well in partial installations for the 2003 and 2004 RHIC runs, enabling detector commissioning, optimization of online calibration and triggering, debugging of subtle electronics problems and initial extraction of physics results. The first paper based on BEMC transverse energy measurements for Au+Au collisions at STAR has already been published [?]; others, based on electron spectra measured with the BEMC, are in preparation. Figure ?? shows an event display for a dijet detected with STAR's TPC and BEMC, together with the spectrum of the ratio of BEMC/total transverse energy for jets reconstructed within the partial BEMC acceptance available during the 2004 p+p RHIC run. Figure ?? shows a typical shower profile measured with the EEMC Shower-Maximum Detector for a π^0 candidate, together with an invariant mass spectrum reconstructed from the SMD and calorimeter tower information for all two-cluster combinations detected with the partial EEMC for several 2004 p+p runs. Analysis groups within STAR are actively working on optimizing the efficiency and background suppression in reconstruction algorithms for jets, π^0 and single photons, electrons and J/ψ . This work is aiming toward first EMC-based spin publications in the latter half of 2005 (based on 2003 and 2004 data) and toward readiness for prompt analysis of results from the anticipated long p+p run during 2005.

In addition to the above hardware upgrades, STAR has enhanced its spin program by the addition of new collaborators over the past few years. Prominent new groups with heavy interest in spin from MIT, LBNL, CalTech and Valparaiso University complement the groups from BNL, ANL, Indiana University, UCLA, Penn State, Texas A&M University, JINR Dubna and IHEP Protvino, who have been instrumental in launching STAR's spin program. With these hardware and manpower additions, the resources are in place to permit STAR to address many of the high priority spin physics goals described in Sec. 2, most prominently the delineation of gluon helicity preferences via a number of reaction channels. Learning curves, but no additional equipment, are anticipated for dealing most effectively with the TPC pileup and BBC occupancy problems that will arise as the p+p luminosity increases at RHIC, and for monitoring beam polarization locally at STAR when the beam spins are oriented longitudinally. However, additional upgrades are needed, as described below, to optimize W detection and the measurement of flavor-dependent sea antiquark polarizations in STAR.

5.1.2 Future STAR upgrades relevant to the spin physics program

The W^\pm production studies central to the envisioned spin physics program at $\sqrt{s} = 500$ GeV strain the capabilities of STAR's Time Projection Chamber, which was designed for heavy-ion collisions to track charged particle momenta up to $p_T \sim 10$ GeV/c. The TPC provides very limited resolution for W daughter leptons up to p_T of 40 GeV/c, especially in the endcap region ($\eta > 1$), where the drifting electrons from charged particle tracks intercept a decreasing fraction of the readout pad rows. Fortunately, the EMC's provide measurements of electron transverse energy with a typical resolution $\sim 4\%$ for $p_T \approx 40$ GeV/c. But the EMC's alone cannot discriminate electrons from positrons, and charge sign determination for the parent W is critical to the physics goal of separating \bar{u} from \bar{d} polarizations in the proton sea (see Sec. ??). Furthermore, comparison of p_T measured from track curvature with E_T measured in the calorimeters provides a powerful method (over and above those available from the EMC's alone) to discriminate the weak W signal from a background of more abundant high- p_T charged hadrons. While STAR's present tracking capabilities are adequate for these tasks in the BEMC region, upgraded tracking is needed for W^\pm production in the endcap region, where the separation of \bar{u} from \bar{d} polarizations is kinematically cleanest (see Sec. ??).

The need for improved endcap tracking is demonstrated in Fig. ?? by simulations of 30 GeV/c pion tracks, whose sagitta in the middle of the endcap region is ≈ 2.5 mm. With the TPC alone, the charge sign is misidentified about 15-20% of the time. The addition of three space points measured with $\sim 50\mu\text{m}$ resolution near the vertex and two with $\sim 100\mu\text{m}$ resolution just in front of the endcap would completely remove the charge misidentification problem, and would provide 30% relative p_T resolution at the high momenta of relevance to the W production program. Improvements of this order are quite feasible with the arrangement of detectors outlined below, and would greatly enhance the physics impact of STAR's W measurements.

The conceptual designs presently under consideration for upgraded forward tracking in STAR are illustrated in Figs. ?? and ??. Space points near the vertex would be provided by an array of annular silicon strip detectors placed just downstream of STAR's vertex tracking devices, and inside the TPC's inner field cage. Position would be measured just upstream of the EEMC in an array of 3-layer Gas Electron Multiplier (GEM) chambers [?] tiling the pseudorapidity region $1 < \eta \leq 2$, over the full azimuthal acceptance. The silicon disks would be integrated with and extend the coverage of a new barrel silicon-strip tracker envisioned to surround, and provide high-rate tracks pointing to, a high-resolution Advanced Pixel Sensor micro-vertex detector (the Heavy Flavor Tracker, HFT). The primary motivation for the latter two subsystems comes from the needs for high-quality measurements of slightly displaced vertices associated with heavy quark production in STAR heavy-ion collisions. Research and development on the type and size of GEM chambers that would be needed for the endcap tracker is ongoing in a collaboration between STAR and PHENIX. In addition to the ability to obtain the needed spatial resolution with thin chambers that would fit within the narrow space available in front of the EEMC, the GEM technology provides fast detectors whose tracking information would significantly alleviate ambiguities from pileup tracks in the TPC anticipated at the ultimate RHIC pp luminosities.

STAR groups from MIT and Lawrence Berkeley National Laboratory are leading a collaboration of several institutions (including ANL, BNL, IUCF, Yale and Zagreb) in planning the inner and endcap tracking upgrades. (The heavy-ion-driven HFT upgrade project is led by the LBNL group.) The conceptual design of the silicon disks and GEM chambers is not yet complete, with

still important open issues to address concerning the optimal tradeoffs between coverage and cost, the resolution impact of material in the endcap of the TPC, etc. The present plan is to proceed with this upgrade in two stages: first, the silicon barrel would be proposed, with the goal of fabricating and installing it in STAR in time for an FY09 RHIC run, when it would be needed to get optimal usage from the HFT (which is to be proposed on a slightly faster timeline); the silicon disks and GEM chambers would be proposed about a year later, with the goal of installation for a long 500 GeV pp run in FY10, when we would hope to collect a large fraction of the statistics needed for the W production A_L measurements. The rough funding scope anticipated for these two phases is \sim \$6M and \$5M, respectively, with contingency. Both DOE and NSF funding are likely to be sought to share the costs for the second phase.

Additional planned STAR upgrades driven by the heavy-ion research program will also have significant benefits for the p+p spin measurements. An extension of STAR's forward electromagnetic calorimetry coverage beyond the EEMC, to $2.5 \lesssim \eta \lesssim 4$, has been proposed to the NSF in January 2005. This Forward Meson Spectrometer (FMS) (see Fig. ??) would be constructed from existing lead-glass counters to replace and expand the present Forward Pion Detector on STAR's west side. The primary motivation for the FMS is to probe the contributions of gluons to nuclear structure at very low Bjorken x (down to $x \sim 0.001$) in proton-nucleus or deuteron-nucleus collisions at RHIC. Gluon saturation models predict a suppression of hadron production at moderate p_T in this forward rapidity region, which can be tested by systematic measurements for mesons reconstructed from their daughter photons. Furthermore, Color Glass Condensate approaches to high-density QCD treat the gluons at such low x and at moderate momentum transfers as a classical field, from which parton scattering will result in mono-jet, rather than traditional di-jet, products. The FMS will allow searches for the onset of such mono-jet events as a function of the rapidity interval between correlated pairs of hadrons (*e.g.*, π^0) detected in coincidence within STAR's extended EMC coverage.

Similar coincidences between leading hadrons in polarized p+p collisions permit study of parton-parton scattering spin sensitivities under conditions where one can vary the subprocess (qq vs. qg vs. gg) contributions in a controlled manner. As described in Sec. 2.4.1 of this document, $A_L L$ measurements for such di-hadron events are part of the planned approach toward unraveling the polarization of gluons in a polarized proton. Furthermore, A_N measurements for mesons in the FMS, in coincidence with other jet fragments, will probe the origin of the large single-spin asymmetries already found (see Fig. ??) for forward π^0 production, as described in Sec. 2.5. Thus, the spin program would make substantial use of the FMS once it is installed in STAR.

It is hoped the FMS will be installed for the FY07 RHIC run, when STAR's present plans call for the next long d+Au collision run. Total funding needed for this project is about \$1M, with most of this to be supplied by NSF. The project is being led by STAR physicists from Penn State University, BNL and IHEP Protvino, with additional participation by LBNL and Texas A&M University. The lead-glass cells would be taken from the existing STAR west-side FPD and from an available supply owned by the Protvino group.

The case for other STAR upgrades driven by the heavy-ion program is summarized in the STAR Decadal Plan [?]. An upgrade to the TPC front-end readout electronics and to the STAR Data Acquisition system will increase the event rate capability from the present ~ 100 Hz to ~ 1000 Hz, by 2007-8. This upgrade has two significant implications for the spin program: the

FEE upgrade will free up space directly in front of the EEMC, needed for eventual installation of the endcap GEM tracker described above; the DAQ upgrade will permit collection of large data samples for abundant reaction channels, such as inclusive jet or π^0 production at moderate p_T , without introducing sizable and undesirable dead time for the rarer channels, such as direct photon production. The barrel Time-Of-Flight detector proposed to improve (by 2008) STAR's particle identification up to ~ 3 GeV/c in heavy-ion collisions will also aid the spin program, *e.g.*, in permitting clean identification of charged pions, and hence of the ρ -meson invariant mass region, for studies of transversity via interference fragmentation functions (see Sec. 2.5). The Heavy Flavor Tracker mentioned above, while again driven by studies of the unique matter produced in heavy-ion collisions, will also permit improved identification in STAR of the production of heavy quarks in p+p collisions, providing access to gluon polarization and to possible spin effects from the quark mass (explicit chiral-symmetry-breaking) terms in the QCD Lagrangian.

In summary, the completion of the barrel and endcap EMC's, following the addition of the Beam-Beam Counters and Forward Pion Detectors, has brought STAR to full readiness to exploit the anticipated long polarized pp collision runs during the 2005-9 period. A significant upgrade to STAR's forward tracking capabilities is still needed to optimize its W^\pm production program in 500 GeV pp runs anticipated for the 2009-12 period. Other STAR upgrades planned for the next several years, primarily to enhance its capabilities in studying heavy-ion collisions, will have significant side benefits for the spin program.

5.2 Other experiments

5.2.1 PP2PP running with the current setup

With a small modification requiring rotation of RP2 and RP4 to horizontal orientation, the present experimental setup, see Fig. ?? is suitable for additional measurements in an extended $|t|$ -range. At $\sqrt{S} = 200$ GeV one can use the capacity of existing power supplies to run with the accelerator optics of $\beta^* = 20$ m. The $\beta^* = 20$ m tune at $\sqrt{S} = 200$ GeV makes it possible to extend the kinematic coverage to a lower $|t|$ of $0.003 < |t| < 0.020$ (GeV/c)². At $\sqrt{S} = 500$ GeV the optics with $\beta^* = 10$ m will be used, allowing measurements up to $|t| \approx 0.12$ (GeV/c)².

The result obtained in Run 2003 is shown in Fig. ?. With the modification of the setup described above one will be able to improve the A_N measurements, measure $A_N N$. This will certainly help resolve the issue of the presence of the hadronic spin flip amplitude at high energies.

5.2.2 Brahms

The BRAHMS detector is well suited to explore the x_f and p_T dependence of the single spin asymmetries for identified charged hadrons. The forward spectrometer is operated at 2.3° and 4.0° for these measurements. The spectrometer has a total bending power of 7.2 Tm and a momentum resolution of $\delta p/p \approx 1 - 2\%$. The particle identification in the current setup using a Ring Imaging Cherenkov (RICH) detector allows for π^\pm identification for momenta up to 35 GeV/c. Operating the RICH at a lower pressure of the radiative gas of C_4F_{10} will allow for π identification up to ≈ 50 GeV/c. Identification of kaons is possible up to a somewhat lower momentum

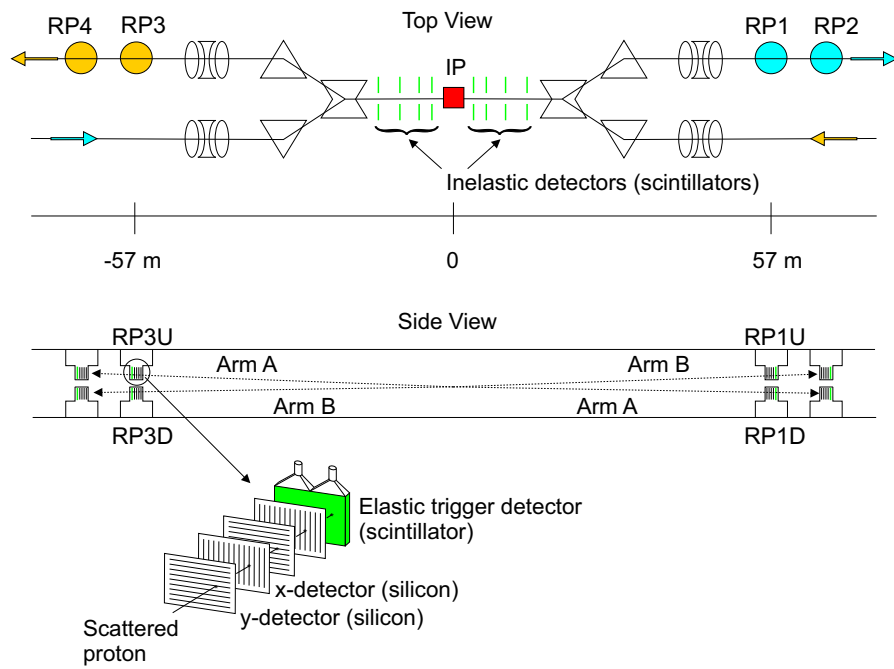


Figure 44: Layout of the pp2pp experiment. Note the detector pairs RP1, RP2 and RP3, RP4 lie in different RHIC rings. Scattering is detected in either one of two arms: Arm A is formed from RP3U and RP1D. Conversely, Arm B is formed from RP3D and RP1U.

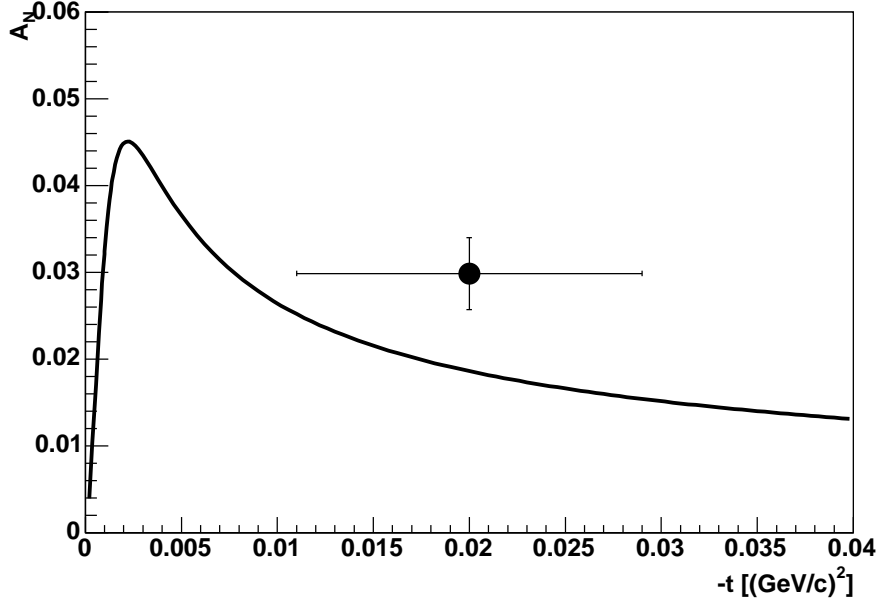


Figure 45: The single spin analyzing power A_N for the full interval. Vertical error bars show statistical errors, horizontal error bars show the $|t|$ -range. Solid curve corresponds to theoretical calculations without hadronic spin flip[?].

than π_S due to the lower yields. During Run-4 a small sample of π^+ and π^- data were taken. Preliminary data are shown in section XX. The coverage in p_T vs x_f is shown in Fig. yy. Future measurements in Run-5 and possible Run-6 will extend this coverage towards larger x_f and p_T . The required delivered luminosity for such measurements are $2 - 4 pb^{-1}$.

5.2.3 Jet Target Experiment

With the addition of recoil detectors at more forward angles, covering 85° to 60° , and with forward detectors near the beam direction, the t range of the polarized atomic hydrogen jet target in RHIC can be substantially increased. The present detectors cover $-t=0.0015$ to 0.03 $(GeV/c)^2$ for polarimetry. A proposal is being considered to reach $-t=1$ $(GeV/c)^2$. The experiment would make very precise measurements of the cross section and transverse spin asymmetries ($d\sigma/dt$, A_N , A_{NN}) for pp elastic scattering, from the Coulomb nuclear interference region to the intermediate t region dominated by Pomeron scattering, and into the diffraction dip region around $-t \sim 1$ $(GeV/c)^2$. With the use of additional RF transitions, the jet target can also produce deuteron beams. Deuterons would allow the study of polarized pd scattering, comparing the protons scattering from protons ($I=1$) to deuterons ($I=0$). Another attractive possibility with this facility is to scatter light ion beams from the polarized jet target. This allows the study of nuclear effects in polarized pA scattering.

5.2.4 Large Acceptance New RHIC Detector

A group of RHIC experimenters has been considering a proposal for a large aperture detector based on ...

6 Spin plan schedule (Gerry)

In the charge, we were requested to consider two running schedules: 10 and 5 physics weeks on spin per year. These follow, showing *example* plans. We emphasize that we expect that the actual run plan will be developed from the experiment beam use proposals. Our consideration of these scenarios should not suggest that we advocate a change to this successful approach.

A key issue is the completion of experiment hardware to run the W physics program. The required hardware are the muon trigger improvements for PHENIX, and a forward tracker for STAR. The PHENIX improvements are being proposed to NSF (\$1.8M for resistive plate chambers) and to the Japan Society for Physical Sciences (\$1.0M for muon tracking readout electronics), with a planned completion for the 2008 RHIC run. The STAR tracker is planned to be proposed to DOE (estimated \$5M) in 2006, and to be complete for the 2010 run.

The example plan below for the 10 physics week/year case is "technically driven". The plan assumes that the funding is received, and the work is completed as planned. For the 5 week plan, the delay in reaching luminosity goals for $\sqrt{s}=200$ GeV delays the start of the W running considerably, by greater than three years. An early completion of the W hardware is less of an issue for this case.

A second key issue is machine performance. We assume that we reach the polarization goal of 70% in 2006. For luminosity, we assume in the example plan that we reach two thirds of the "maximum" luminosity (see section 3). This assumption is discussed there.

A third key issue is experiment availability, in which we include up time, live time, and the fraction of the collision vertex accepted by the experiment. This results in "recorded luminosity" for each experiment. We have taken the up time to be 70% for each experiment, as has been achieved. The live time for PHENIX is 90%, due to multi-event buffering; the live time for STAR is 50%. The online data selection adjusts thresholds, for example the lower p_T requirement, to reach these live time levels. The PHENIX vertex acceptance for the 200 GeV running is 60%, requiring the vertex to be within 20 cm of the IP. We have used this acceptance also for 500 GeV. The STAR vertex acceptance contains all collisions. The overall factor for recorded/delivered luminosity for both experiments is 35%. The physics sensitivities shown in section 2 also include apparatus acceptance and event selection acceptance.

6.1 10 physics weeks

Table 6 shows the example spin plan for 10 physics weeks per year, with a *technically driven* schedule. The 200 GeV running continues through 2008, with a total of 300 pb^{-1} delivered,

and 100 pb^{-1} recorded luminosity by both PHENIX and STAR. By the year 2009, the PHENIX muon triggering improvements are complete, and the STAR forward tracking is partially in place, and complete for the 2010 run. The year 2009 is considered an engineering run, for both the accelerator and the experiments. By the completion of the year 2012, for 500 GeV, 800 pb^{-1} luminosity is delivered, and 300 pb^{-1} recorded by each experiment. These luminosities and polarizations provide the physics sensitivities presented in section 2.

Table 5: RHIC spin example schedule, 10 physics weeks per year, technically driven.

Fiscal year	Spin Weeks	CME(GeV)	P	L(pb^{-1})	Remarks
2002	8	200	0.15		First pol. pp collisions! Transverse spin
2003	10	200	0.27		Spin rotators commissioned, first helicity measurements
2004	1	200	0.4		New betatron tune developed, first jet absolute meas. P
2005	9	200	0.5	10-20	$A_{LL}(\pi^0, jet)$, also 500 GeV studies
2006	10	200	0.7		AGS Cold Snake commissioned, NEG vacuum coating complete
2007	0				
2008	20	200	0.7		Direct γ , completes goal for 200 GeV running
2009	10	500	0.7		PHENIX muon arm trigger installed, eng. run
2010	10	500	0.7		STAR forward tracker installed, W physics
2011	10	500	0.7		
2012	10	500	0.7		Completes 500 GeV goal

6.2 5 physics weeks

Table ?? gives the example spin plan for 5 physics weeks per year, which we have interpreted to mean 10 physics weeks each two years to reduce the end effects. As has been presented in section 3, the delay in the RHIC spin physics results is actually greater than a factor of two, compared to 10 physics weeks each year. This is due to an assumed "turn-on" period of reaching the instantaneous luminosity maximum that is based on our experience, from the heavy ion program. In any case, the programs are stretched out to over 6 years for the gluon polarization measurements at 200 GeV, and an additional 6 years or more for the W physics program. The proposed measurements would be completed in 2018 or later.

Table 5.2 RHIC spin example schedule, 5 physics weeks per year.

Table 6: RHIC spin example schedule, 10 physics weeks per year.

Fiscal year	Spin Weeks	CME(GeV)	P	L(pb ⁻¹)	Remarks
2005	9	200	0.5	10-20	$A_{LL}(\pi^0, jet)$, also 500 GeV studies
2006-2007	10	200	0.7		AGS Cold Snake commissioned, NEG vacuum coating complete
2008-2009	10	200	0.7		Direct γ
2010-11	10	200	0.7		completes goal for 200 GeV running
20012-13	10	500	0.7		PHENIX muon arm trigger installed, eng. run
2014-2015	10	500	0.7		STAR forward tracker installed, W physics
2016-2017	10	500	0.7		
2018-2019	10	500	0.7		Completes 500 GeV goal

7 Summary (Gerry)

In this document we have described the RHIC spin research plan, responding to the request by the Department of Energy Office of Nuclear Physics. We were requested to cover 1) the science, 2) the requirements for the accelerator, 3) the resources that are needed and timelines, and 4) the impact of a constant effort budget to the program.

1) The science is presented in section 2. Here we have emphasized measuring gluon polarization and anti-quark polarization in the proton. RHIC will provide the first sensitive measurements of each. We believe this is an exciting program, which addresses the structure of matter.

2) The accelerator requirements are presented in section 3. We are well along in reaching the polarization requirement of 70%, and anticipate reaching this goal in 2006, for 200 GeV running. To reach this goal for 500 GeV running will require releveling the machine, which is planned. Reaching the luminosity goal will be challenging. We must store 2×10^{11} polarized protons in 110 rf bunches in each RHIC ring and collide them. Limits of betatron tune shift and of electron cloud formation will be tested. For the physics sensitivities presented, we have used a luminosity of 2/3 of the calculated maximum.

3) The required experiment resources are presented in section 4. The PHENIX and STAR detectors are complete for the gluon polarization program. Both need improvements to be ready for the W physics program. These are described in the section. For a "technically driven" program, where the improvements are funded and completed as proposed, the PHENIX detector will be ready for W physics in 2009, and the STAR detector in 2010.

There are also important planned upgrades for the heavy ion and spin programs that greatly extend the range of spin physics, and these are also described in section 4.

4) The impact of a constant effort budget is presented in section 5, where we compare the two

plans, as requested in the charge to the RHIC Spinplan Group:

"I ask that you consider two RHIC Spin running scenarios: 1) 5 spin physics data taking weeks per year (averaged over two years using the combined fiscal year concept); 2) 10 spin physics data taking weeks per year. These two scenarios will give appropriate indications of the physics goals that can be met over a period of years without involving the Group in difficult funding and cost scenarios that are not central to the calculation of physics accomplishments over time." (Appendix A)

The plan with 10 spin physics weeks per year, the technically driven plan, completes the gluon polarization measurements and the W physics measurements by 2012.

The plan with 5 spin physics weeks per year completes this program in 2019 or later. With this plan RHIC runs 25% of the year on average (we assume 10 spin physics weeks per two year cycle).

Acknowledgments

Based on work over many years by accelerator physicists, experimenters and theorists in the RHIC Spin Collaboration and colleagues in the STAR, PHENIX, BRAHMS, and pp2pp experiments.

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